PWSCF

and

diagonalization





ELECTRONS

```
call electron_scf
do iter = 1, niter
call c_bands --> C_BANDS
call sum_band --> SUM_BAND
call mix_rho
call v_of_rho
end do iter
```



PWSCF

call read_input_file (input.f90)

call run_pwscf



SETUP

defines grid and other dimensions, no system specific calculations yet

INIT_RUN

- call pre_init
- call allocate_fft
- call ggen
- call allocate_nlpot
- call allocate_paw_integrals
- call paw_one_center
- call allocate_locpot
- call allocate wfc
- call openfile
- call hinit0
- call potinit
- call newd
- call wfctinit



ELECTRONS

```
call electron_scf
do iter = 1, niter
call c_bands --> C_BANDS
call sum_band --> SUM_BAND
call mix_rho
call v_of_rho
end do iter
```



C BANDS

do ik = 1, nks call get_buffer (evc) call init_us_2 (vkb) call diag_bands --> DIAG_BANDS call save_buffer end do ik

DIAG BANDS

DAVIDSON (isolve=0) $hdiag = g2 + vloc_avg + Vnl_avg$ call cegterg or pcegterg

CG (isolve=1) hdiag = 1 + g2 + sqrt(1+(g2-1)**2)call rotate_wfc call ccgdiagg



Step 4 : diagonalization





Diagonalization of $H\ensuremath{\mbox{KS}}$ is a major step in the scf solution of any system.

In pw.x two methods are implemented:

<u>Davidson diagonalization</u>
 efficient in terms of number of Hpsi required
 memory intensive: requires a work space up to

 (1+3*david) * nbnd * npwx
 and diagonalization of matrices up to
 david*nbnd x david*nbnd
 where david is by default 4, but can be reduced to 2

•<u>Conjugate gradient</u>

-memory friendly: bands are dealt with one at a time.
-the need to orthogonalize to lower states makes it intrinsically sequential and not efficient for large systems.

Davidson Diagoalization

- •Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$ •Eigenpairs of the reduced Hamiltonian $\tilde{H}_{ij} = \langle \phi_i^{(n)} | H_{KS} | \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \phi_i^{(n)} | S | \phi_j^{(n)} \rangle$
- •Build the correction vectors $|\tilde{\phi}_i^{(n)}\rangle$ $|\tilde{\phi}_i^{(n)}\rangle = (H_{diag} - \varepsilon_i S_{diag})^{-1}(H_{KS} - \varepsilon_i S)|\phi_i^{(n)}\rangle$

•Build an extended reduced Hamiltonian

$$\tilde{H}_{ij} = \langle \phi_i^{(n)} / \tilde{\phi}^{(n)} | H_{KS} | \phi_j^{(n)} / \tilde{\phi}_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \phi_i^{(n)} / \tilde{\phi}_i^{(n)} | S | \phi_j^{(n)} / \tilde{\phi}_j^{(n)} \rangle$$

•Diagonalize the small 2nbnd x 2nbnd reduced Hamiltonian to get the new estimate for the eigenpairs $(\tilde{H} - \varepsilon \tilde{S})v = 0 \longrightarrow \{ |\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)} \}$

•Repeat if needed in order to improve the solution \rightarrow 3nbnd x 3nbnd \rightarrow 4nbnd x 4nbnd ... \rightarrow <u>nbnd x nbnd</u> <u>Davidson diagonalization</u>
efficient in terms of number of Hpsi required
memory intensive: requires a work space up to (1+3*david) * nbnd * npwx
and diagonalization of matrices up to david*nbnd x david*nbnd
where david is by default 4, but can be reduced to 2

•routines

- regterg , cegterg <u>real/cmplx</u> <u>eigen</u> i<u>ter</u>ative <u>g</u>eneralized
- h_psi, s_psi, g_psi
- rdiaghg, cdiaghg real/cmplx diagonalization <u>H</u> generalized

Conjugate Gradient

•For each band, given a trial eigenpair:

 $\{|\phi_i^{(n)}\rangle,\varepsilon_i\}$

•Minimize the single particle energy

 $E(|\phi_i\rangle) = \langle \phi_i | H_{KS} | \phi_j \rangle$

by (pre-conditioned) CG method

subject to the constraints

$$\langle \phi_i | S | \phi_j \rangle = \delta_{ij}, \quad \forall j \le i$$

.... see attached documents for more details

•Repeat for next band until completed

•<u>Conjugate gradient</u>

-memory friendly: bands are dealt with one at a time.

