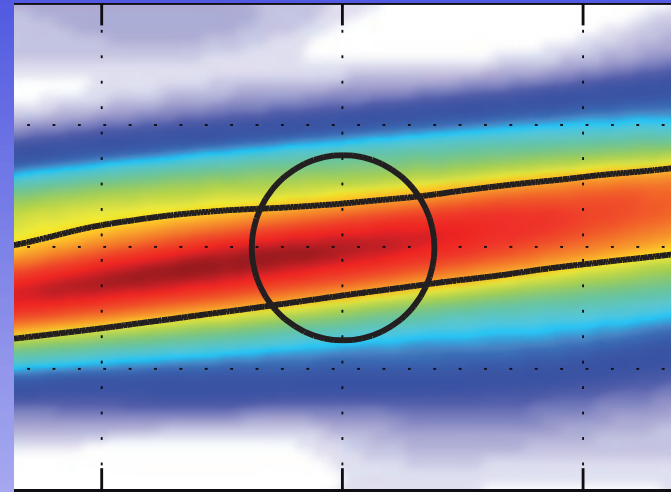
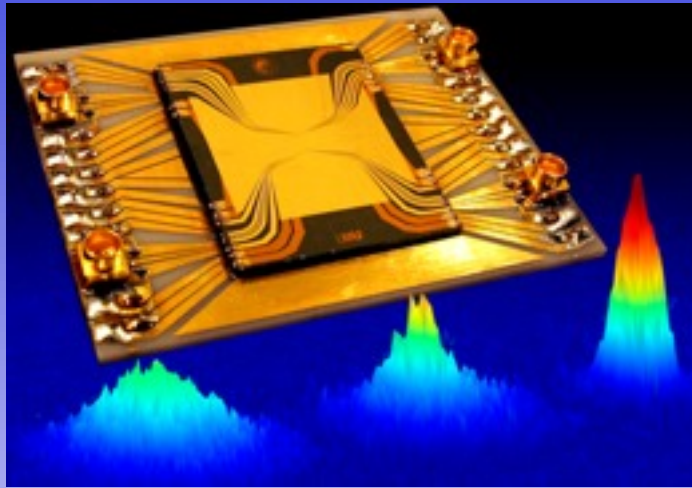


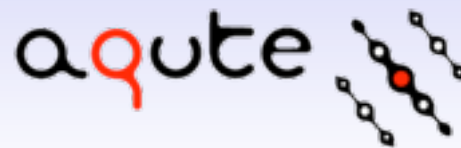
# Atom-chip-based generation of entanglement for quantum metrology



Philipp Treutlein

Max F. Riedel, Pascal Böhi, Jad C. Halimeh,  
Roman Schmied, and Theodor W. Hänsch

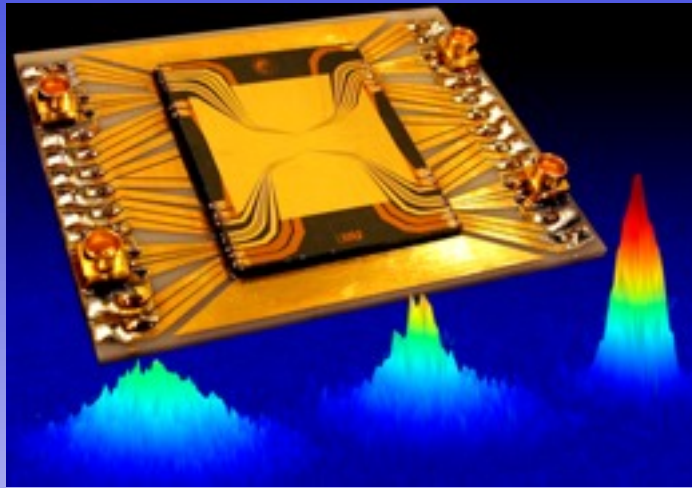
theory collaboration: Yun Li and Alice Sinatra (LKB/ENS)



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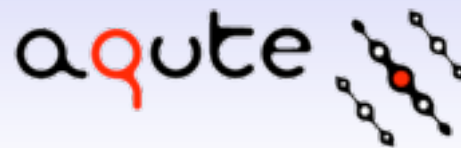


see poster by  
Alice Sinatra

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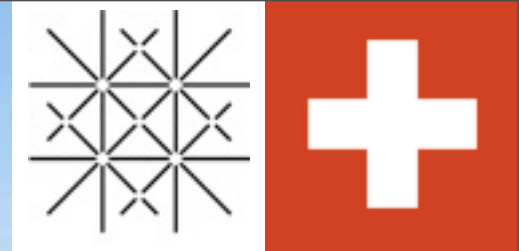


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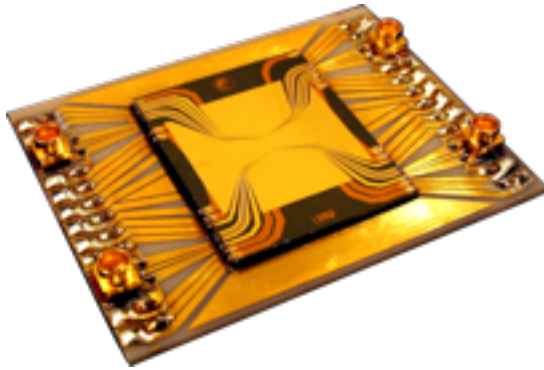




# September 2010: move to the UNIVERSITY OF BASEL Switzerland



Tuesday, June 8, 2010



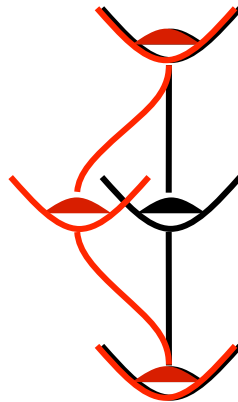
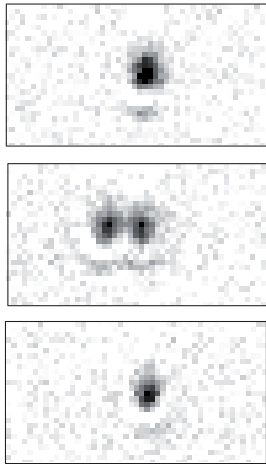
## Ultracold atoms on atom chips

- basic principles
- experimental setup

## Chip-based atomic clocks and interferometers

- BEC interferometer with internal-state labeling

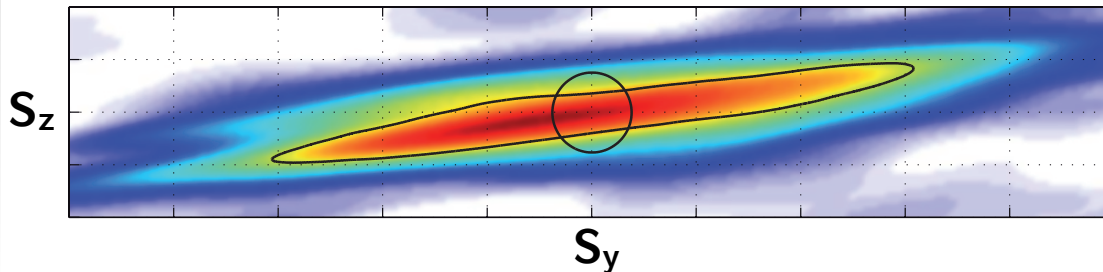
P. Böhi et al., *Nature Physics* 5, 592 (2009).



## Spin-squeezing and multi-particle entanglement

- control of collisions in state-dependent potential
- spin-squeezing on the clock transition
- Wigner function reconstruction
- experimental test of multi-particle entanglement

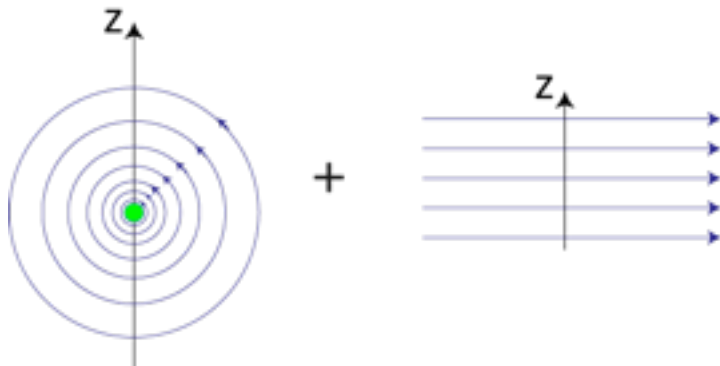
M. F. Riedel et al.,  
*Nature* 464, 1170 (2010).





# Atom chips

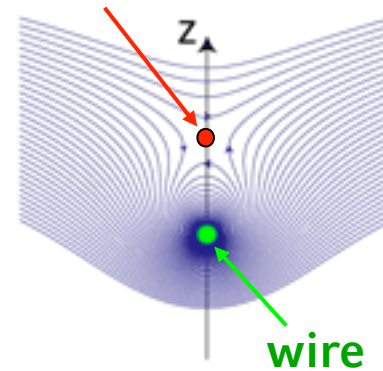
## Basic principle: the wire trap



magnetic field  
of wire

homogeneous  
field

atoms are trapped in  
B-field minimum



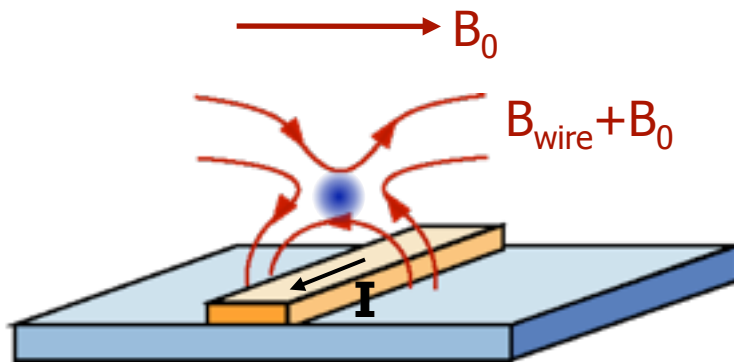
trap (2D)

## Zeeman interaction

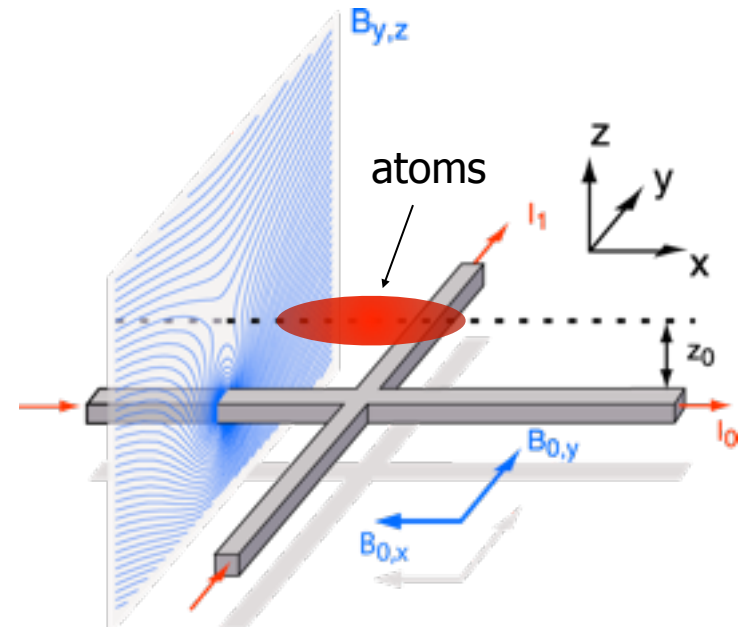
$$H_Z = -\mu \cdot B$$

## Wire field gradient

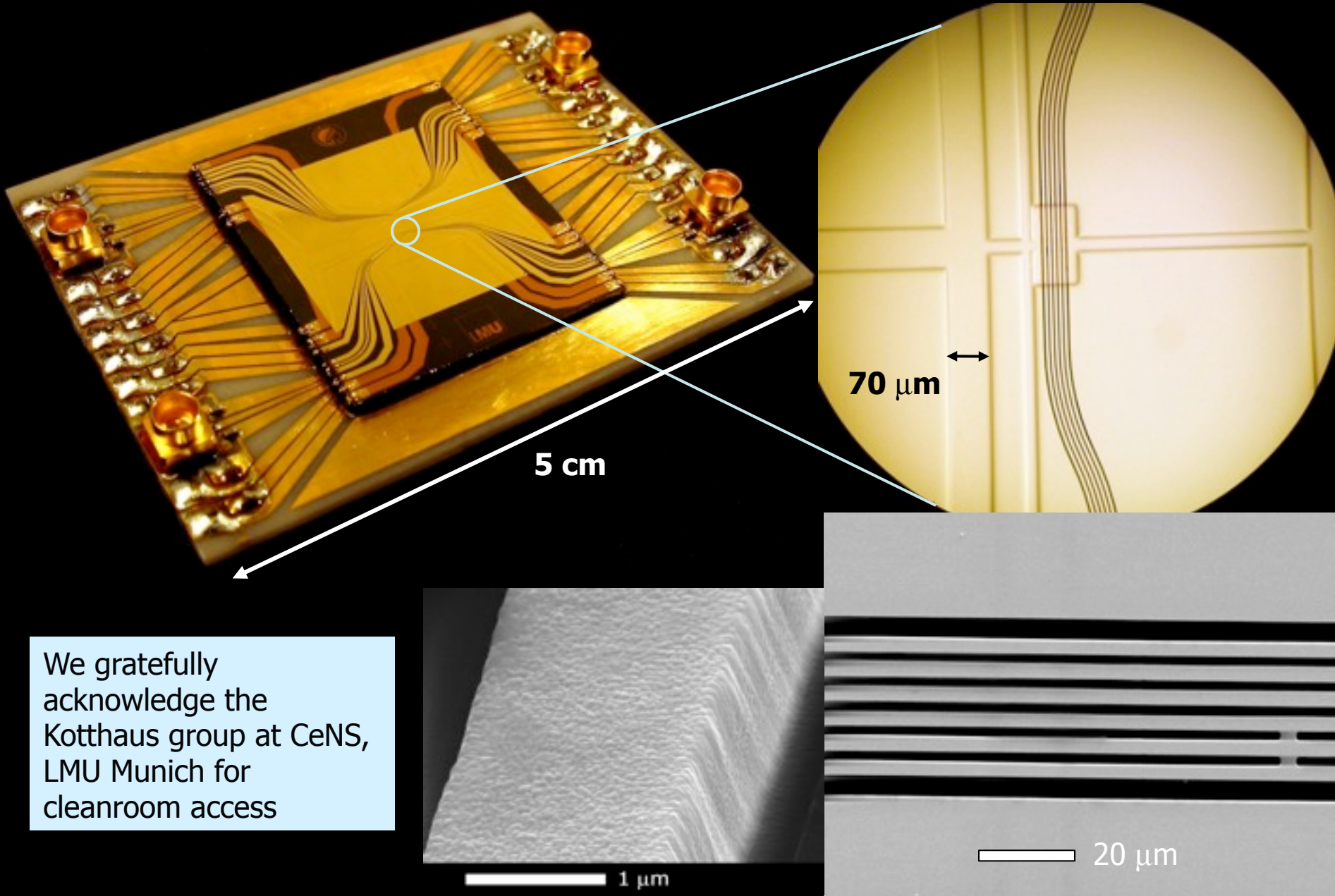
$$\frac{\partial B}{\partial z} = -\frac{\mu_0}{2\pi} \cdot \frac{I}{z^2}$$



high trap frequencies: up to 1 MHz  
small atom-surface distance: < 1  $\mu\text{m}$  possible

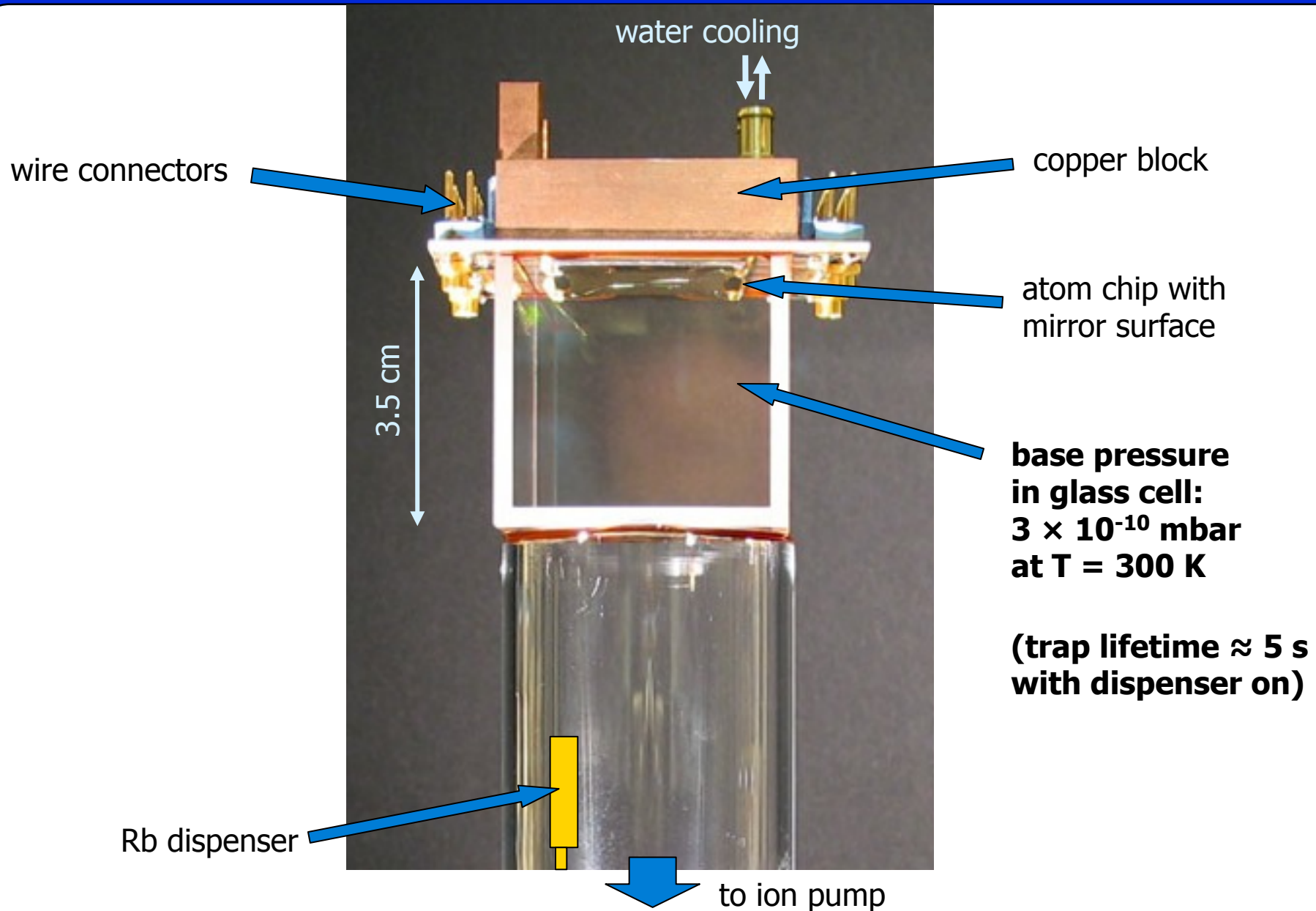


# State of the art: multi-layer atom chips



We gratefully acknowledge the Kotthaus group at CeNS, LMU Munich for cleanroom access

