

Superfluidity of a 2D Bose gas

(arXiv:1205.4536v1)

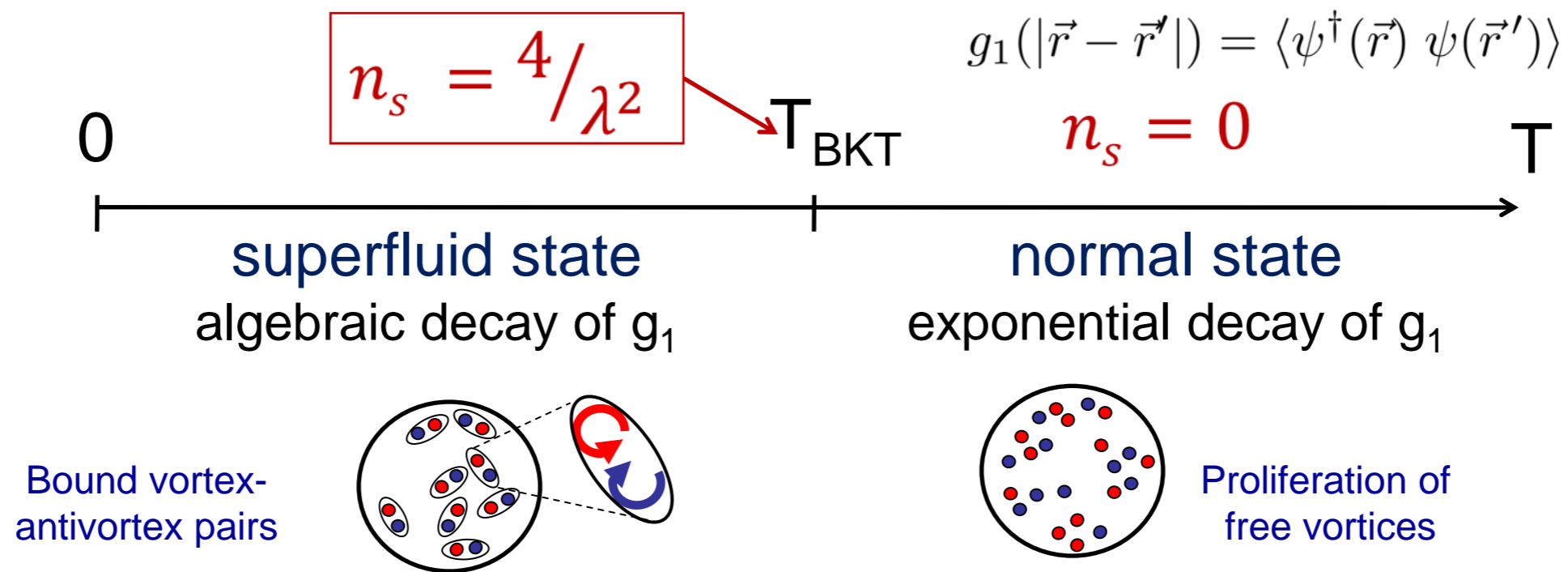
Christof Weitenberg,
Rémi Desbuquois, Lauriane Chomaz, Tarik Yefsah,
Julian Leonard, Jérôme Beugnon, Jean Dalibard

Trieste 18.07.2012



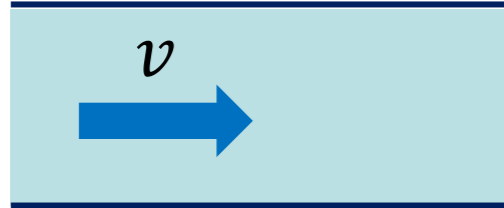
Phase transitions in 2D

- Mermin-Wagner theorem: no long ranged order in low dimensions (no breaking of continuous symmetry). In particular no BEC at finite T in uniform system.
- But Berezinskii-Kosterlitz Thouless (BKT) transition to state with quasi-long-ranged order below T_{BKT} (algebraic decay of coherence).