-the need to orthogonalize to lower states makes it intrinsically sequential and not efficient for large systems.

•<u>routines</u>

- rcgdiagg , ccgdiagg <u>r</u>eal/<u>c</u>mplx <u>CG</u> <u>diag</u>onalization <u>g</u>eneralize
- h_1psi, s_1psi
 - * preconditioning

Parallel Orbital update method

and some thoughts about

-bgrp parallelization-ortho parallelization-task parallelization

in pw.x

Some recent work on an alternative iterative methods

A PARALLEL ORBITAL-UPDATING APPROACH FOR ELECTRONIC STRUCTURE CALCULATIONS *

XIAOYING DAI[†], XINGAO GONG[‡], AIHUI ZHOU[†], AND JINWEI ZHU[†]

Abstract. In this paper, we propose an orbital iteration based parallel approach for electronic structure calculations. This approach is based on our understanding of the single-particle equations of independent particles that move in an effective potential. With this new approach, the solution of the single-particle equation is reduced to some solutions of independent linear algebraic systems and a small scale algebraic problem. It is demonstrated by our numerical experiments that this new approach is quite efficient for full-potential calculations for a class of molecular systems.

arXiv:1405.0260v2 [math.NA] 20/11/2014

A PARALLEL ORBITAL-UPDATING BASED OPTIMIZATION METHOD FOR ELECTRONIC STRUCTURE CALCULATIONS *

XIAOYING DAI[†], ZHUANG LIU[‡], XIN ZHANG[§], AND AIHUI ZHOU[¶]

Abstract. In this paper, we propose a parallel optimization method for electronic structure calculations based on a single orbital-updating approximation. It is shown by our numerical experiments that the method is efficient and reliable for atomic and molecular systems of large scale over supercomputers.

arXiv:1510.07230v1 [math.NA] 25/10/2015

ParO in a nutshell

Algorithm 1.1.

- 1. Given initial data $(\lambda_i^{(0)}, u_i^{(0)}) \in \mathbb{R} \times H_0^1(\Omega)$ with $(u_i^{(0)}, u_j^{(0)})_{\Omega} = \delta_{ij}, (i, j = 1, 2, \dots, N)$, define \mathcal{T}_0 and V_0 , and let n = 0
- 2. Construct \mathcal{T}_{n+1} and V_{n+1} based on an adaptive procedure to $(\lambda_i^{(n)}, u_i^{(n)})$. 3. For $i = 1, 2, \dots, N$, find $u_i^{(n+1/2)} \in V_{n+1}$ satisfying

$$a(U^{(n)}; u_i^{(n+1/2)}, v) = \lambda_i^{(n)}(u_i^{(n)}, v) \ \forall v \in V_{n+1}$$

in parallel.

4. Project to eigenspace: find $(\lambda^{(n+1)}, u^{(n+1)}) \in \mathbb{R} \times \tilde{V}_{n+1}$ satisfying $||u^{(n+1)}||_{0,\Omega} = 1$ and

$$a(U^{(n+1/2)}; u^{(n+1)}, v) = \lambda^{(n+1)}(u^{(n+1)}, v) \quad \forall v \in \tilde{V}_{n+1}$$

to obtain eigenpairs $(\lambda_i^{(n+1)}, u_i^{(n+1)})(i = 1, 2, \dots, N)$. 5. Let n = n + 1 and go to Step 2. Here $\tilde{V}_{n+1} = span \{u_1^{(n+1/2)}, u_2^{(n+1/2)}, \dots, u_N^{(n+1/2)}\}, U^{(n)} = (u_1^{(n)}, u_2^{(n)}, \dots, u_N^{(n)}), U^{(n+1/2)} = (u_1^{(n+1/2)}, u_2^{(n+1/2)}, \dots, u_N^{(n+1/2)}), and a(\cdot; \cdot, \cdot) is the nonlinear variational form associated the Kohn-Sham equation defined in Section 2.2.$

arXiv:1405.0260v2 [math.NA] 20/11/2014

ParO as I understand it

•Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

•Solve in parallel the *nbnd* linear systems

$$(H_{KS} + \lambda S) |\tilde{\phi}_i^{(n)}\rangle = (\varepsilon_i^{(n)} + \lambda) S |\phi_i^{(n)}\rangle$$

•Build the reduced Hamiltonian

$$\tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} | S | \tilde{\phi}_j^{(n)} \rangle$$

•Diagonalize the <u>small</u> *nbnd x nbnd* reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \longrightarrow \{ |\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)} \}$$