# Compact glass cell vacuum chamber





# Production of Bose-Einstein condensates

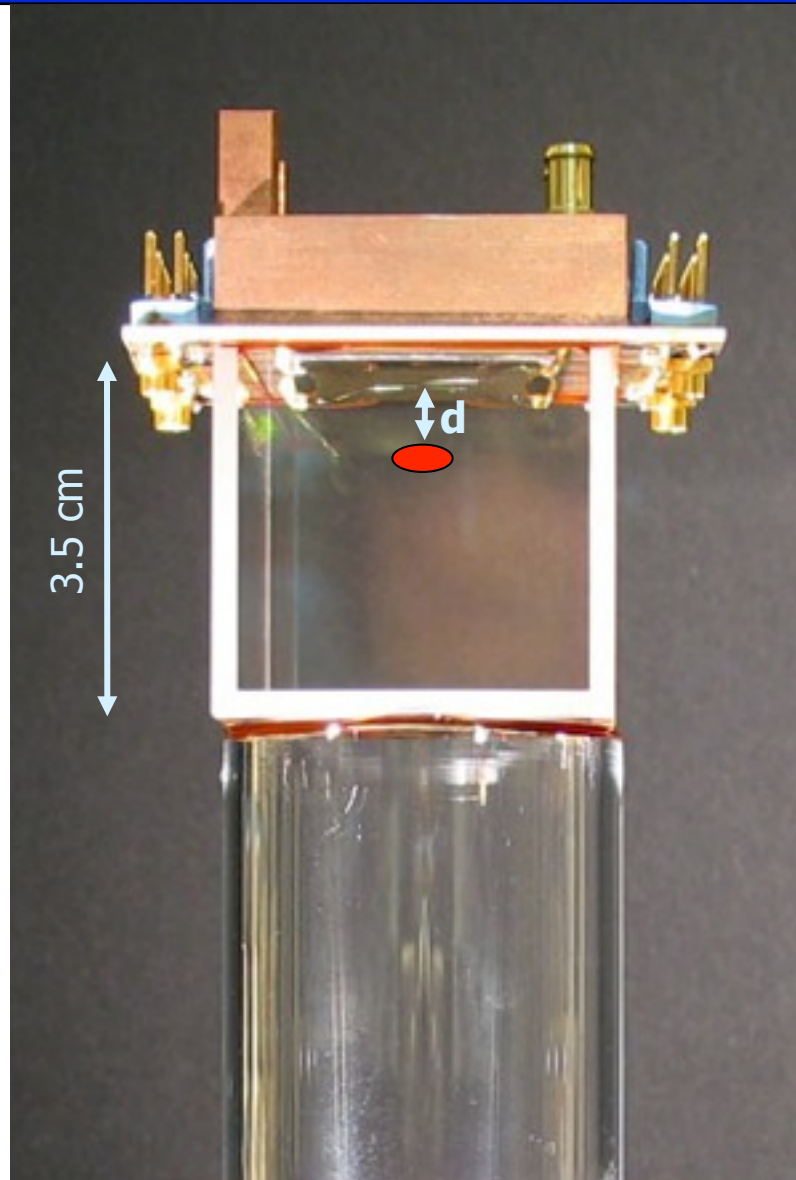
## BEC sequence:

- mirror-MOT
- optical molasses
- optical pumping
- magnetic trap
- transport atoms
- evaporative cooling to BEC

**all inside the same glass cell**

pressure:  
 $3 \times 10^{-10}$  mbar

experimental  
cycle: 10 s



# Production of Bose-Einstein condensates

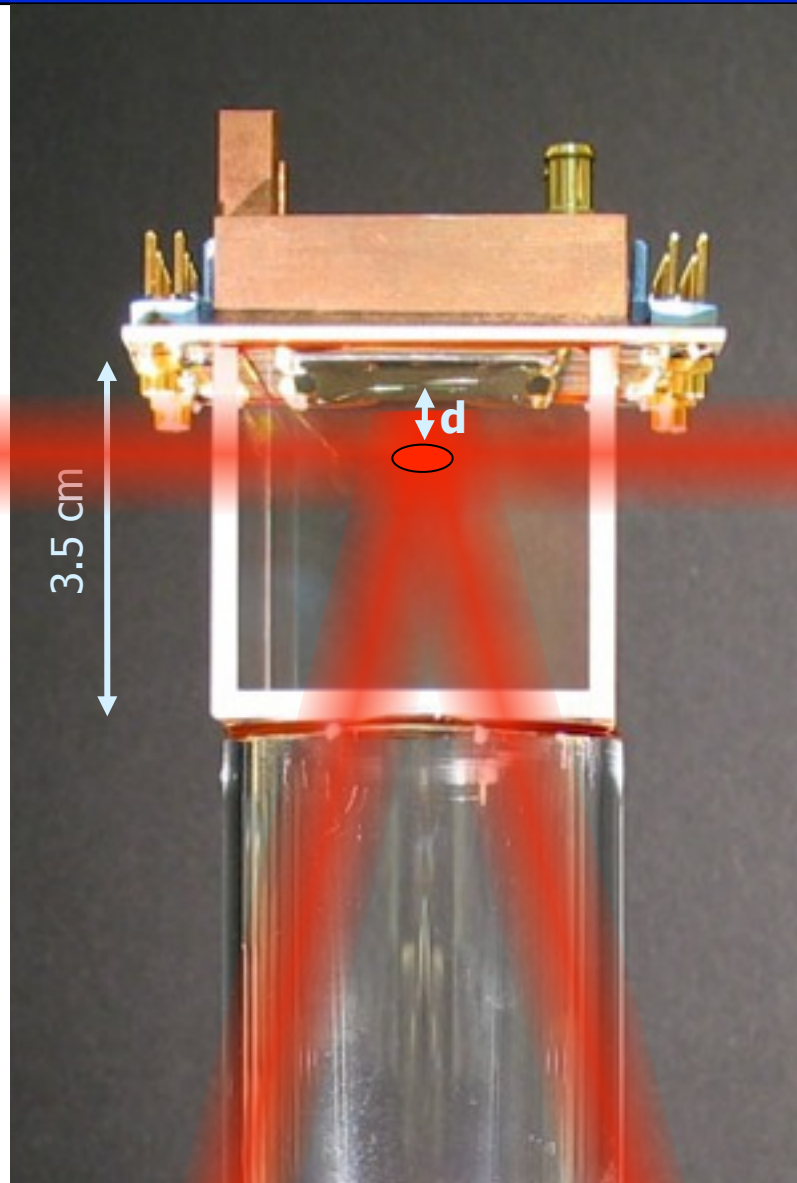
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 $3 \times 10^{-10}$  mbar

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# Absorption imaging

detection beam

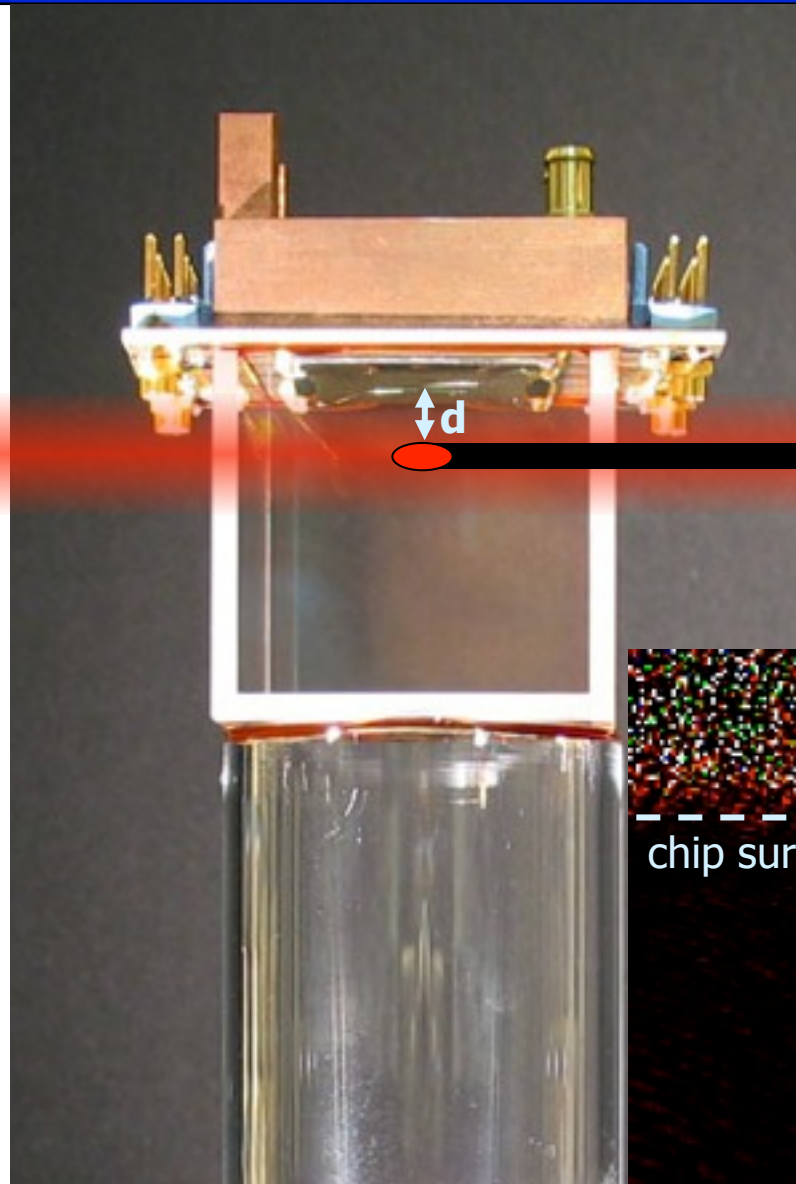
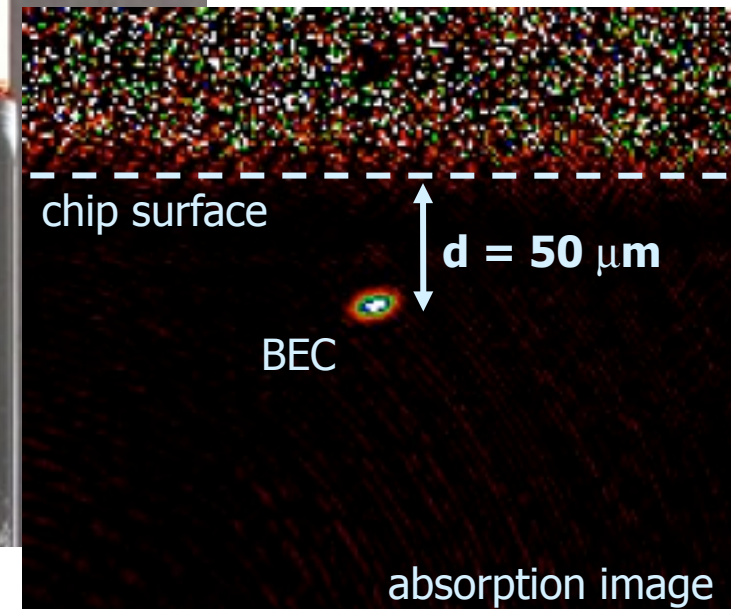
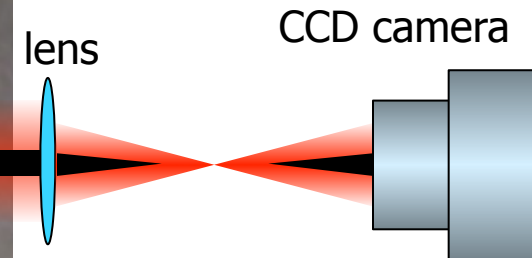


image shadow cast by atoms onto CCD camera

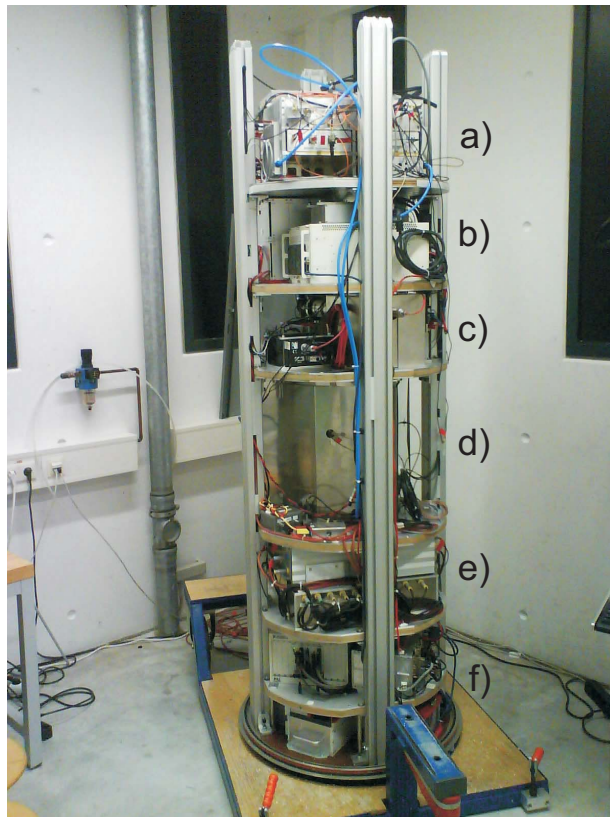




# Portable atom-chip setups

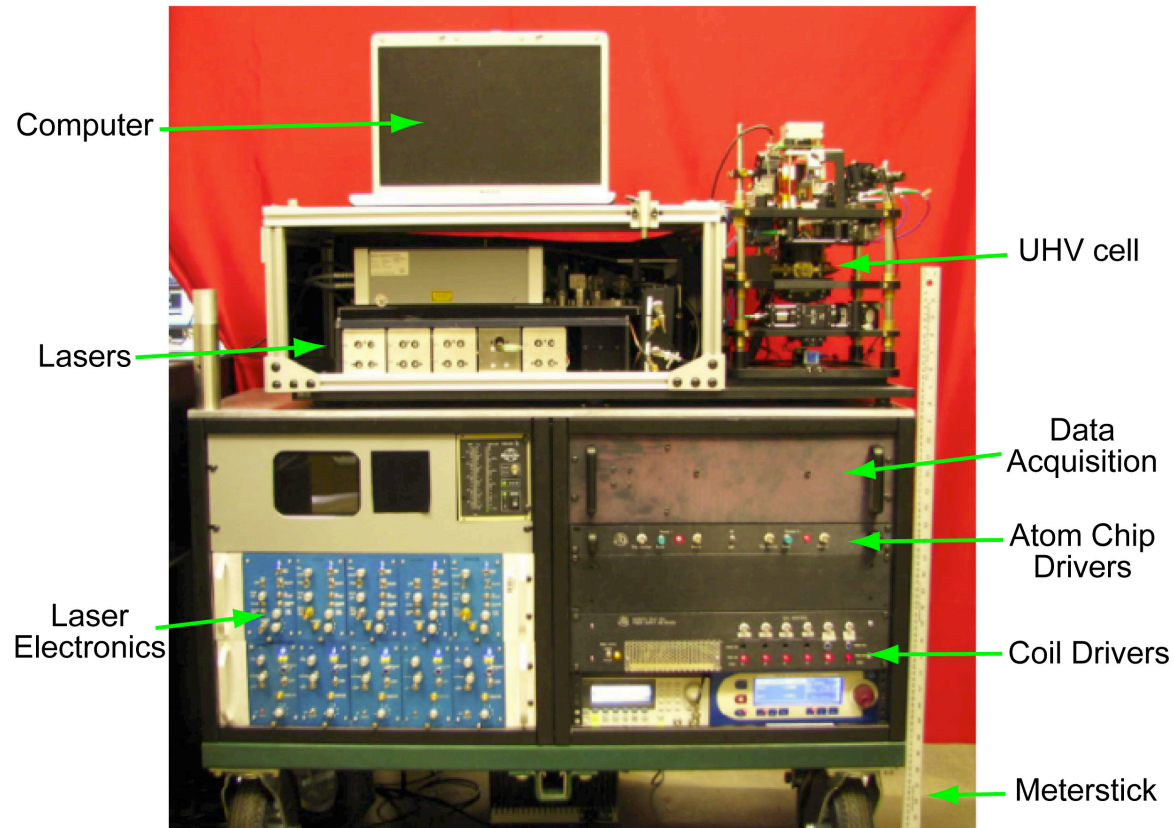
## QUANTUS project

A. Vogel et al.,  
Appl. Phys. B 84, 663 (2006).



## D. Anderson's group, Boulder

D. M. Farkas et al., arXiv:0912.0533 (2009).

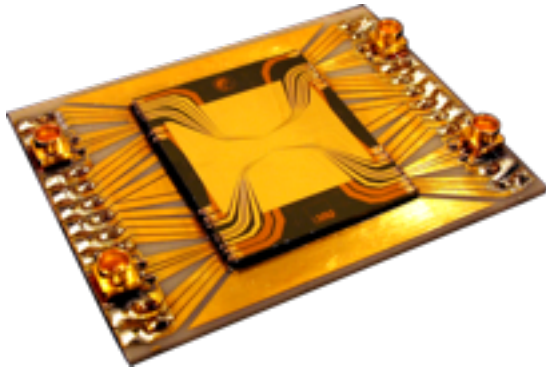


Key components  
commercially available:



ColdQuanta

[www.coldquanta.com](http://www.coldquanta.com)



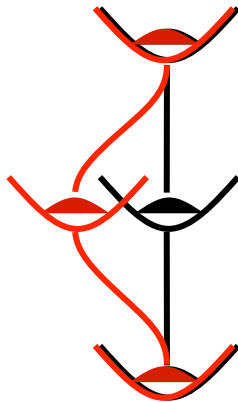
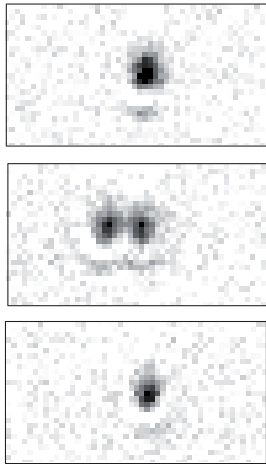
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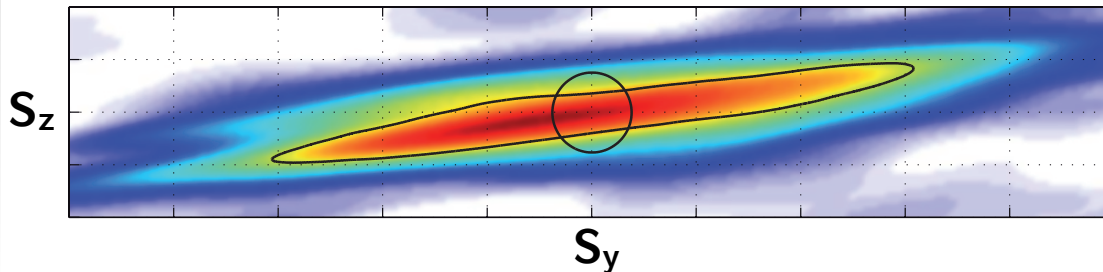
P. Böhi et al., *Nature Physics* 5, 592 (2009).



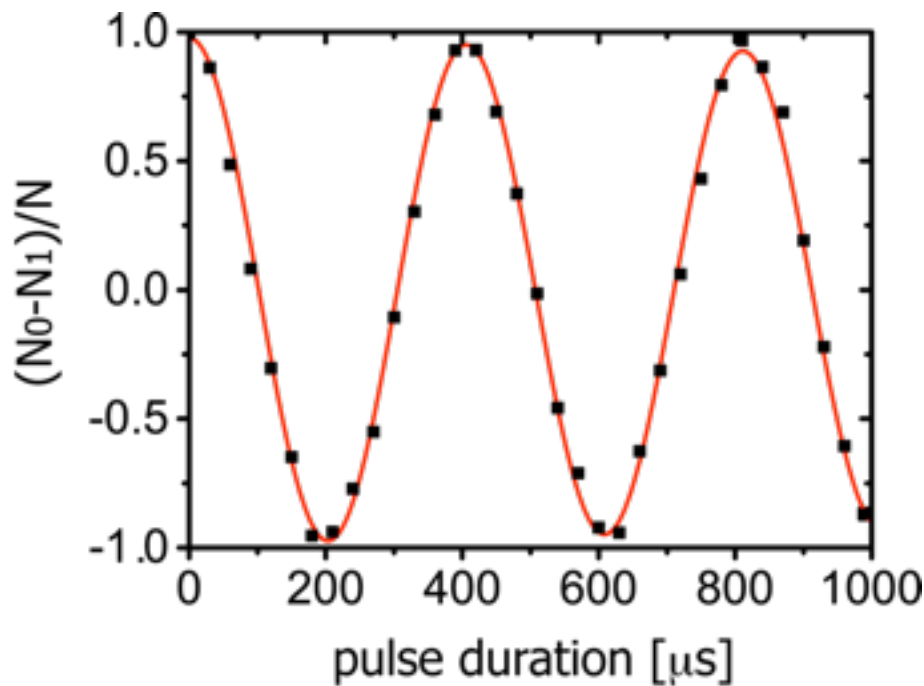
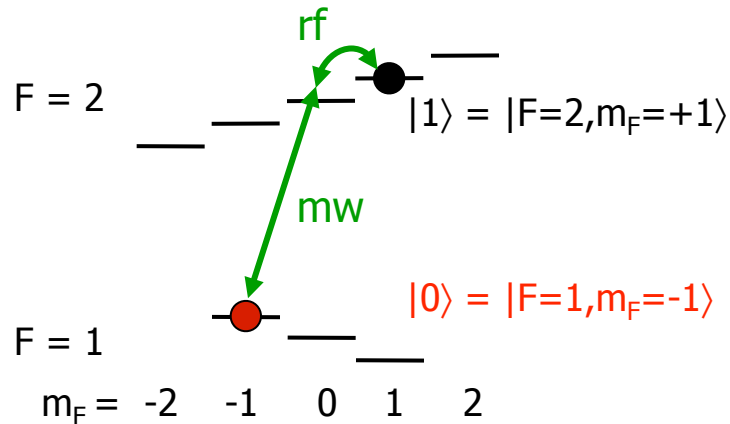
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M. F. Riedel et al.,  
*Nature* 464, 1170 (2010).