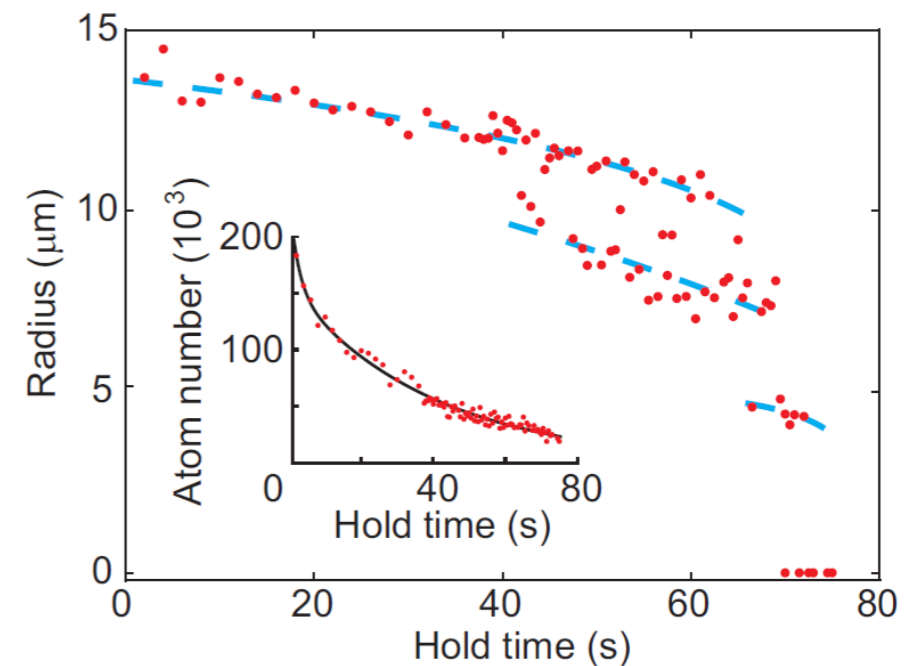
- Coherence and microscopic nature of the transition were studied with cold atoms. **But a direct proof of superfluidity is missing with these systems.**

Superfluidity (in 2D and 3D)



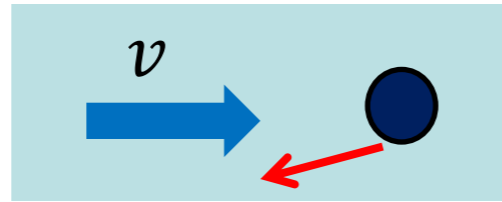
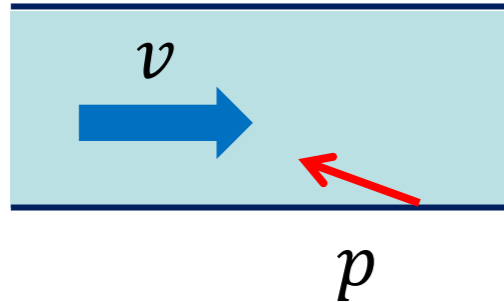
Fluid flows with velocity v with respect to container walls or impurity

- Flow without friction, persistent current.
- Metastable state: equilibrium means fluid at rest ($v = 0$).
- Height of energy barrier depends on v and T and disappears at a critical velocity v_c .
- Decay time can be very long.
- Elementary excitations: phonons, vortices, rotons...



Persistent current in 3D toroidal trap
Moulder et al., arXiv:1112.0334 (2012)

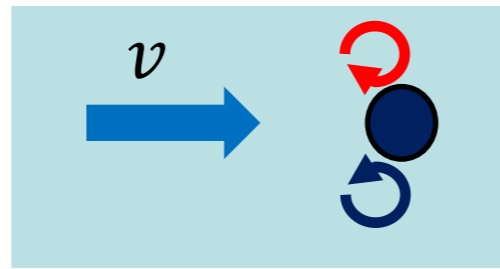
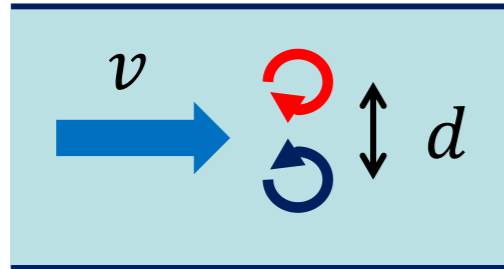
Landau criterion



An excitation with momentum p opposes the flow.