A variant of ParO method

•Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

•Solve in parallel the *nbnd* linear systems

$$(H_{KS} + \lambda S) |\tilde{\phi}_i^{(n)}\rangle = (\varepsilon_i^{(n)} + \lambda) S |\phi_i^{(n)}\rangle$$

•Build the reduced Hamiltonian from both $|\tilde{\phi}_i^{(n)}\rangle \& |\phi_i^{(n)}\rangle$ $\tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | S | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle$

•Diagonalize the small $2nbnd \times 2nbnd$ reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \longrightarrow \{ |\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)} \}$$

A variant of ParO method (2)

•Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

•Solve in parallel the *nbnd* linear systems

 $\begin{pmatrix} H_{KS} - \varepsilon_i^{(n)}S + \alpha S | \phi_i^{(n)} \rangle \langle \phi_i^{(n)} | S \rangle | \tilde{\phi}_i^{(n)} \rangle = -(H_{KS} - \varepsilon_i^{(n)}S) | \phi_i^{(n)} \rangle \\ \bullet \text{Build the reduced Hamiltonian from both } | \tilde{\phi}_i^{(n)} \rangle & \| \phi_i^{(n)} \rangle \\ \tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | S | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle$

•Diagonalize the small $2nbnd \times 2nbnd$ reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \longrightarrow \{ |\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)} \}$$

A variant of ParO method (3)

•Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

•Solve in parallel the *nbnd* linear systems

$$\begin{split} \left(H_{KS} - \varepsilon_{i}^{(n)}S + \alpha S |\phi_{i}^{(n)}\rangle \langle \phi_{i}^{(n)}|S \right) |\tilde{\phi}_{i}^{(n)}\rangle &= -(H_{KS} - \varepsilon_{i}^{(n)}S) |\phi_{i}^{(n)}\rangle \\ \bullet \text{Build the reduced Hamiltonian from } \tilde{|}\phi_{i}^{(n)}\rangle &= |\phi_{i}^{(n)}\rangle + |\tilde{\phi}_{i}^{(n)}\rangle \\ \tilde{H}_{ij} &= \langle \tilde{|}\phi_{i}^{(n)}|H_{KS}|\tilde{|}\phi_{j}^{(n)}\rangle, \quad \tilde{S}_{ij} &= \langle \tilde{|}\phi_{i}^{(n)}|S|\tilde{|}\phi_{j}^{(n)}\rangle \end{split}$$

•Diagonalize the <u>small</u> *nbnd x nbnd* reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \longrightarrow \{ |\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)} \}$$

Memory requirements for ParO method

•Memory required is nbnd * npwx + [nbnd*npwx] in the original ParO method or when $\tilde{\phi}_i^{(n)}$ are used.

•Memory required is 3 * nbnd * npwx + [2*nbnd*npwx] if both $|\tilde{\phi}_i^{(n)}\rangle \& |\phi_i^{(n)}\rangle$ are used.

•Could be possible to reduce this memory and/or the number of h_psi involved by playing with the algorithm.

Comparison with the other methods

•NOT competitive with Davidson at the moment

•Timing and number of h_psi calls similar to cg on a single bgrp basis. It scales !

216 Si atoms in a SC cell : Timing



216 Si atoms in a SC cell : Timing



Not only Silicon: BaTiO3 320 atms, 2560 el



Not only Silicon: BaTiO3 320 atms, 2560 el



Comparison with the other methods

•NOT competitive with Davidson at the moment

•Timing and number of h_psi calls similar to CG on a single bgrp basis. It scales well with bgrp parallelization!

TO DO LIST

- •Profiling of a few relevant test cases
- •Extend band parallelization to other parts

•Understand why h_psi is so much more efficient in the Davidson method.

•See if number of h_psi can be reduced

bgrp parallelization

•We should use bgrp parallelization more extensively distributing work w/o distributing data (we have R&G parallelization for that) so as to scale up to more processors.

We can distribute different loops in different routines (nats, nkb, ngm, nrxx, ...). Only local effects: incremental!
A careful profiling of the code is required.

ortho/diag parallelization

•It should be a sub comm of the pool comm (k-points) not of the bgrp comm.

•Does it give any gain ? Except for some memory reduction I saw no gain (w/o scalapack).

task parallelization

•Only needed for very large/anisotropic systems, intrinsically requiring many more processors than planes.

Is not a method to scale up the number of processors for a "small" calculation (should use bgrp parallelization for that).
Should be activated also when m < dffts%nogrp