# Clock states of $^{87}\text{Rb}$ in a magnetic trap



## Qubit / clock states of $^{87}\text{Rb}$

- both magnetically trappable
- nearly identical potentials
- coherence lifetime  $> 5$  s possible (chip-based clock, thermal atoms)
- limitation: 2-body loss in  $|2,1\rangle$  (long lifetime only at low density)

## Rabi oscillations on the two-photon transition

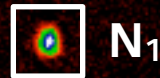
- atom number  $N = 1000$
- image both states in one shot  $\rightarrow$  normalize to total  $N$
- observed contrast  $C = 0.98$



# Detection system

$F = 2$   
(4.5 ms TOF)

detection noise:  
 $\pm 10$  atoms r.m.s.  
for  $N = 1000$



Ramp (30 ms) to relaxed detection trap (36 Hz/114 Hz),  
200  $\mu\text{m}$  from chip,  $B_0 = 3.0$  G

Resonant absorption imaging, pulse: 40  $\mu\text{s}$ ,  $I = 0.8 I_s$ ,  
cloud size 15x20  $\mu\text{m}^2$ ,  $\text{OD}_{\text{max}} = 1-2$

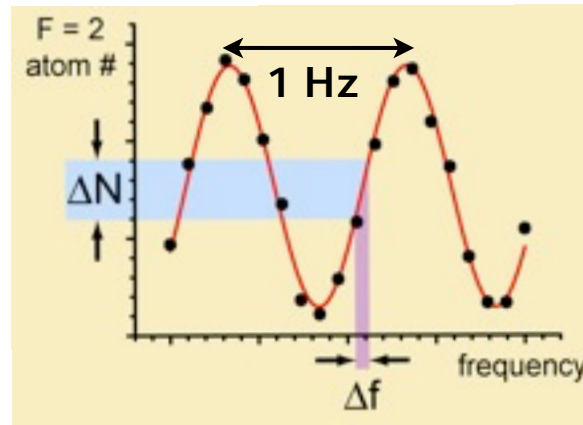
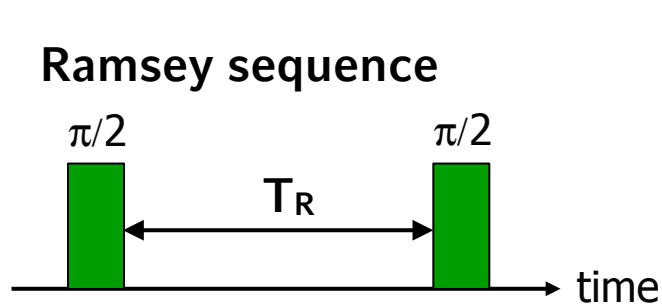
camera QE=0.9, spatial resolution 4  $\mu\text{m}$

$F = 1$   
(6.1 ms TOF)



# Chip-based atomic clock

**Proof-of-principle experiment:** P. Treutlein et al., PRL 92, 203005 (2004).



$T_R = 1\text{ s}$

thermal atoms

frequency stability  
(Allan deviation):  
 $1.7 \times 10^{-11}$  @ 1 s

**Dedicated precision experiment:**

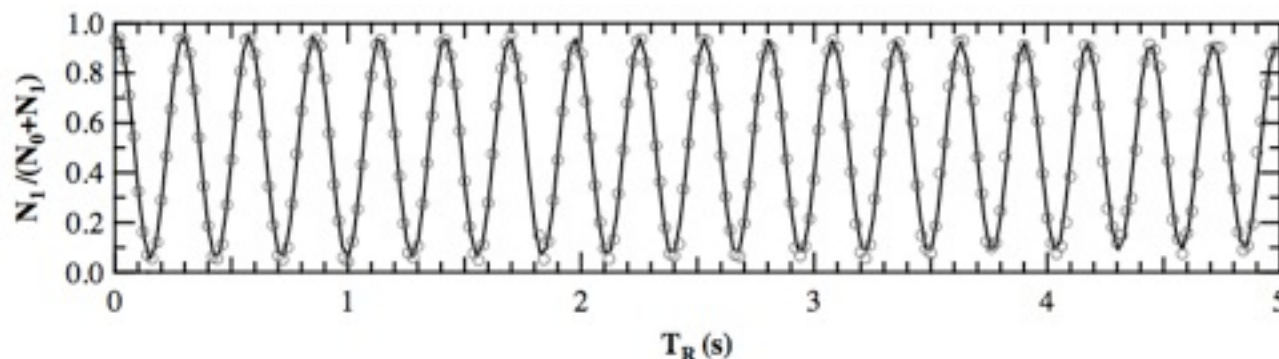
goal: stability  $10^{-13}$  @ 1 s

P. Rosenbusch/J. Reichel, Observatoire/SYRTE/LKB, Paris

C. Deutsch et al., arXiv:1003.5925 (2010).

- $10^4 - 10^5$  atoms
- thermal atoms or BEC
- magnetic shielding
- improved detection

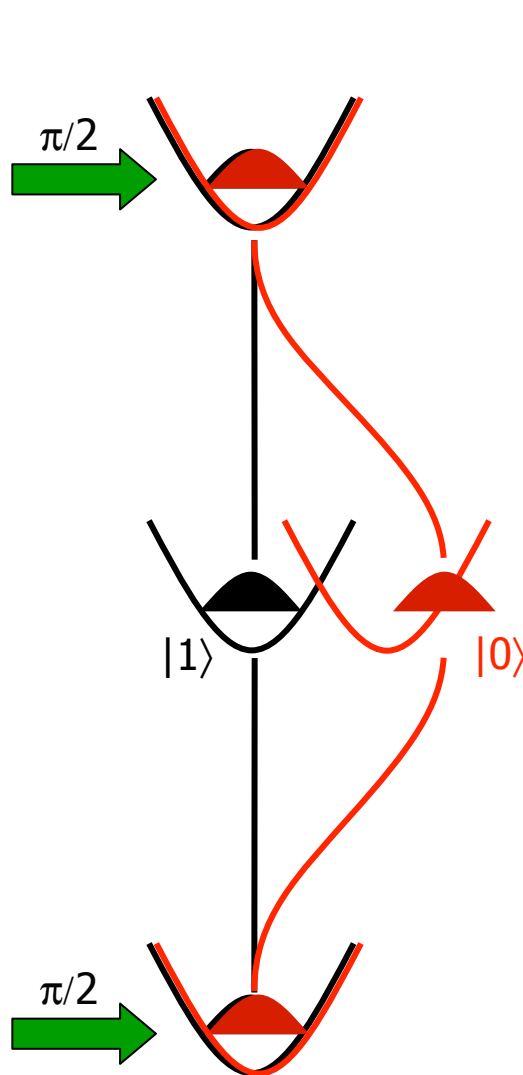
compact clock e.g. for  
satellite navigation



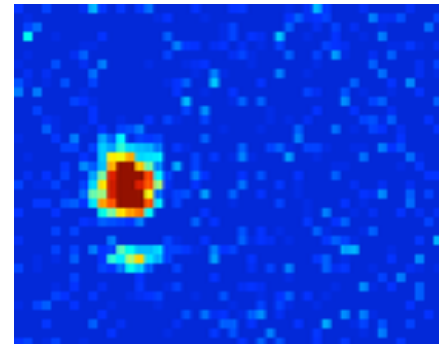
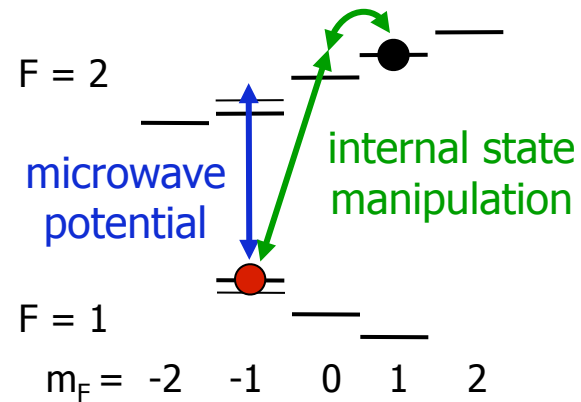


# A trapped-atom interferometer on a chip

microwave  
potential



time



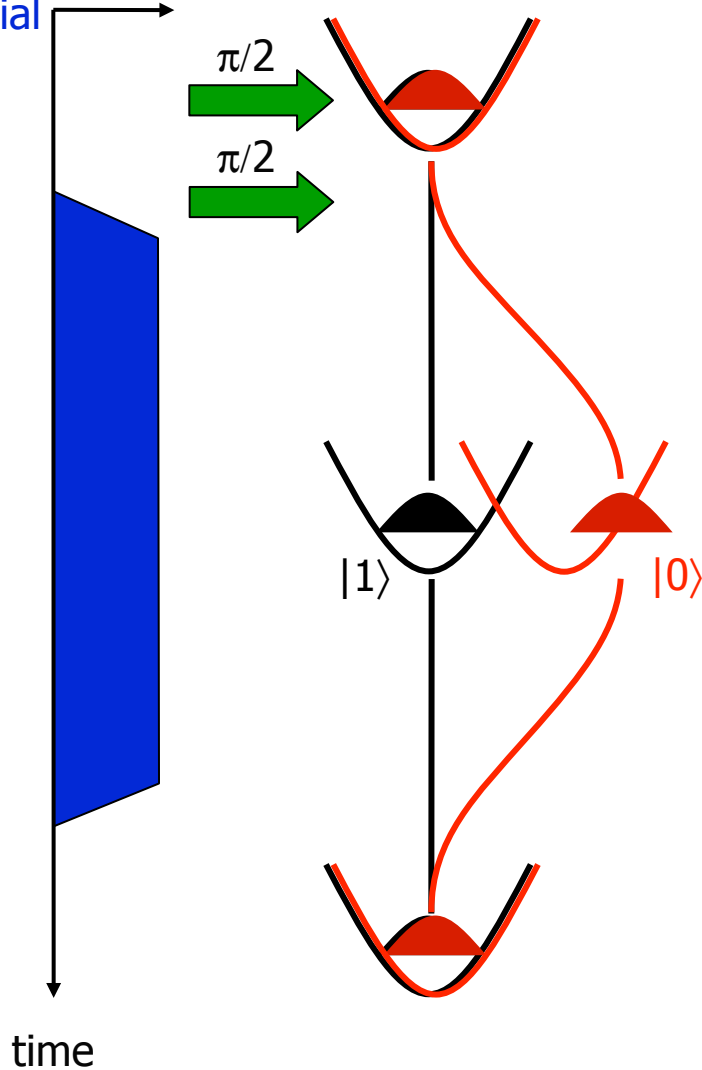
- movie shows in-situ images of BEC with 350 atoms during splitting
- detect both states ( $F=1$  and  $F=2$ )

P. Böhi et al., *Nature Physics* 5, 592 (2009).

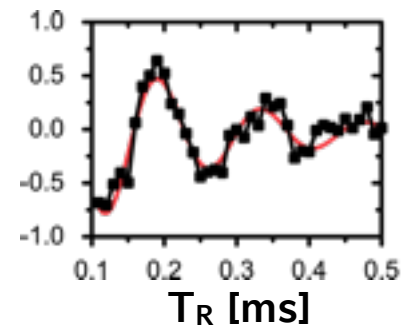


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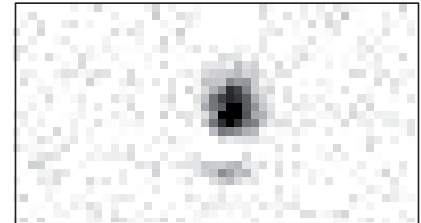
microwave  
potential



Ramsey fringes



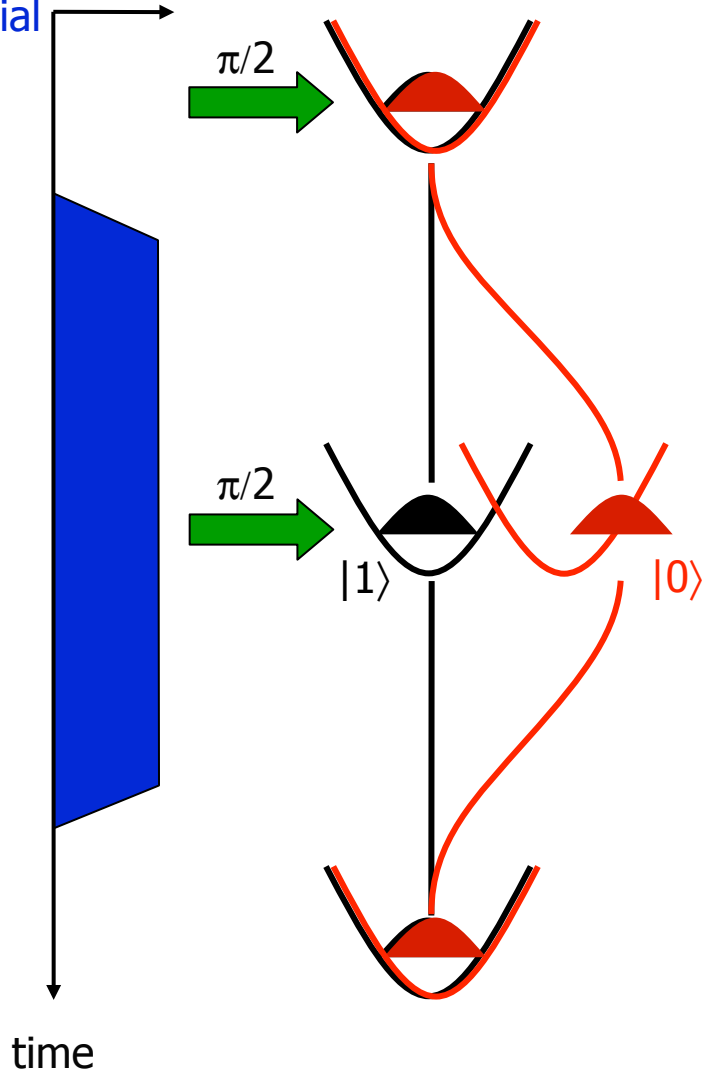
in-situ images



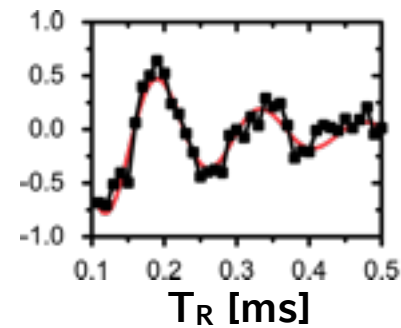
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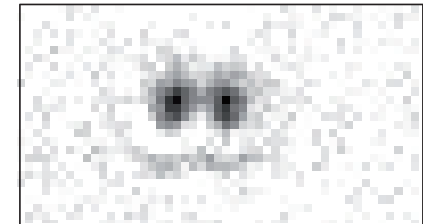
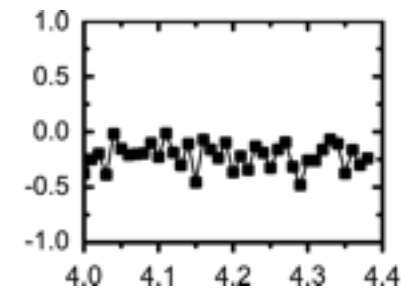
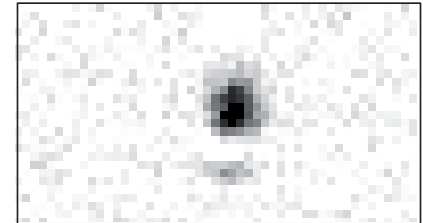
microwave  
potential



Ramsey fringes



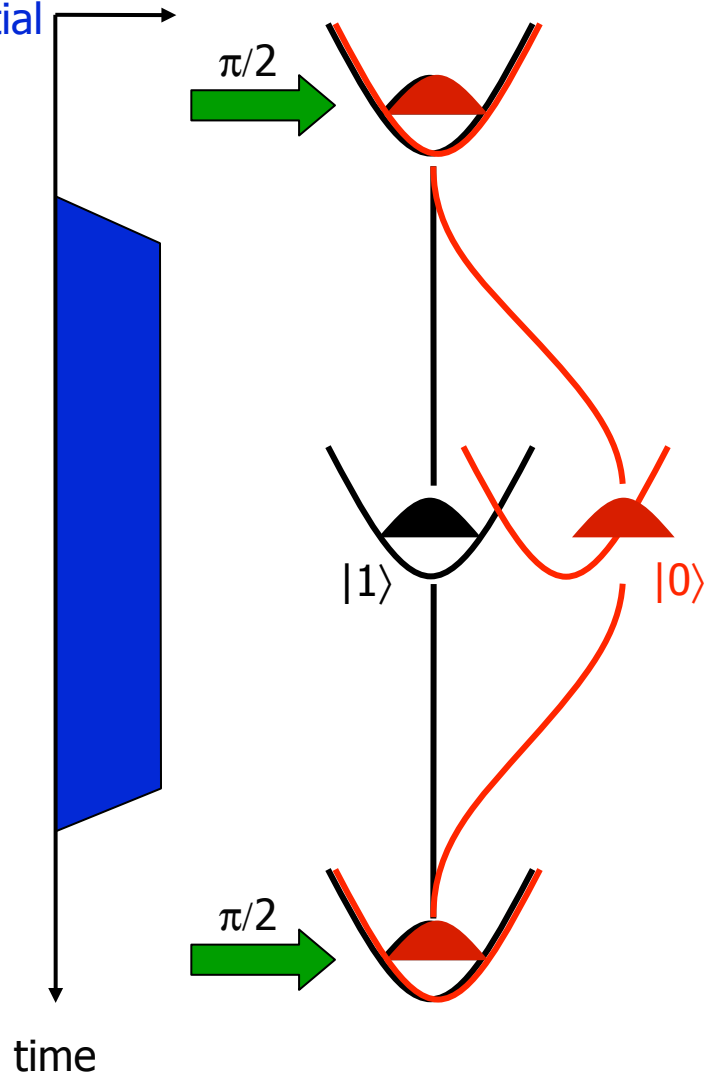
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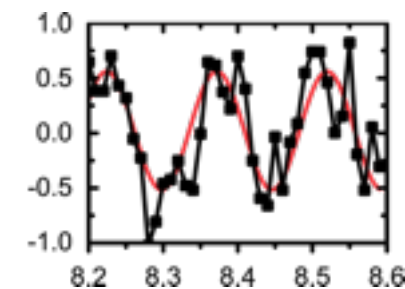
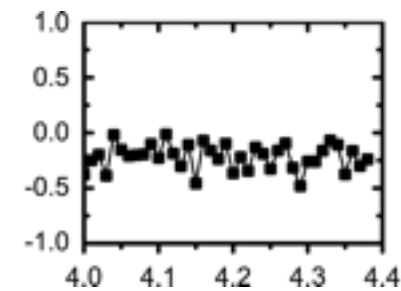
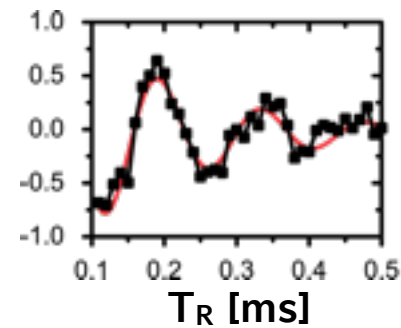
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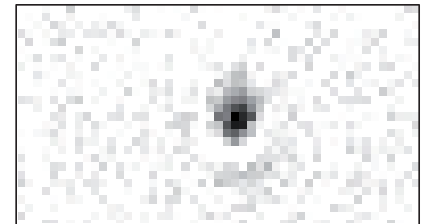
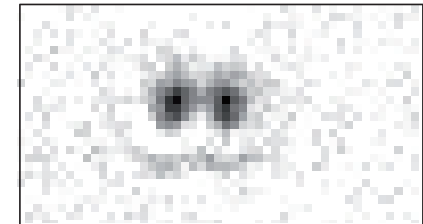
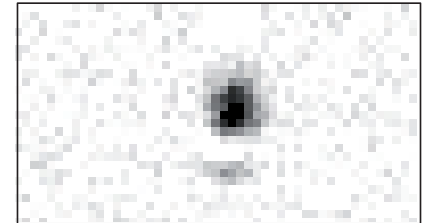
microwave  
potential