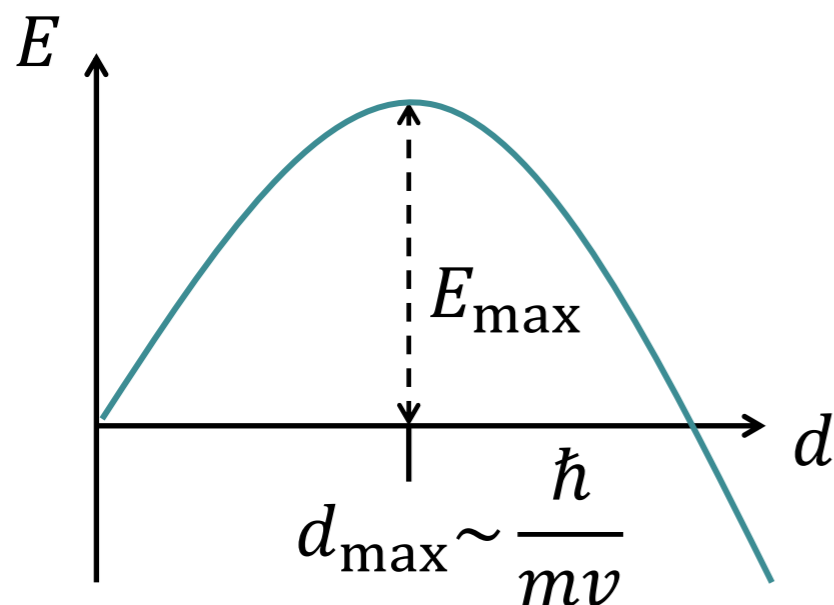
- Derived critical velocity for a dispersion relation $\epsilon(p)$ from consideration of transformation of energy and momentum.
 - Excitation energy in rest frame of container $\epsilon(p) + p \cdot v$
- Landau critical velocity $v_c = \min_p \frac{\epsilon(p)}{p}$
- For free particles $v_c = 0$
 - For Bogoliubov spectrum $v_c = c_s$ (~ 1 mm/s for cold atoms)

Stability of the superfluid flow



The flow can decay via vortex pairs (rings in 3D) that oppose the flow.

- Energy as function of vortex separation d has barrier E_{\max}
- Stability for $\gamma k_B T \lesssim E_{\max} \Rightarrow v < v_c$
- Extended obstacle with diameter w : vortex pair created with $d \sim w$. If $w > d_{\max}$ (or $v > v_c = \hbar/mw$) the flow is unstable.
- Similar argument by Feynman for flow through channel to explain the low $v_c \ll c_s$ observed for liquid helium.



In 3D (vortex ring):

$$E_{\max} = \frac{\rho_s \hbar^3}{m^2 v \hbar^3}$$

$$v_c = \rho_s \frac{\hbar^3}{\gamma m^2 k_B T}$$

In 2D (vortex pair):

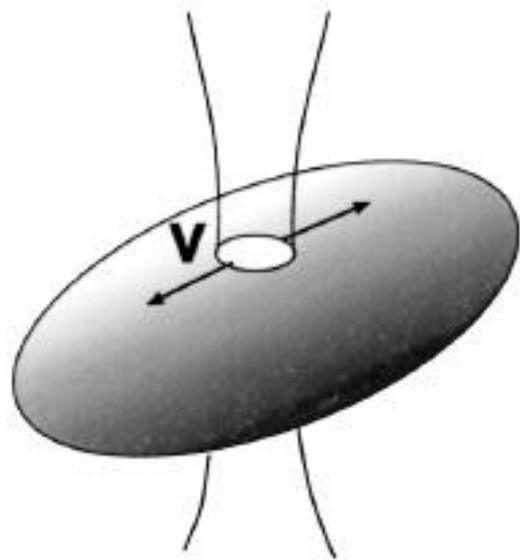
$$E_{\max} \approx \frac{\rho_s^{(2D)} \hbar^2}{m} \ln \frac{\hbar}{m \xi v}$$

$$v_c = c e^{-\gamma / PSD}$$

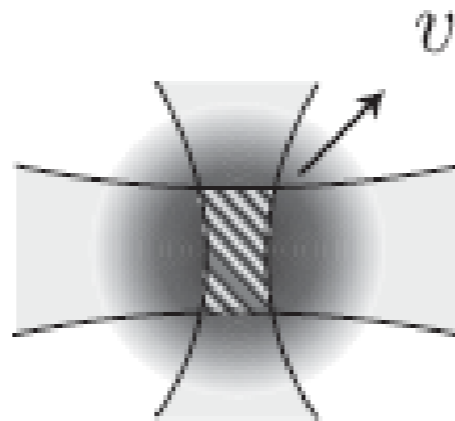
Stirring the superfluid

Experiments with cold atoms (3D systems)