Ramsey fringes

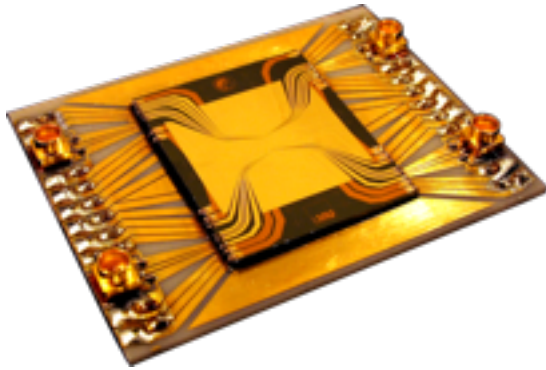


in-situ images



P. Böhi et al., Nature Physics 5, 592 (2009).





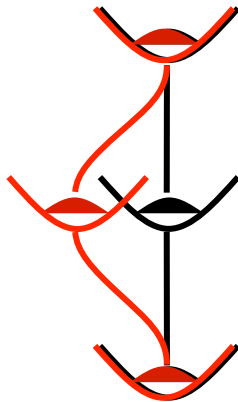
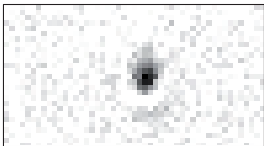
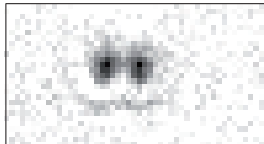
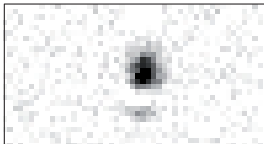
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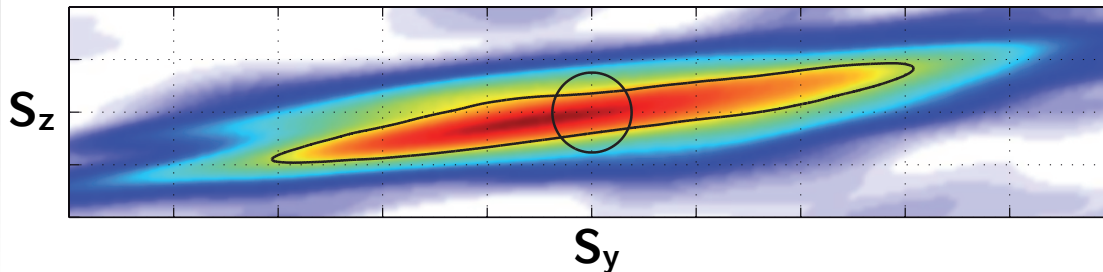
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M. F. Riedel et al.,  
*Nature* 464, 1170 (2010).



# Two-component BEC as a collective spin

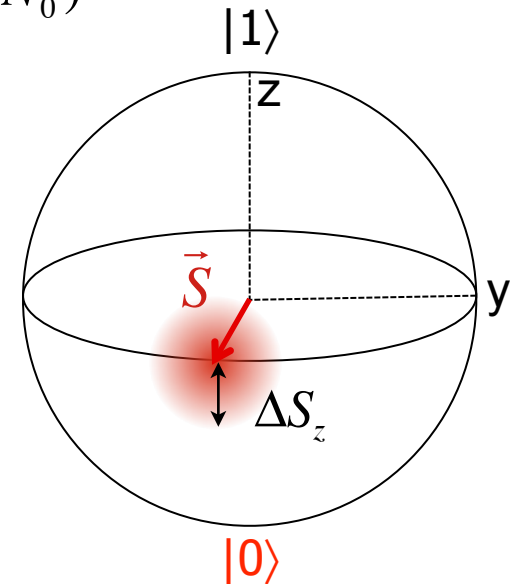
## BEC internal state: collective spin

$$N \times \begin{array}{l} \text{---} \bullet \text{---} |1\rangle \\ \text{---} \bullet \text{---} |0\rangle \end{array} \Rightarrow \vec{S} = \sum_{i=1}^N \vec{s}_i, \quad S = \frac{N}{2} \quad S_z = \frac{1}{2}(N_1 - N_0)$$

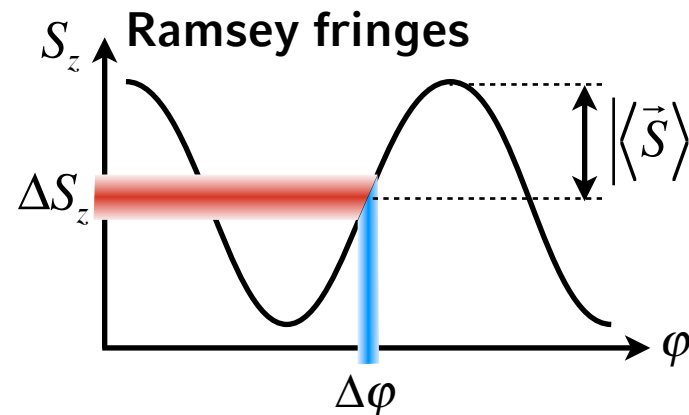
## Coherent spin state (product state):

$$|\Psi\rangle \sim (|0\rangle + |1\rangle)^{\otimes N}, \quad \bar{N}_0 = \bar{N}_1 = N/2$$

$$\Rightarrow |\langle S_x \rangle| = N/2, \quad \Delta S_z = \Delta S_y = \sqrt{N}/2$$



Quantum projection  
noise in a Ramsey  
interferometer



# Spin squeezing and entanglement generation

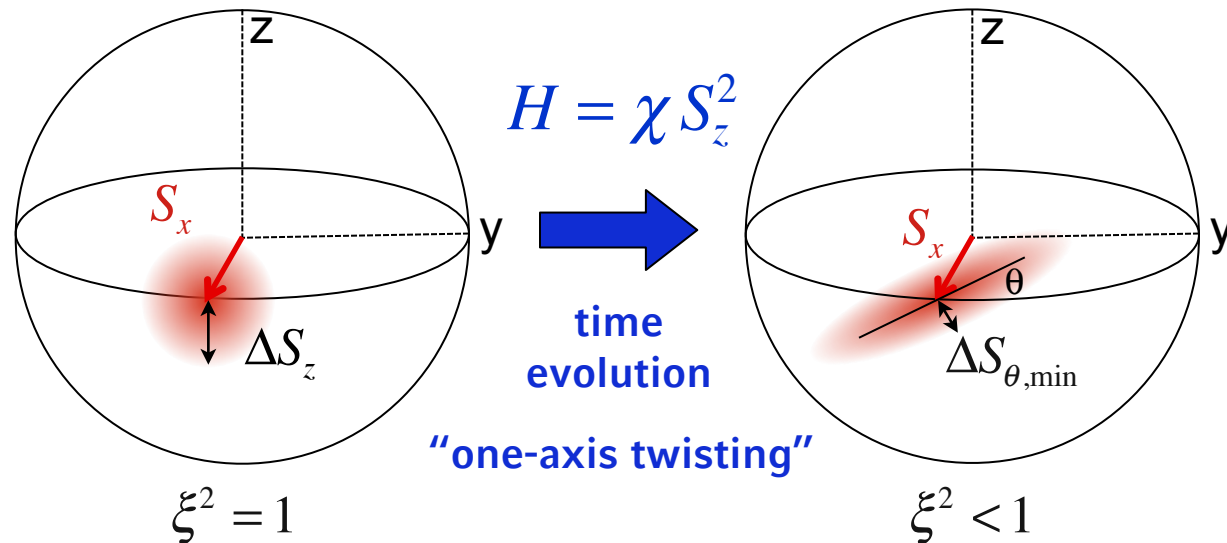
## Spin squeezing/ entanglement through nonlinear dynamics

Kitagawa, Ueda (1993)

Sørensen, Duan, Cirac,  
Zoller (2001)

Sinatra, Castin (2000,  
2008, 2009)

...



## Squeezing/entanglement parameter (Wineland, 1994):

$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

if  $\xi^2 < 1 \Rightarrow$

- useful resource for interferometry beyond standard quantum limit
- atoms entangled

to determine  $\xi$ , measure:

- minimum fluctuations  $\Delta S_{\theta, \min}$
- mean spin (Ramsey contrast)  $\langle S_x \rangle$

recent experiments: Oberthaler (BEC, double well 2008, int. state 2010)  
Polzik, Vuletic (thermal atoms, int. state 2008/09/10), ...



# Control nonlinearity by wave-function engineering

**Hamiltonian:**  
(two-mode model)

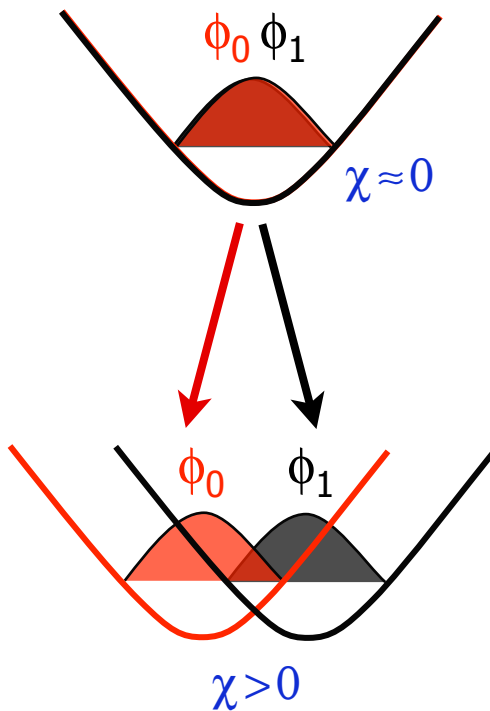
$$H = \delta S_z + \Omega_R S_\phi + \chi S_z^2 \quad \text{nonlinearity due to collisions}$$

$$\chi \sim a_{00} \int |\phi_0|^4 dr^3 + a_{11} \int |\phi_1|^4 dr^3 - 2a_{01} \int |\phi_0|^2 |\phi_1|^2 dr^3$$

(simplification: BEC mode functions  $\phi_0, \phi_1$  independent of  $N_0, N_1$ )

but for  $^{87}\text{Rb}$ :  $a_{00} \sim a_{11} \sim a_{01} \Rightarrow \chi \approx 0$

no convenient Feshbach resonance in magnetic trap



**use state-dependent potential to control interactions via wave function overlap**

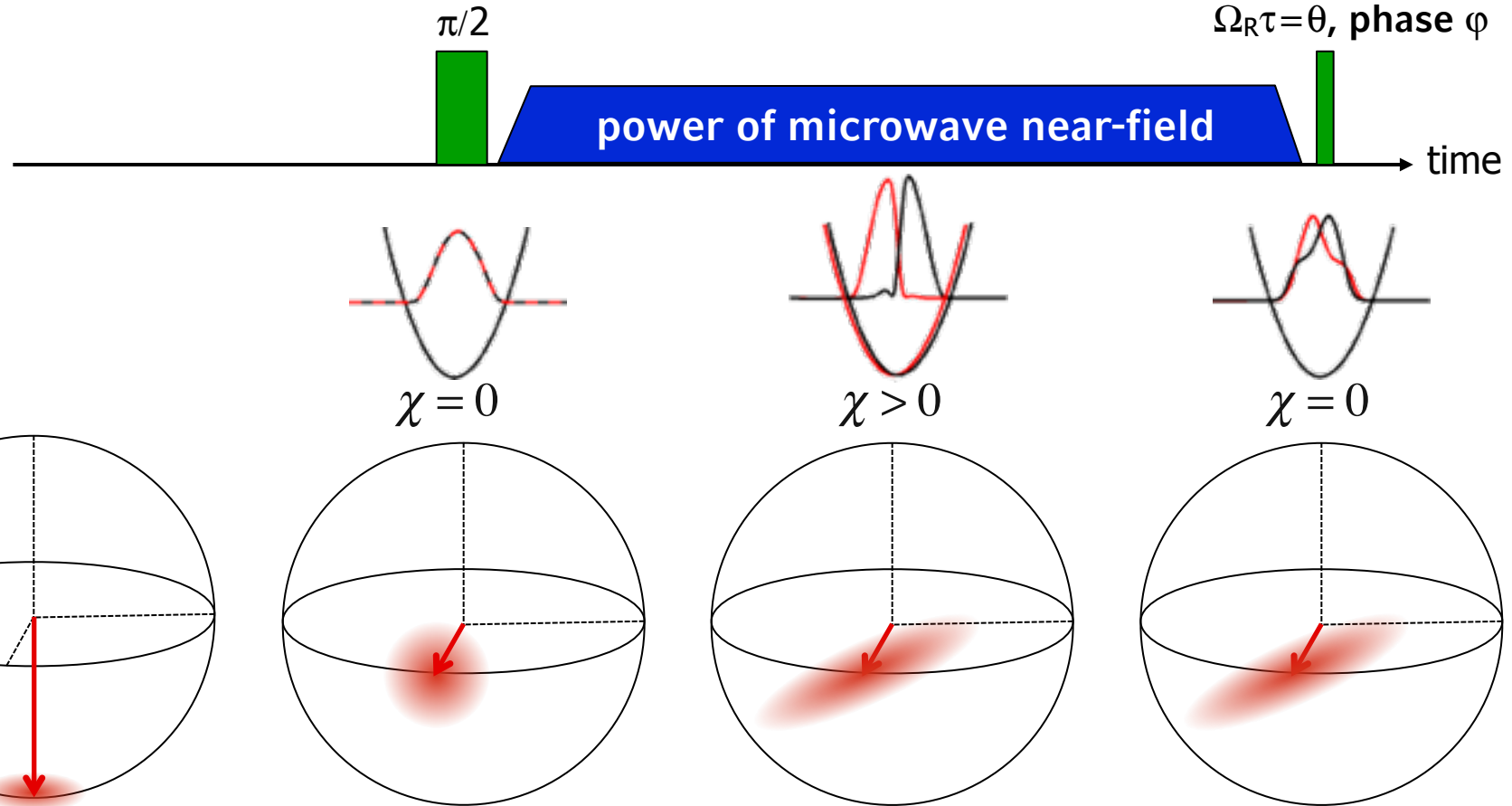
(turn nonlinearity on for well-defined time, avoid oversqueezing)

Y. Li, P. Treutlein, J. Reichel, A. Sinatra, Eur. Phys. J. B 68, 365 (2009).

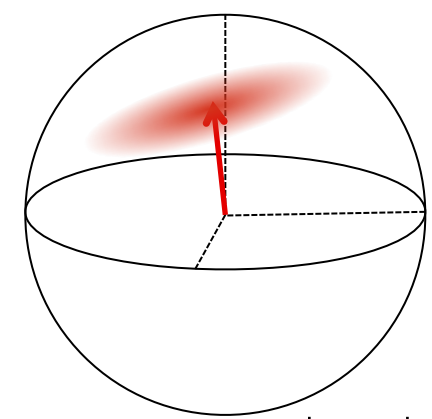
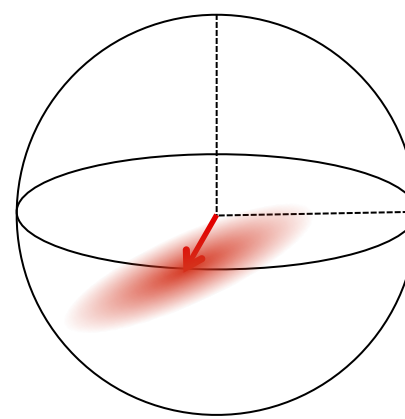
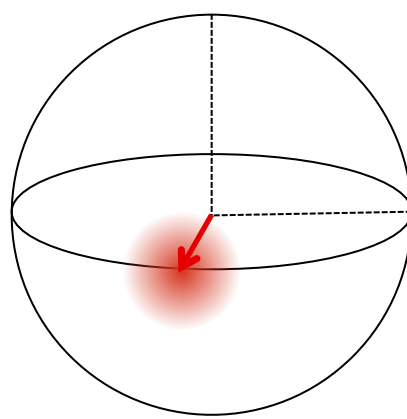
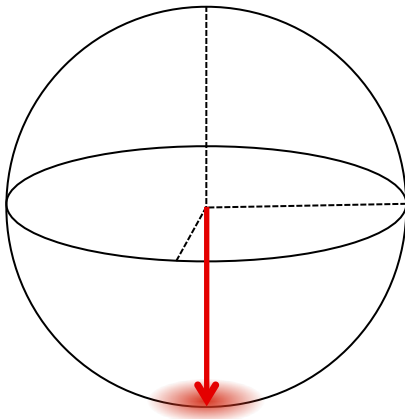
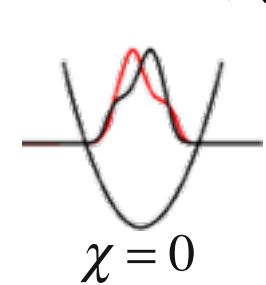
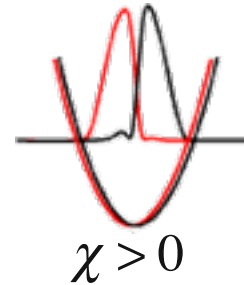
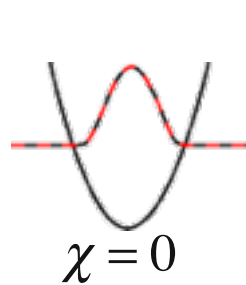
related idea for BEC in TOF: U. Poulsen and K. Mølmer, PRA 65, 033613 (2002).

related ideas for QIP with single atoms: T. Calarco et al., PRA 61, 022304 (2002).

# Experimental sequence for spin squeezing



# Experimental sequence for spin squeezing

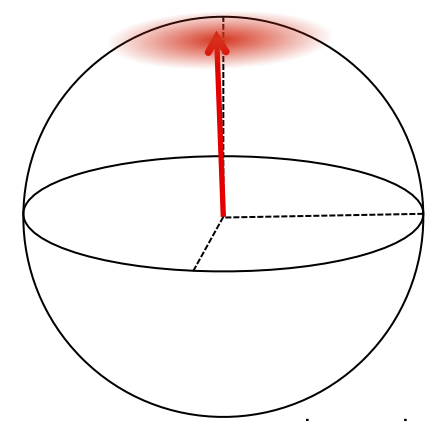
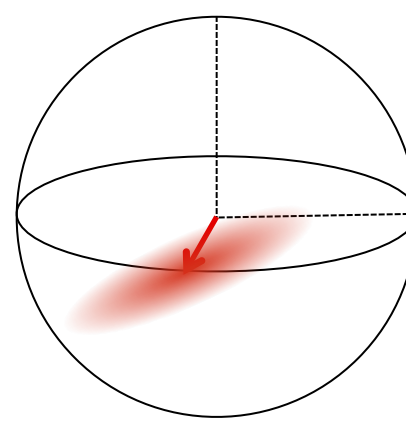
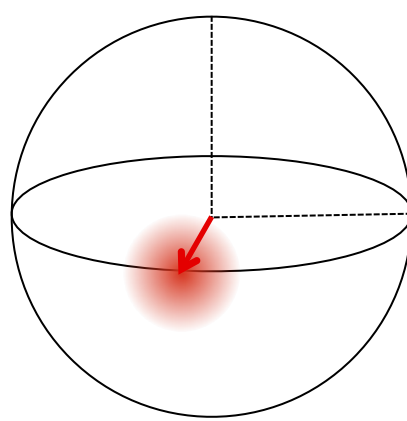
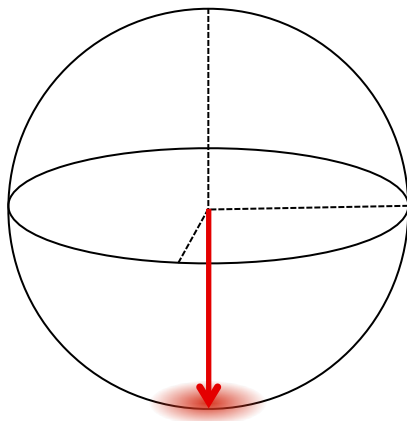
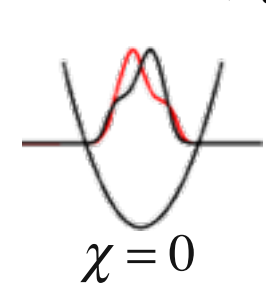
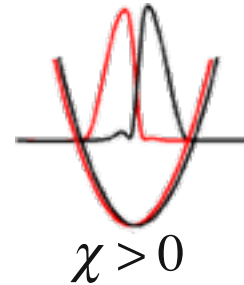
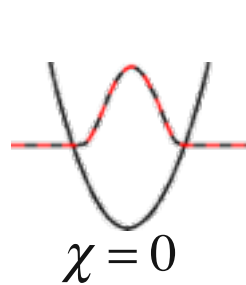


measure  $|\langle S_x \rangle|$   
(Ramsey contrast)

$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$



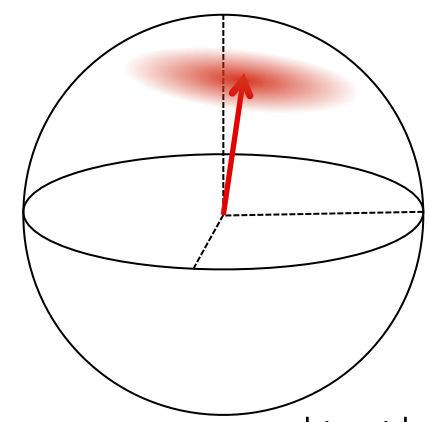
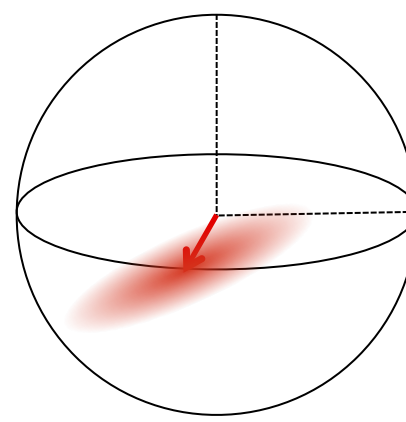
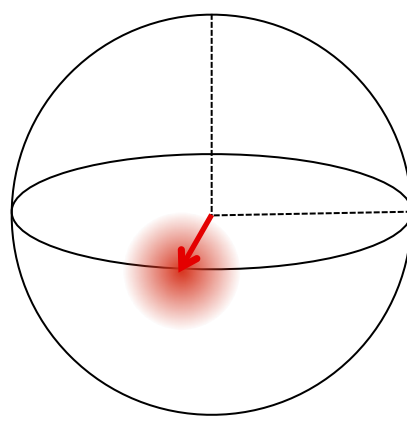
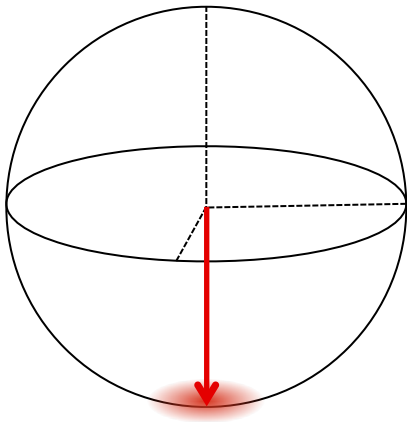
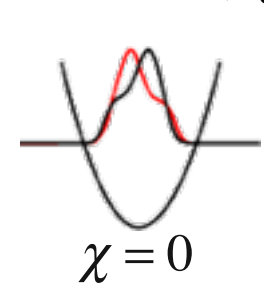
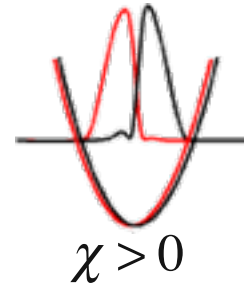
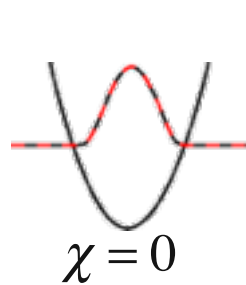
# Experimental sequence for spin squeezing



measure  $|\langle S_x \rangle|$   
(Ramsey contrast)

$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

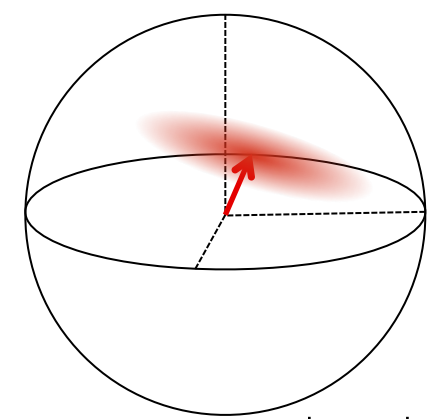
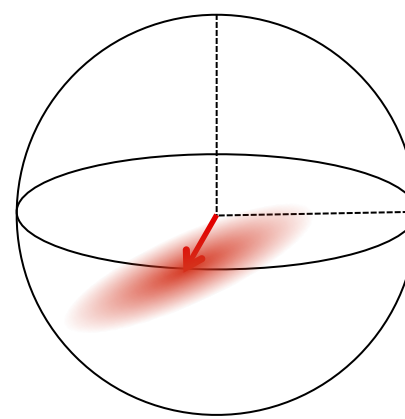
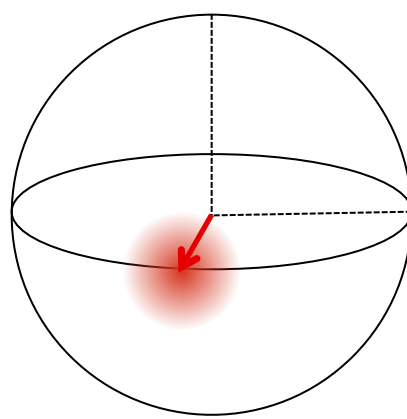
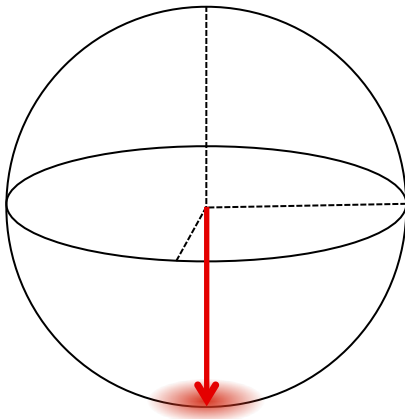
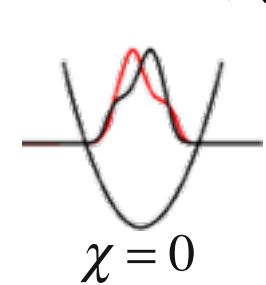
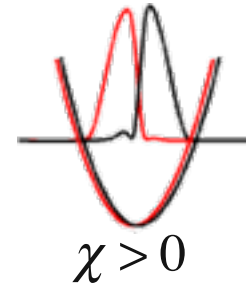
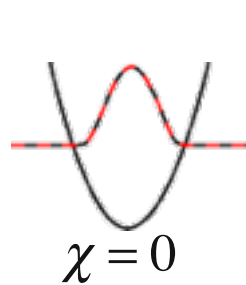
# Experimental sequence for spin squeezing



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# Experimental sequence for spin squeezing

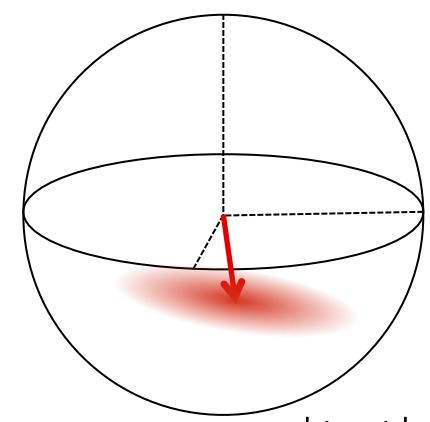
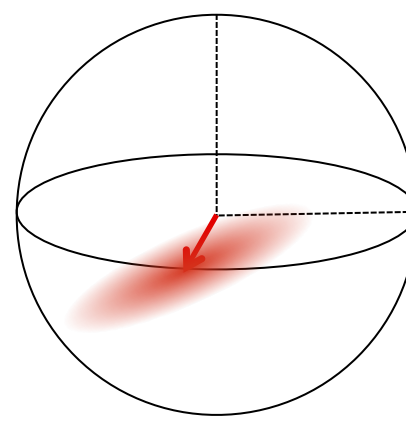
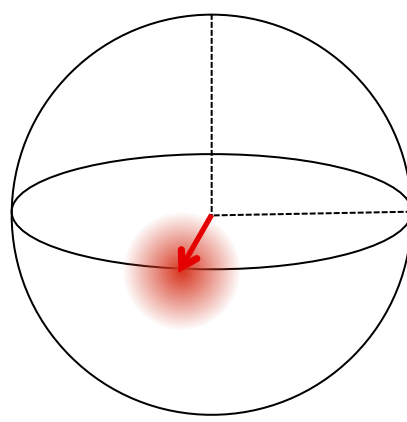
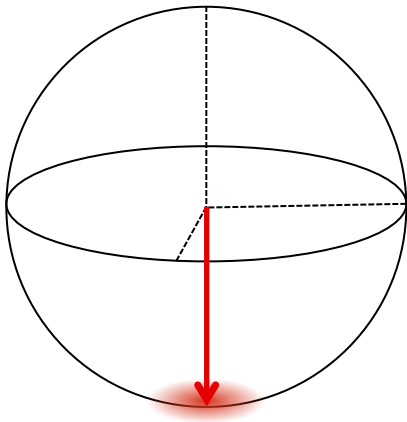
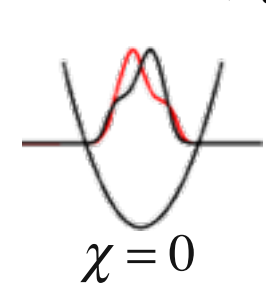
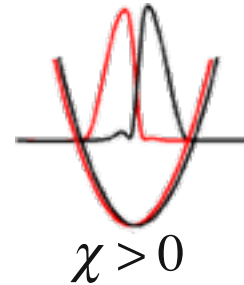
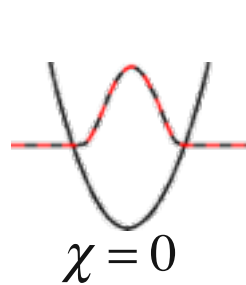


measure  $|\langle S_x \rangle|$   
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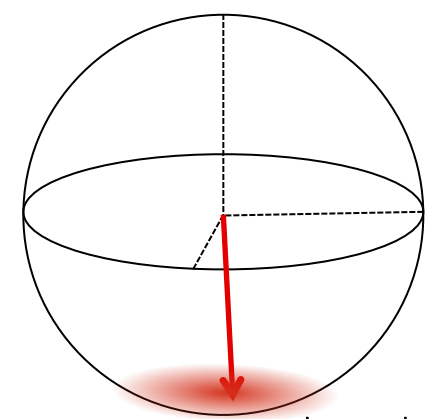
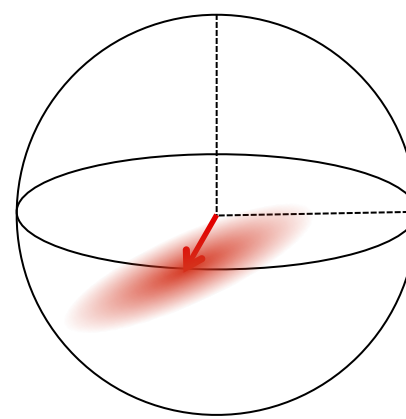
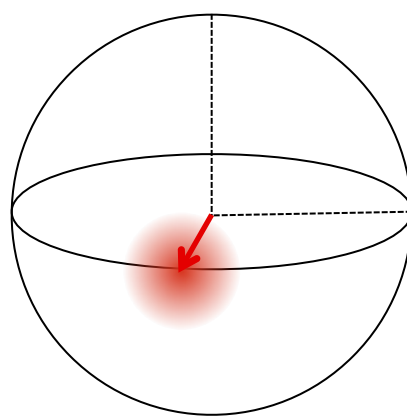
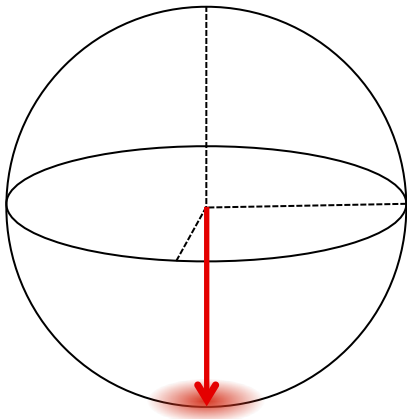
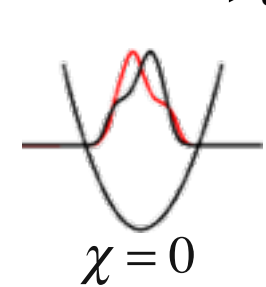
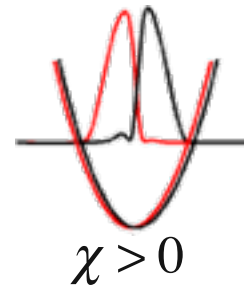
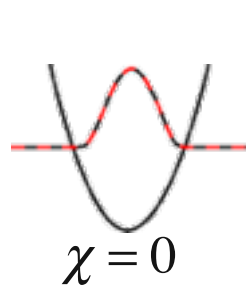
# Experimental sequence for spin squeezing



measure  $|\langle S_x \rangle|$   
(Ramsey contrast)

$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

# Experimental sequence for spin squeezing

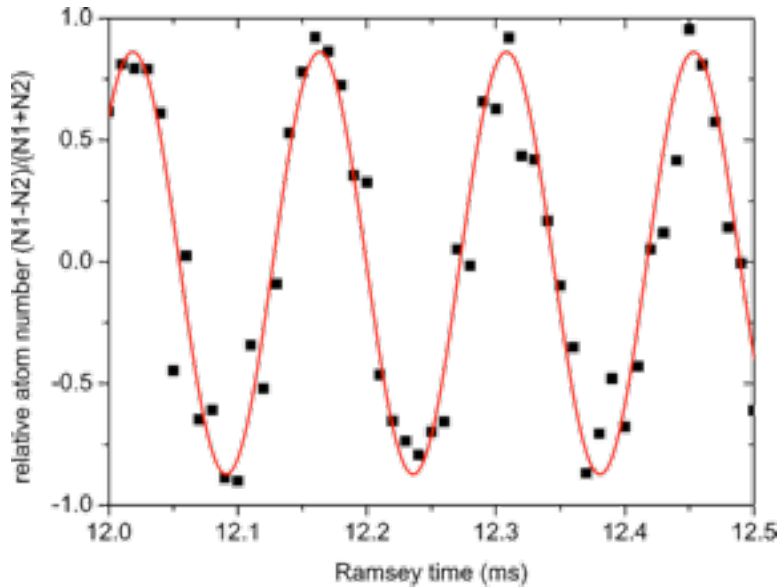


measure  $|\langle S_x \rangle|$   
(Ramsey contrast)

$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

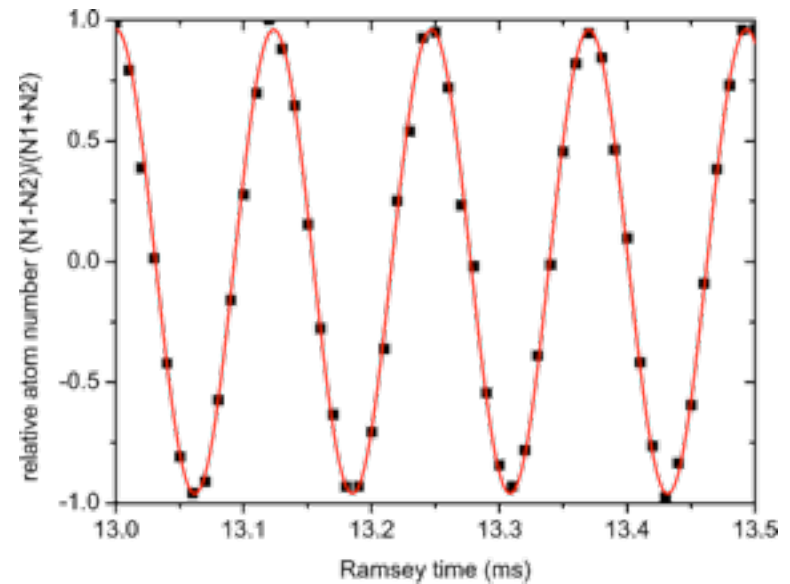
# Contrast of Ramsey fringes

**with splitting** ( $\chi > 0$ , squeezing)



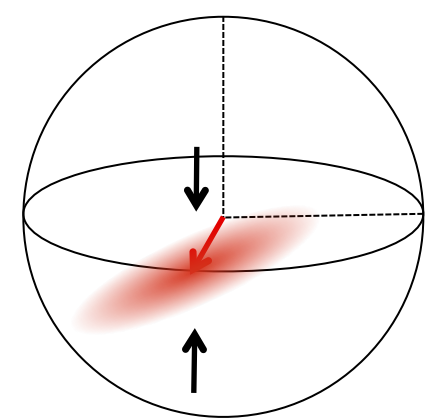
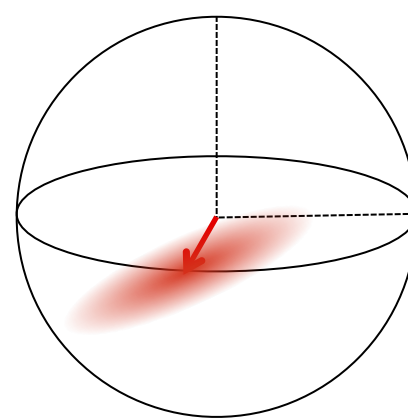
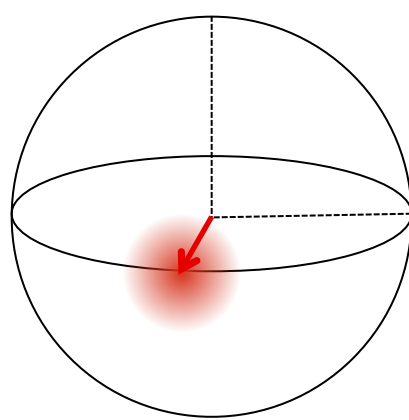
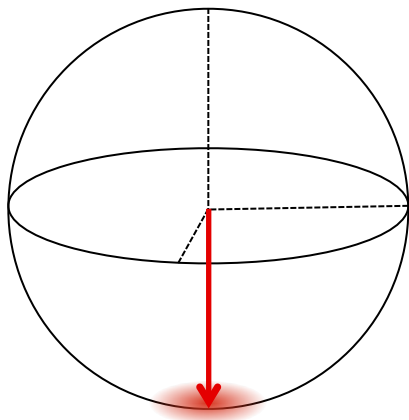
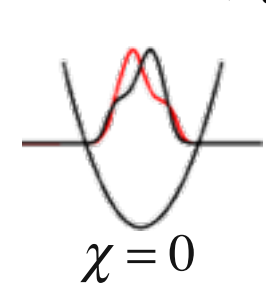
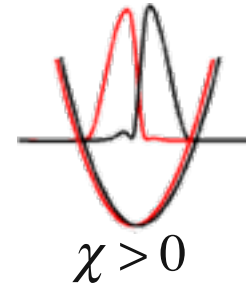
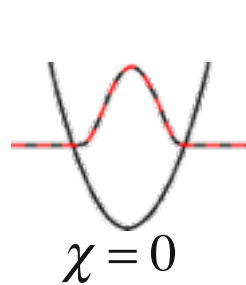
- Ramsey contrast  $C = 0.88 \pm 0.03$