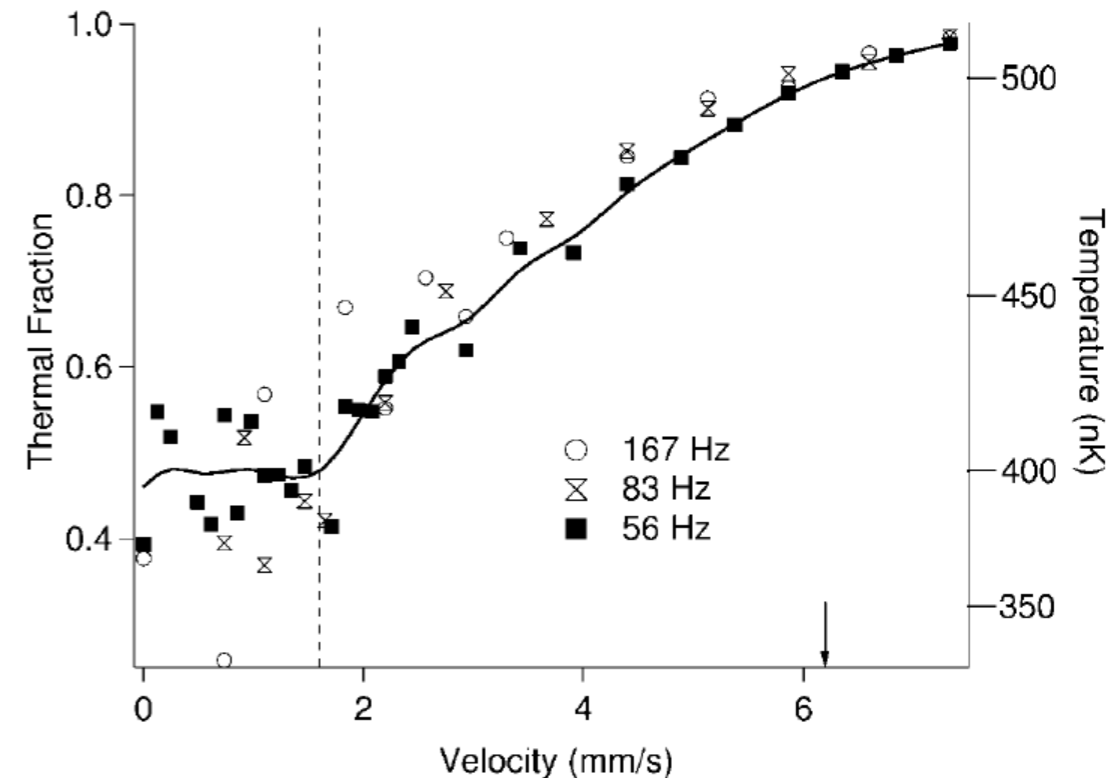
- move a small defect (laser beam or 1D lattice) through the cloud
- observe subsequent heating or excitations
- done for Bosonic and Fermionic systems



Moving laser beam



Moving optical lattice



Transpose technique to 2D system

- Allows smaller defect (no divergence of laser beam inside cloud)
- Allows measurements at fixed density by stirring in a circle

Raman et al., PRL **83**, 2502 (1999)

Raman et al., J. Low Temp. Phys. **122**, 99 (2001)

Neely et al., PRL **104**, 160401 (2010)

Onofrio et al., PRL **85**, 2228 (2000)

Engels and Atherton, PRL **99**, 160405 (2007)

Miller et al., PRL **99**, 070402 (2007)

Preparation of 2D Bose gases

- Superimpose a blue-detuned dipole potential to a degenerate 3D gas (^{87}Rb) in a magnetic TOP trap
- Depump atoms in the side wells with resonant light
- Evaporate with RF to reach degeneracy