**without splitting** ( $\chi = 0$ , reference)



- Ramsey contrast  $C = 0.96 \pm 0.01$
- $|1, -1\rangle$  trap lifetime  $\sim 4$  s
- $|2, 1\rangle$  trap lifetime  $\sim 200$  ms
- superposition  $\sim 250$  ms

# Experimental sequence for spin squeezing

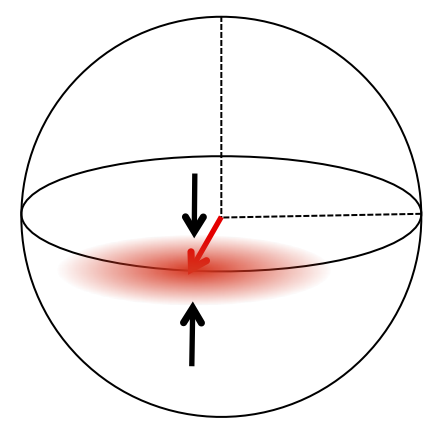
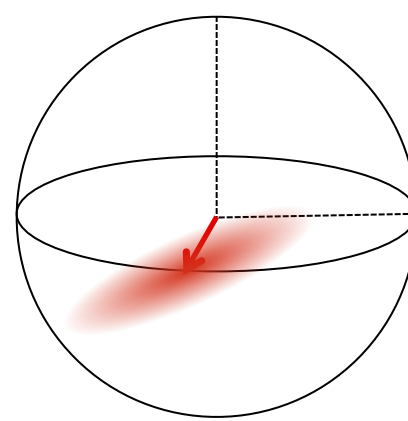
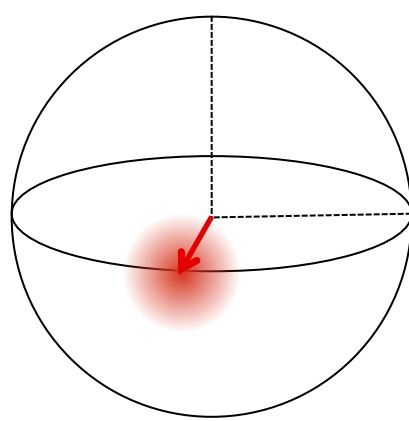
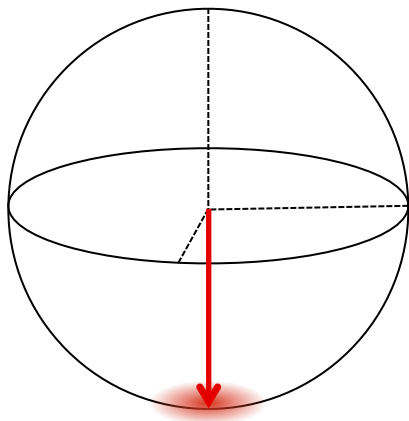
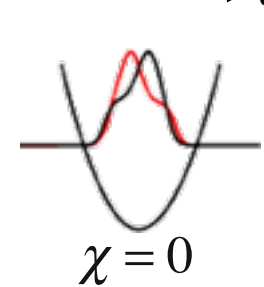
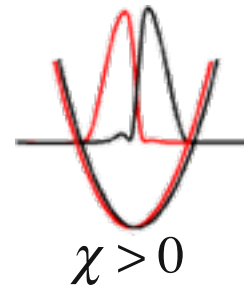
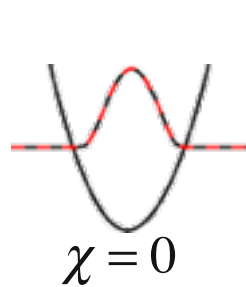


$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

**State tomography:**  
measure  $\Delta S_{\theta}$   
(projection noise)  
after turning for  
several angles  $\theta$



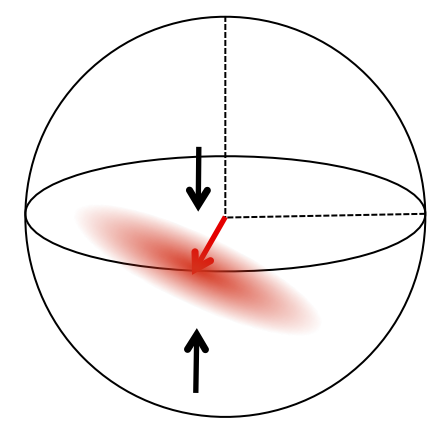
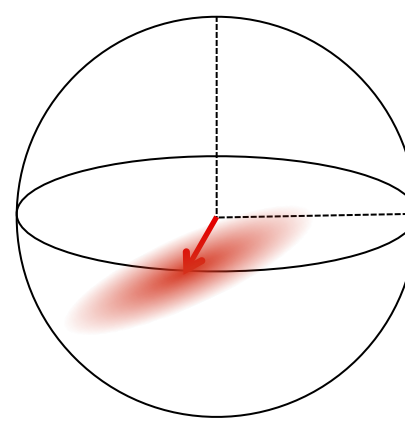
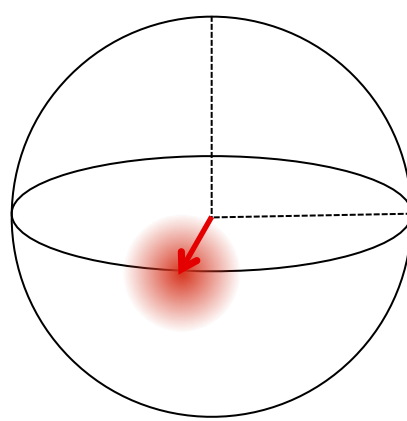
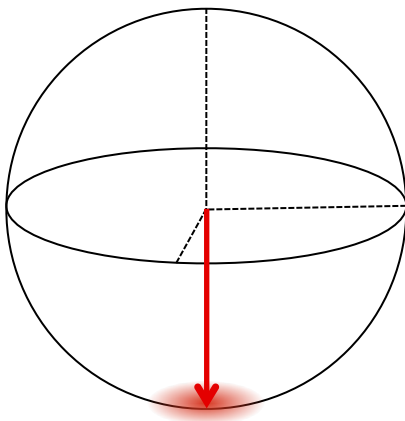
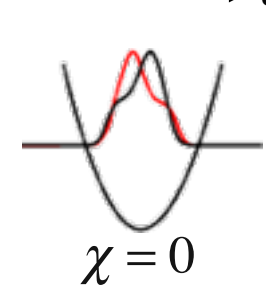
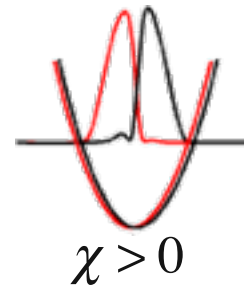
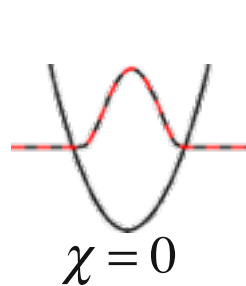
# Experimental sequence for spin squeezing



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**State tomography:**  
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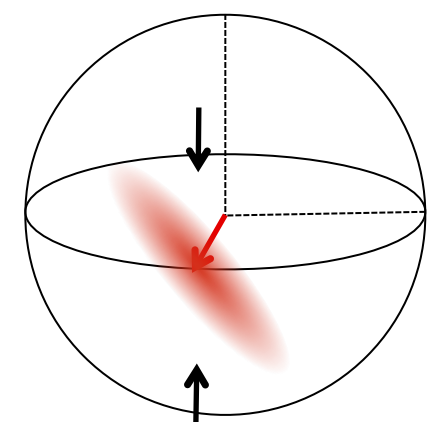
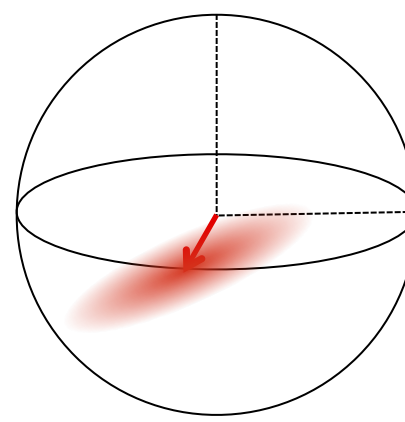
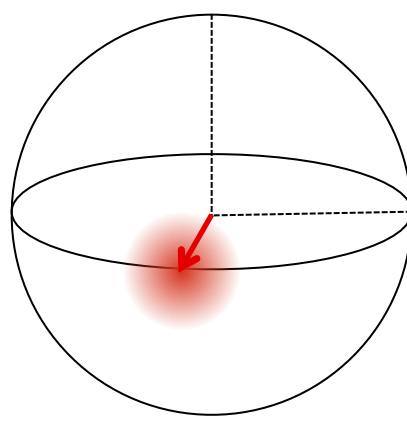
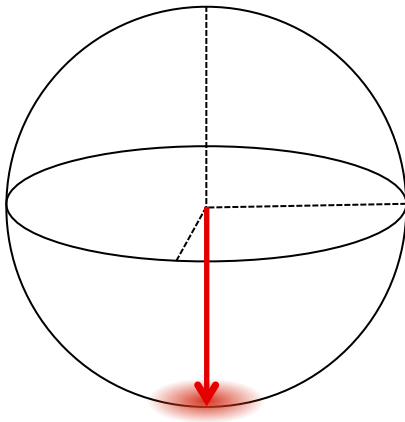
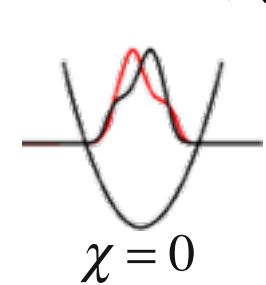
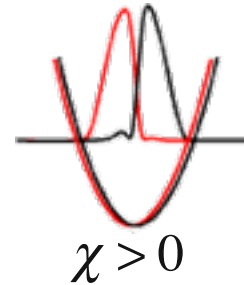
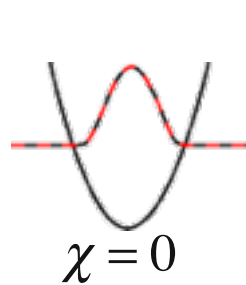
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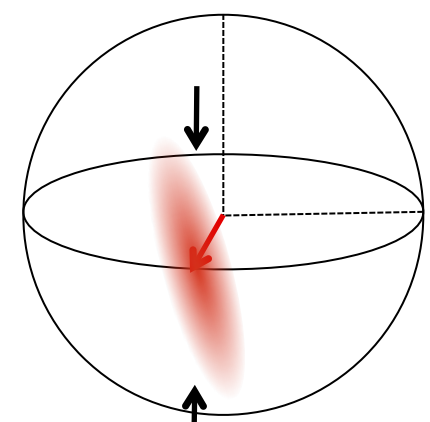
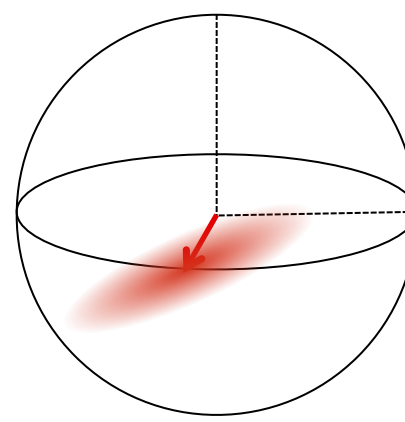
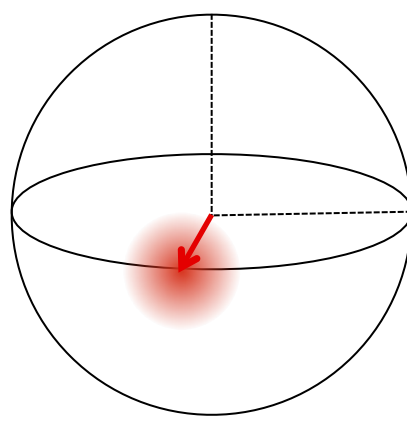
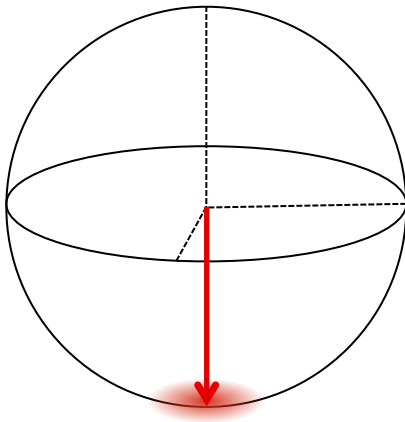
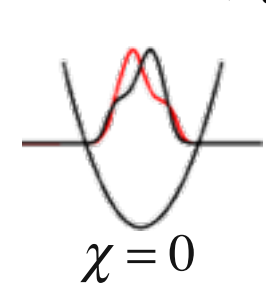
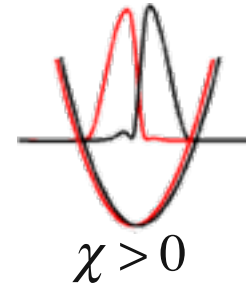
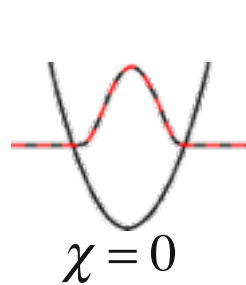
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**State tomography:**  
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several angles  $\theta$

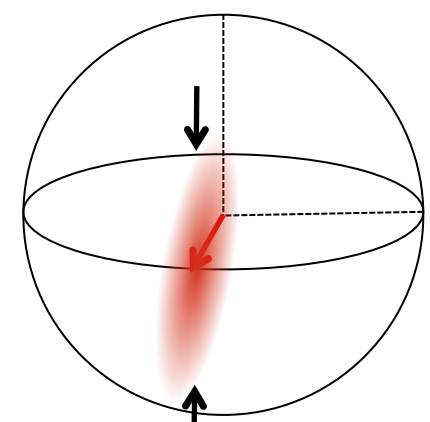
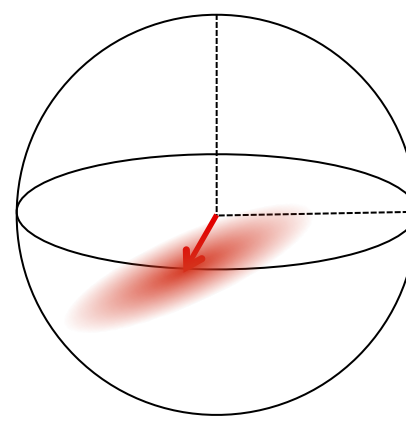
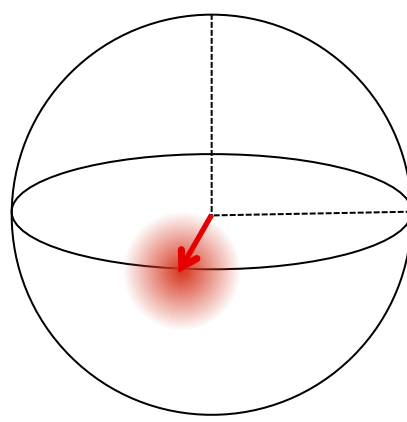
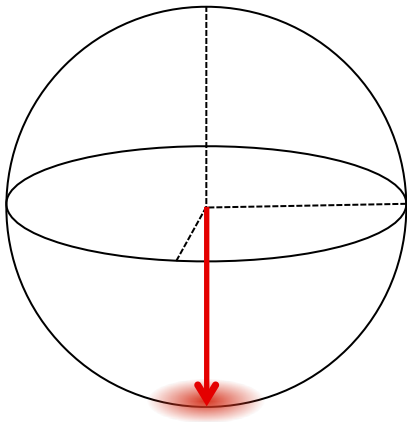
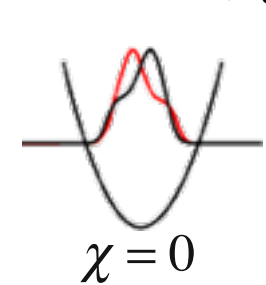
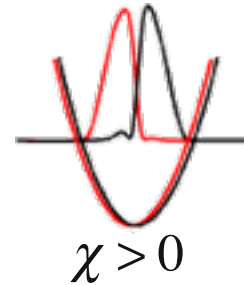
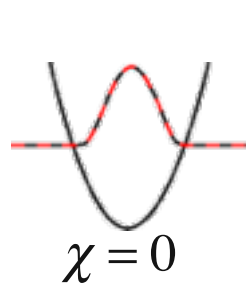
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**State tomography:**  
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# Experimental sequence for spin squeezing

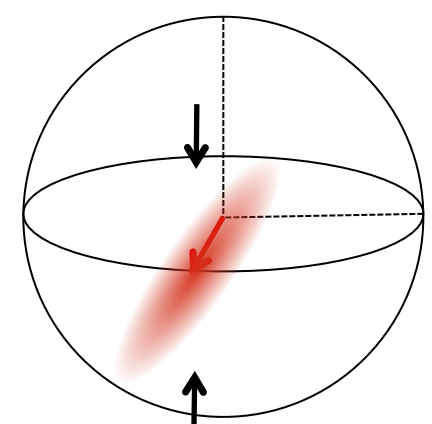
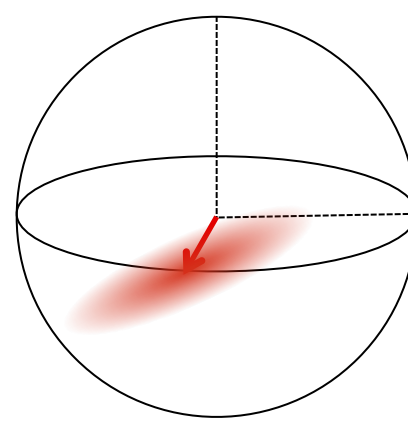
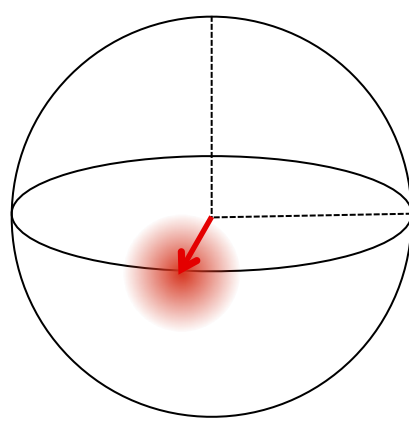
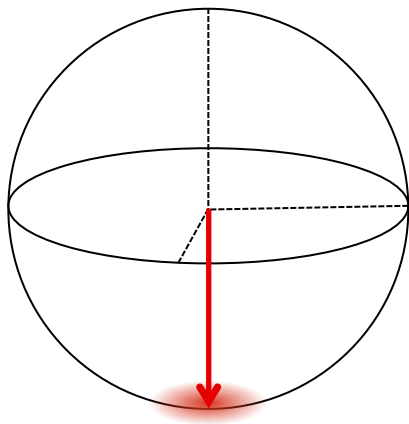
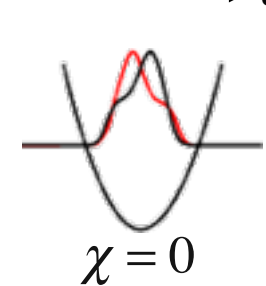
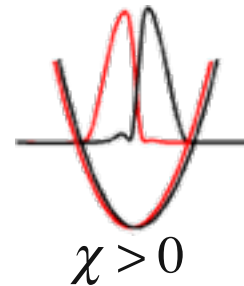
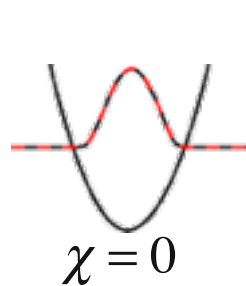


$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

**State tomography:**  
measure  $\Delta S_{\theta}$   
(projection noise)  
after turning for  
several angles  $\theta$



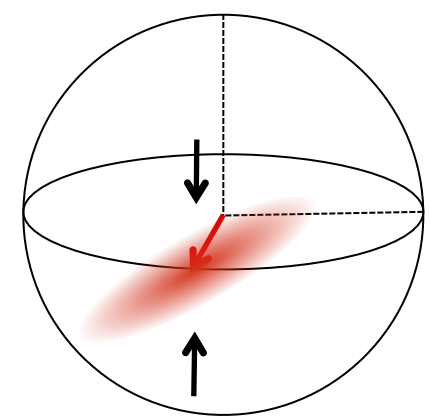
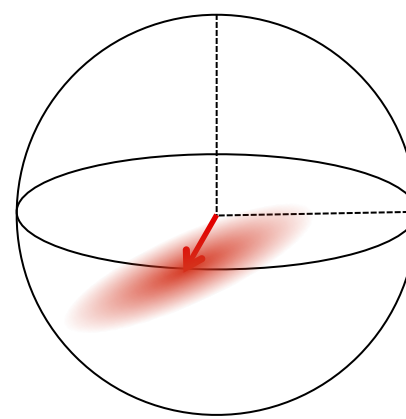
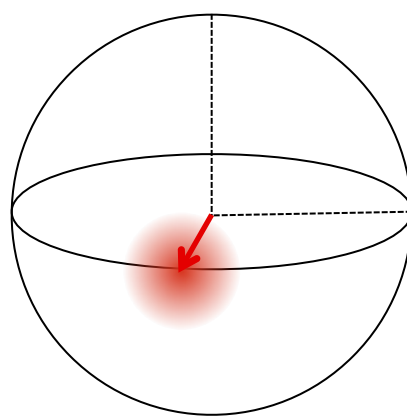
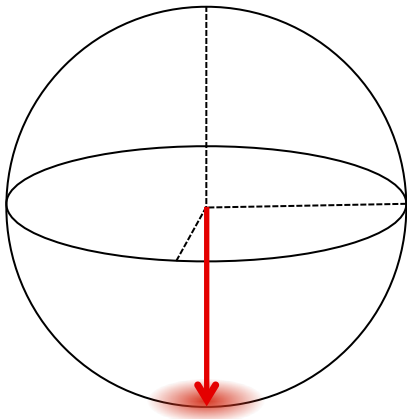
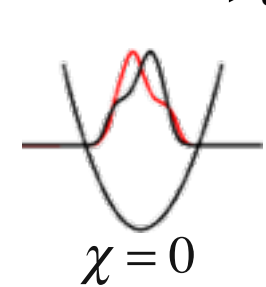
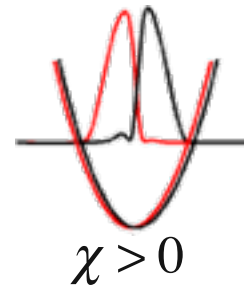
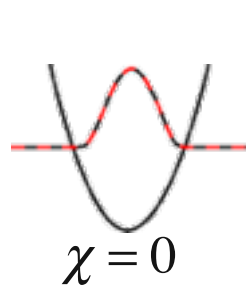
# Experimental sequence for spin squeezing



$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

**State tomography:**  
measure  $\Delta S_{\theta}$   
(projection noise)  
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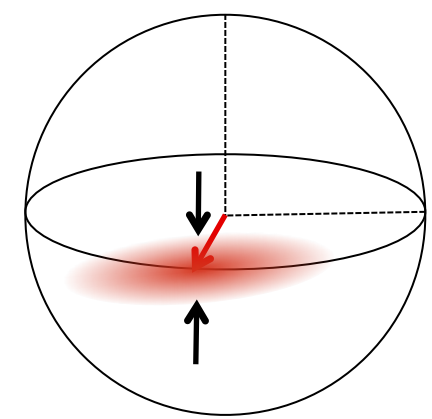
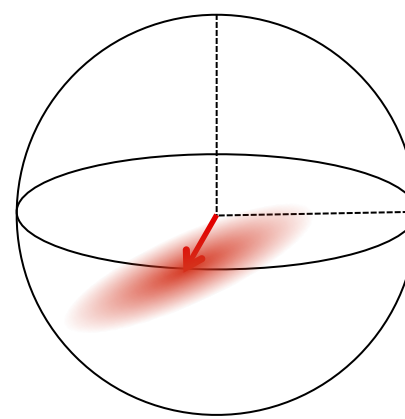
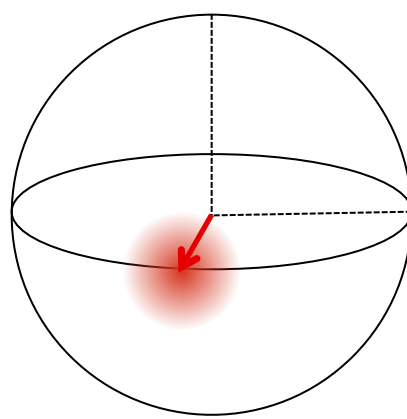
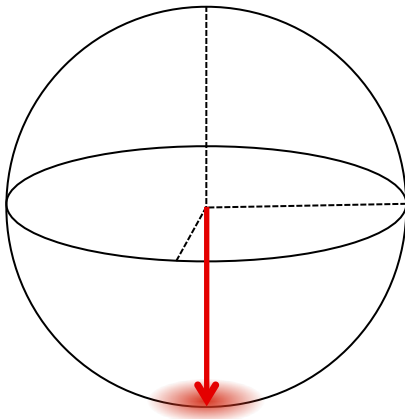
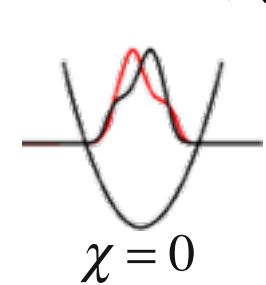
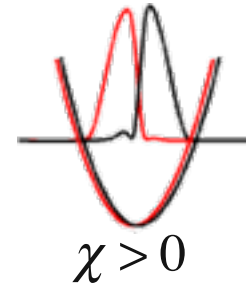
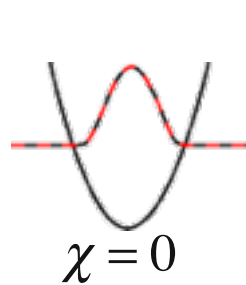
# Experimental sequence for spin squeezing



$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

**State tomography:**  
measure  $\Delta S_{\theta}$   
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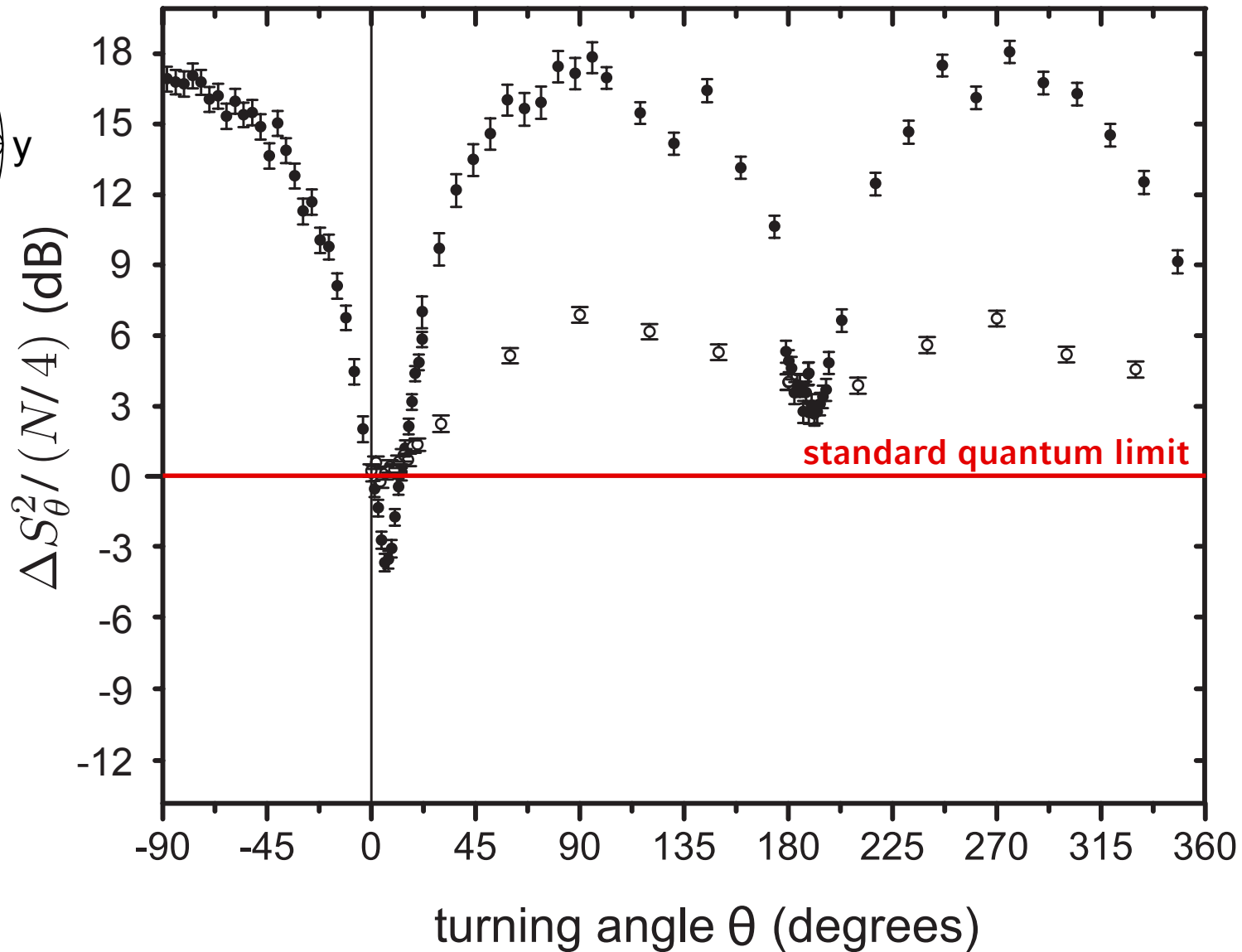
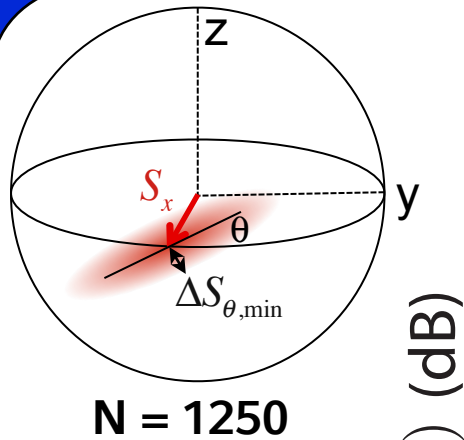
# Experimental sequence for spin squeezing



$$\xi^2 \equiv \frac{2S (\Delta S_{\theta, \min})^2}{\langle S_x \rangle^2}$$

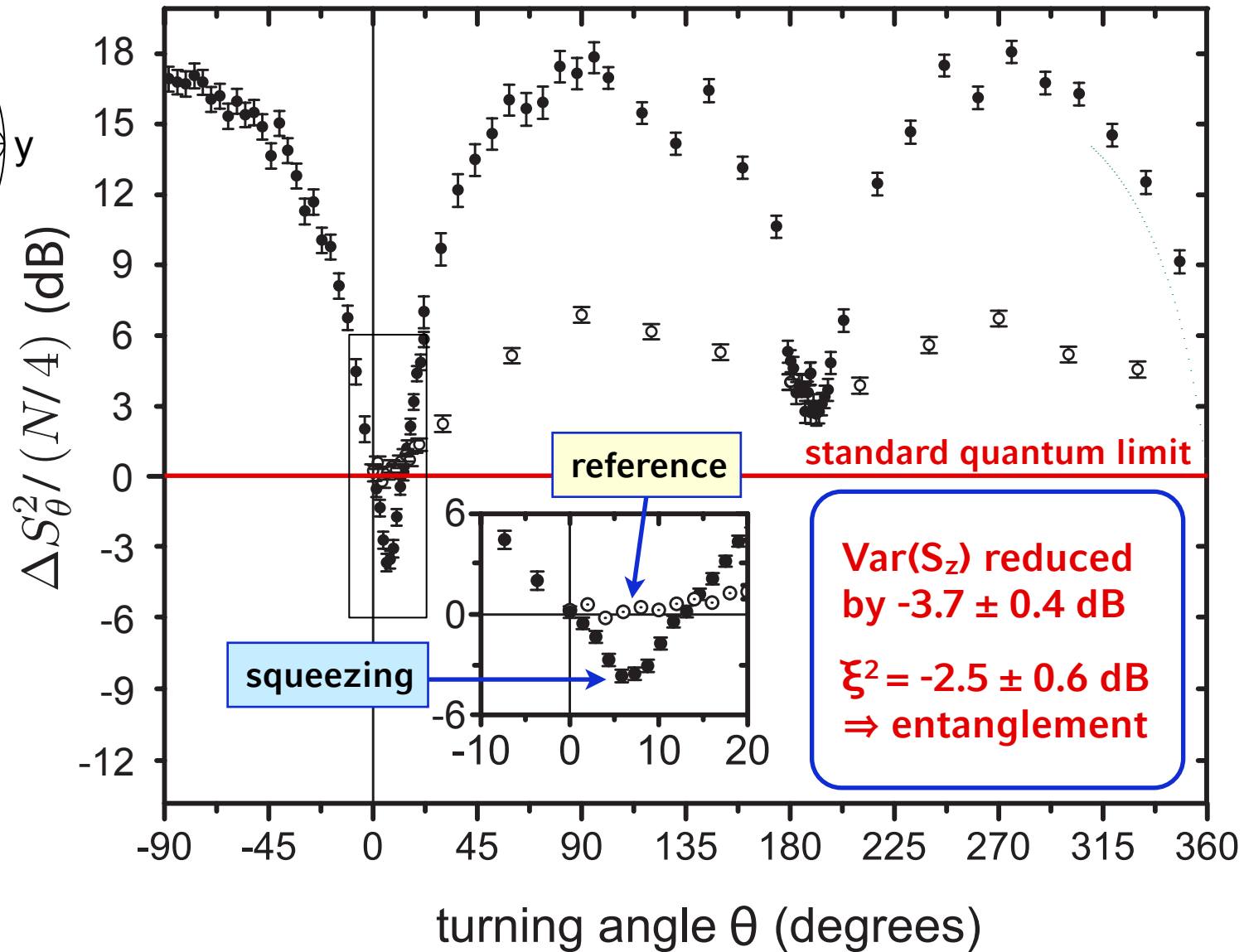
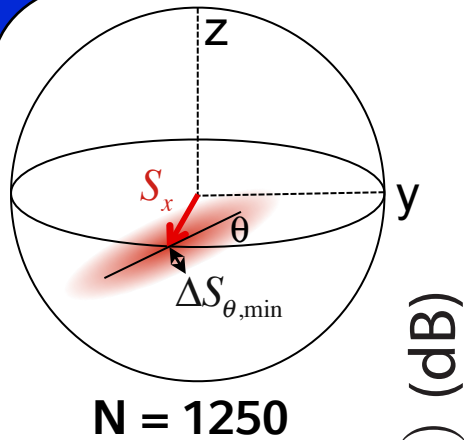
**State tomography:**  
measure  $\Delta S_\theta$   
(projection noise)  
after turning for  
several angles  $\theta$

# Spin squeezing: data



statistics:  
370 shots per  
datapoint

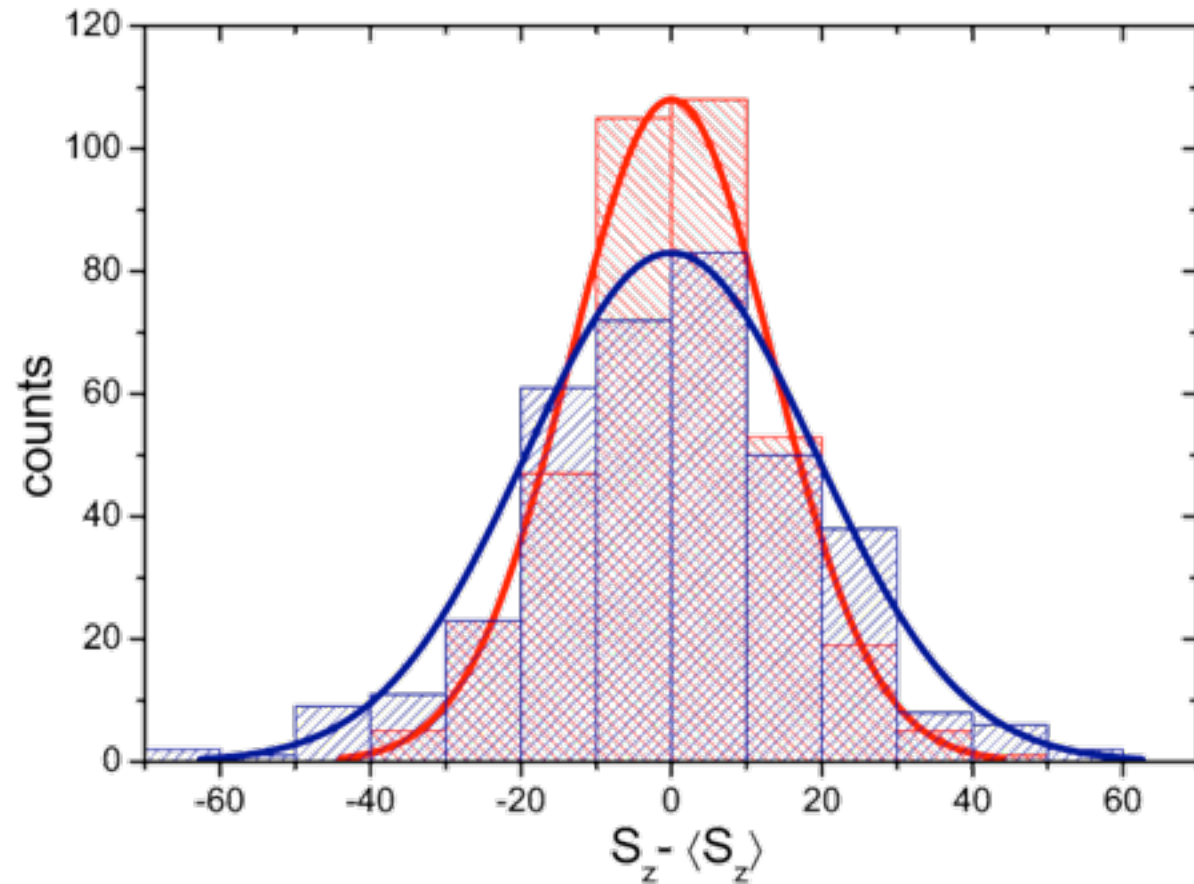
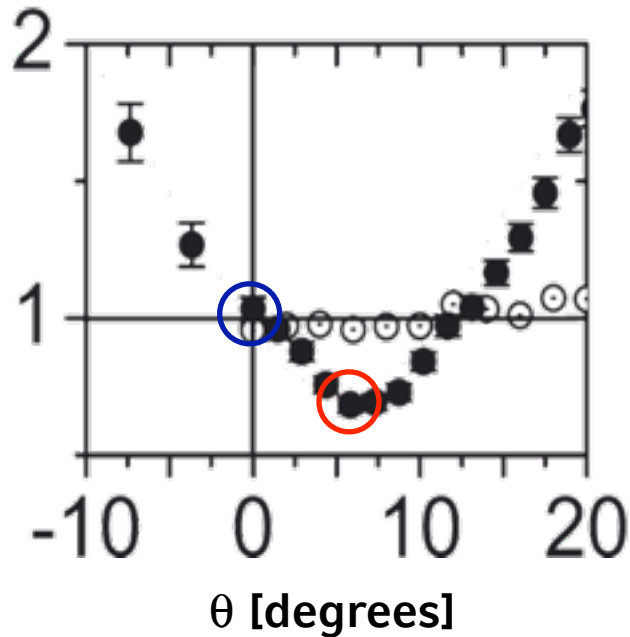
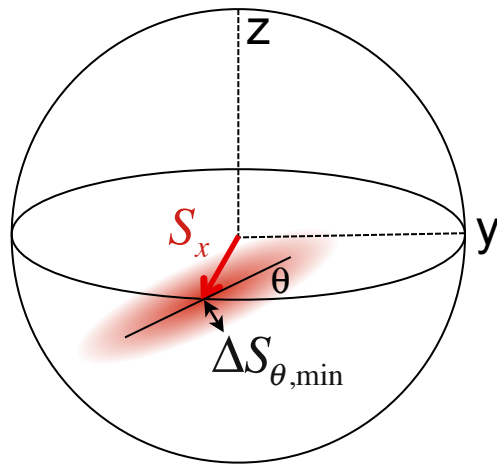
# Spin squeezing: data



statistics:  
370 shots per  
datapoint



# Marginals of the Wigner function

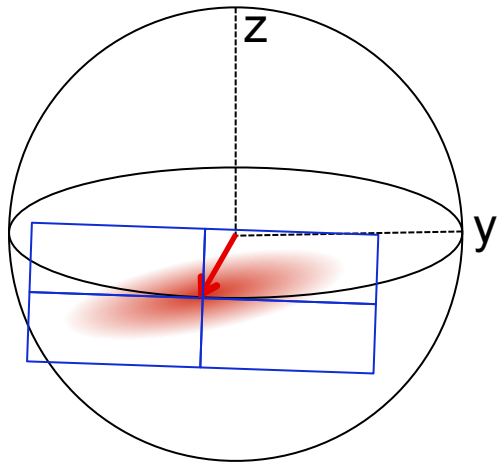
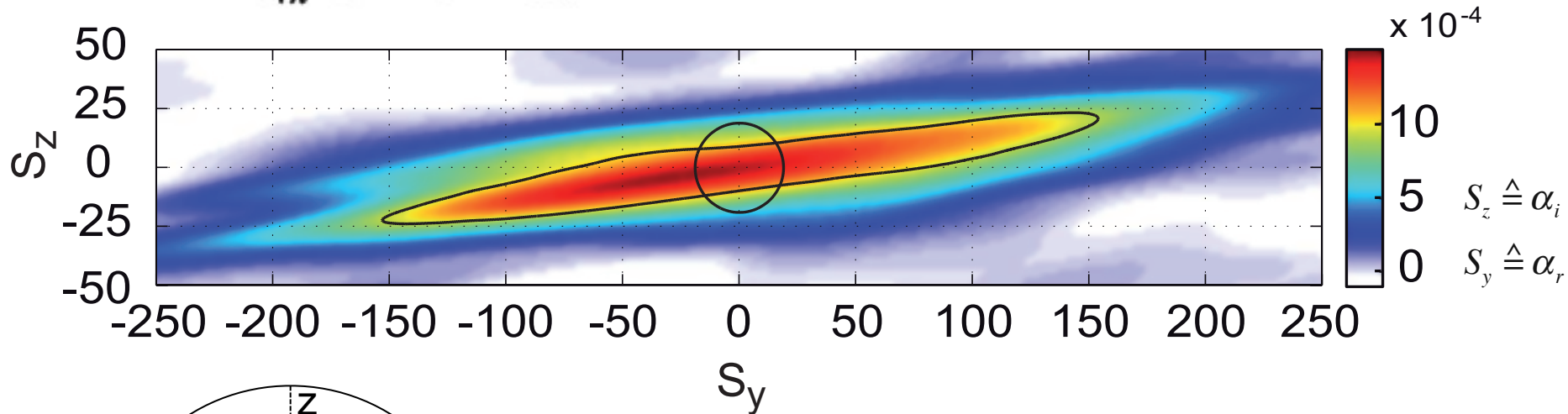


Quantum state tomography

# Quantum state reconstruction

**Wigner function** - obtained by inverse Radon transform

$$W(\alpha_r, \alpha_i, s) = \frac{1}{4\pi^2} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \int_0^\pi w(x, \theta) \exp[s\eta^2/8 + i\eta(x - \alpha_r \cos\theta - \alpha_i \sin\theta)] |\eta| dx d\eta d\theta.$$



local approximation of Bloch sphere by plane

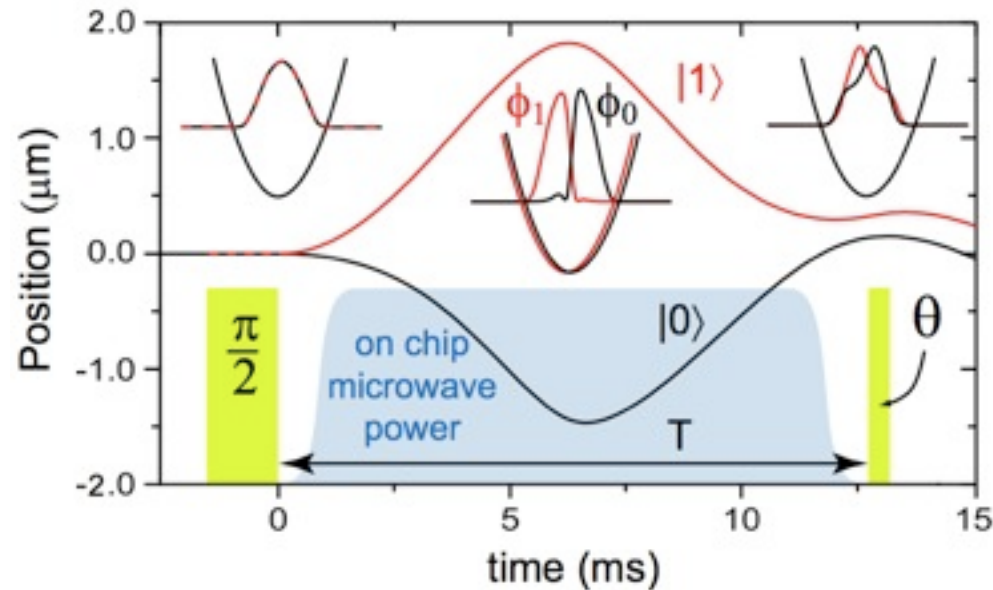
limited resolution due to

- finite angular resolution
- finite resolution in atom number (imaging noise)
- limited amount of data for histograms

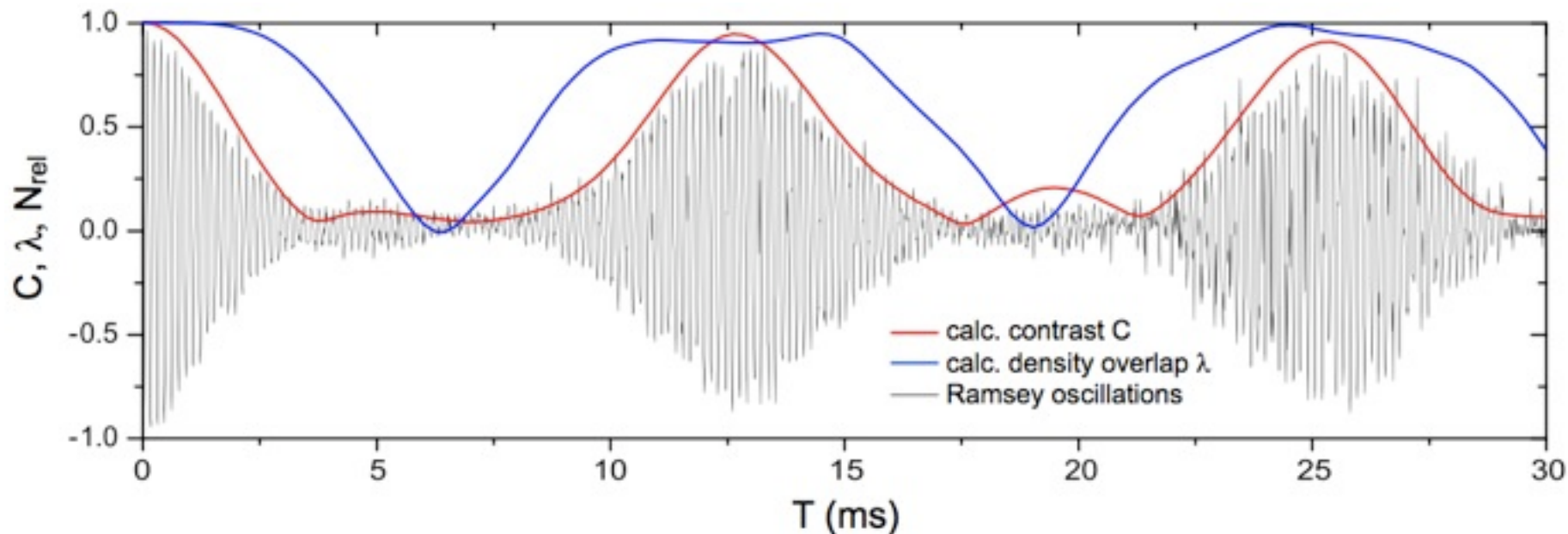
# Simulation: Li Yun and A. Sinatra, ENS Paris

Yun Li et al.,  
EPJB 68, 365 (2009).

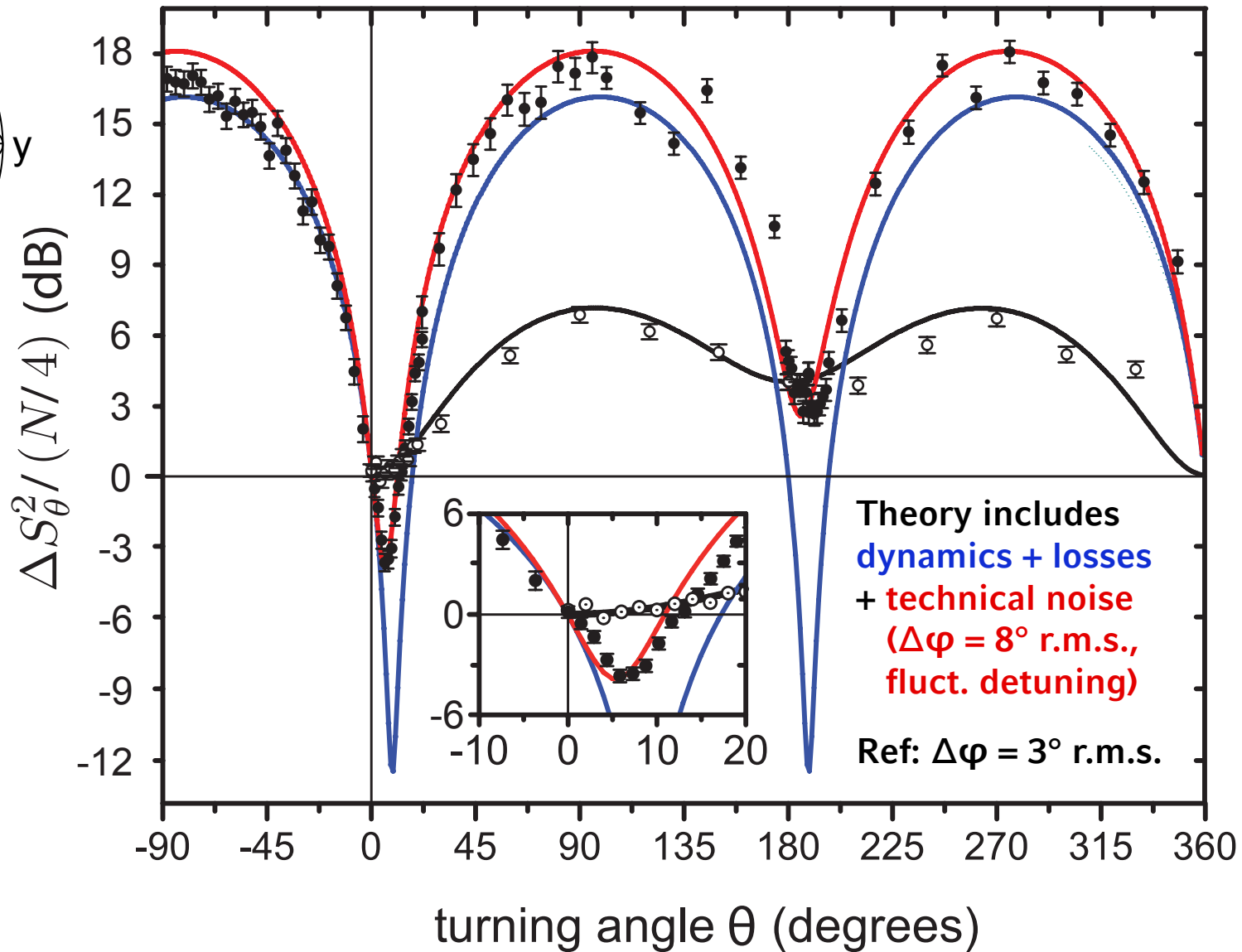
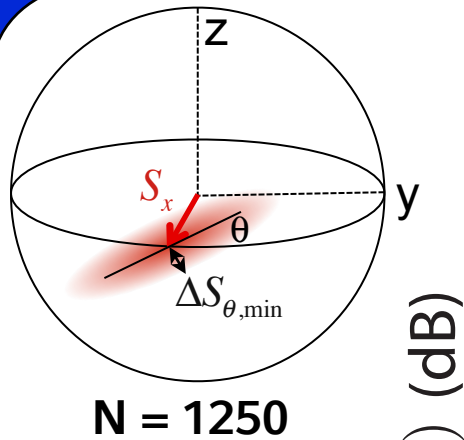
**Spatial evolution  
of the two states**  
( $\phi_0$ ,  $\phi_1$  for  $N_0=N_1=N/2$ )



**Ramsey contrast  
(theory vs. experiment)**



# Spin squeezing: data + theory



# Depth of entanglement in the squeezed BEC

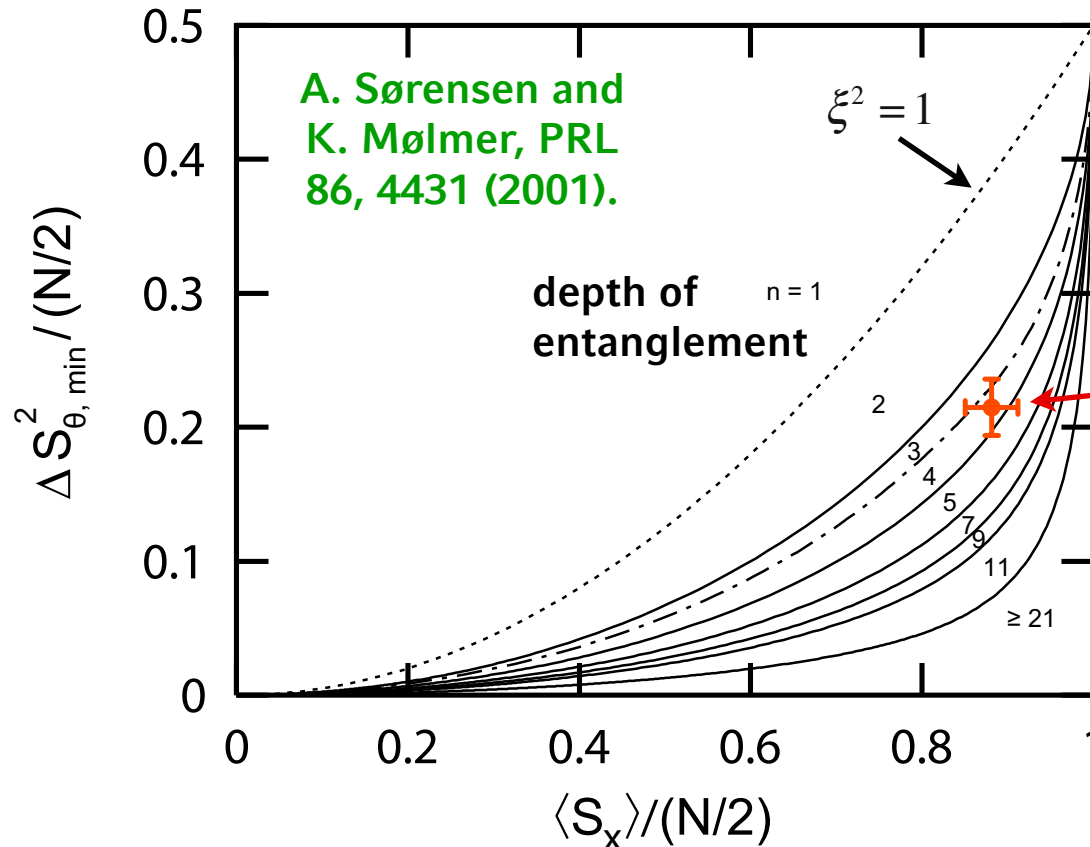
Separable state  
of N atoms:

$$\rho = \sum_k P_k \rho_k^{(1)} \otimes \rho_k^{(2)} \otimes \dots \otimes \rho_k^{(N)} \rightarrow \xi^2 \geq 1$$

A. Sørensen et al.,  
Nature 409, 63 (2001).

How large are the clusters of  
entangled atoms in the BEC at least?

$\rightarrow$  depth of entanglement



our data

$\Rightarrow$  clusters of  $\geq 4$   
entangled particles



- improve squeezing (decrease technical noise)
- study scaling with atom number and temperature
- use squeezed states in atomic clock on a chip  
(relax magnetic trap after squeezing to turn off nonlinearity: squeezing survives for  $\sim 0.6$  s in presence of loss and residual phase diffusion)
- characterize multi-particle entanglement (quantum Fisher information...)
- entanglement of several BECs through collisions
- QIP with single atoms on atom chips

P. Rosenbusch/J. Reichel,  
Observatoire, Paris

A. Smerzi et al., Trento

T. Calarco et al.,  
PRA 61, 022304 (2002).

P. Treutlein et al.,  
PRA 94, 022312 (2006).

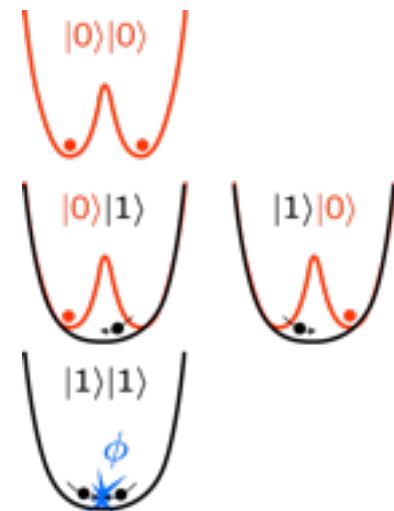
P. Treutlein et al.,  
Fortschr. Phys. 54, 702 (2006).

$$|0\rangle|0\rangle \Rightarrow |0\rangle|0\rangle$$

$$|0\rangle|1\rangle \Rightarrow |0\rangle|1\rangle$$

$$|1\rangle|0\rangle \Rightarrow |1\rangle|0\rangle$$

$$|1\rangle|1\rangle \Rightarrow e^{i\phi}|1\rangle|1\rangle$$



# Munich atom chip team - P. Treutlein / T. W. Hänsch



Collaborations: L. Yun, A. Sinatra, and J. Reichel (ENS Paris), M. Lukin (Harvard), D. König and J. Kotthaus (LMU Munich), M. Ludwig and F. Marquardt (LMU), K. Hammerer, K. Stannigel, C. Genes, M. Wallquist, and P. Zoller (Innsbruck)