Typical cloud :

$$T = 65 - 120 \text{ nK}$$

$$N = 3 - 9 \cdot 10^4$$

$$\mu = k_B \cdot 35 - 60 \text{ nK}$$

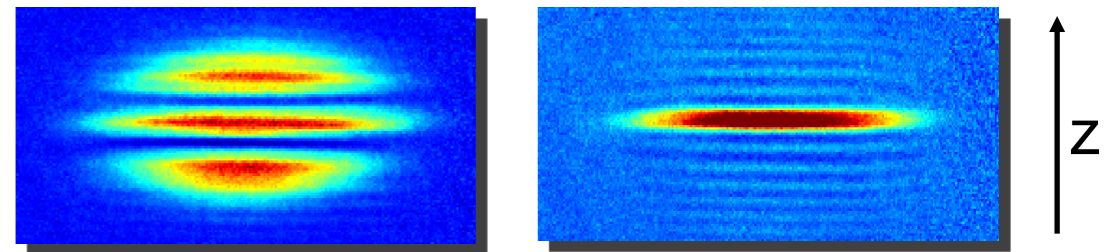
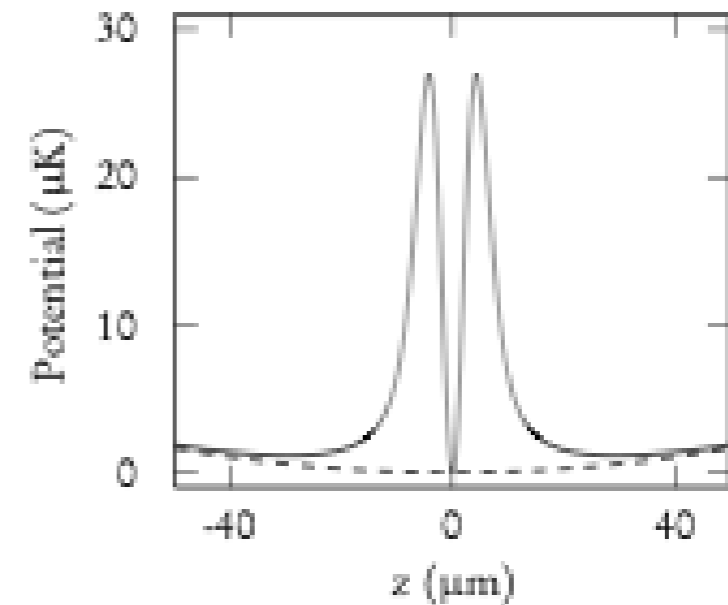
$$\omega_r = 2\pi \cdot 25 \text{ Hz}$$

$$\omega_z = 2\pi \cdot 1.4 \text{ kHz}$$

$$\tilde{g} = \sqrt{8\pi} a/l_z = 0.093$$

$$\xi \sim 0.5 \text{ } \mu\text{m}$$

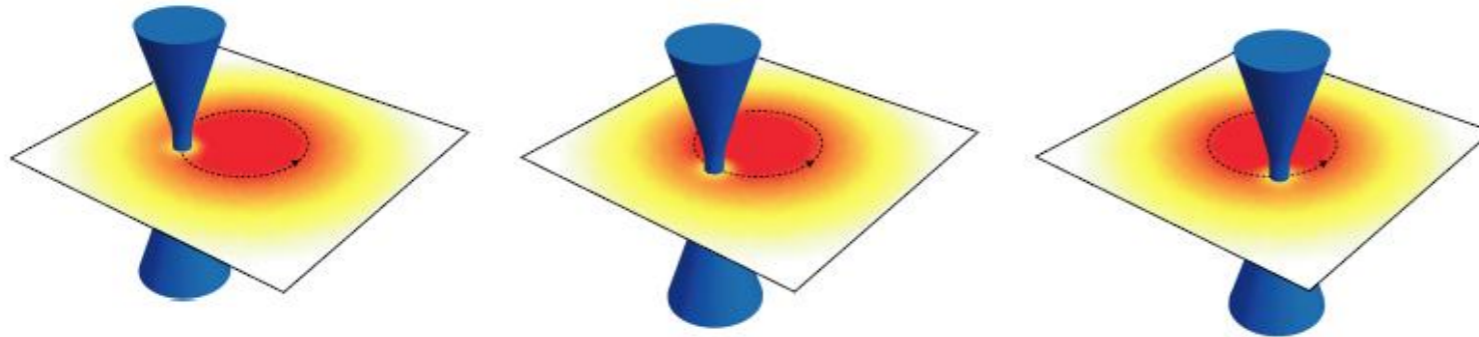
magnetic TOP trap
+
optical confinement



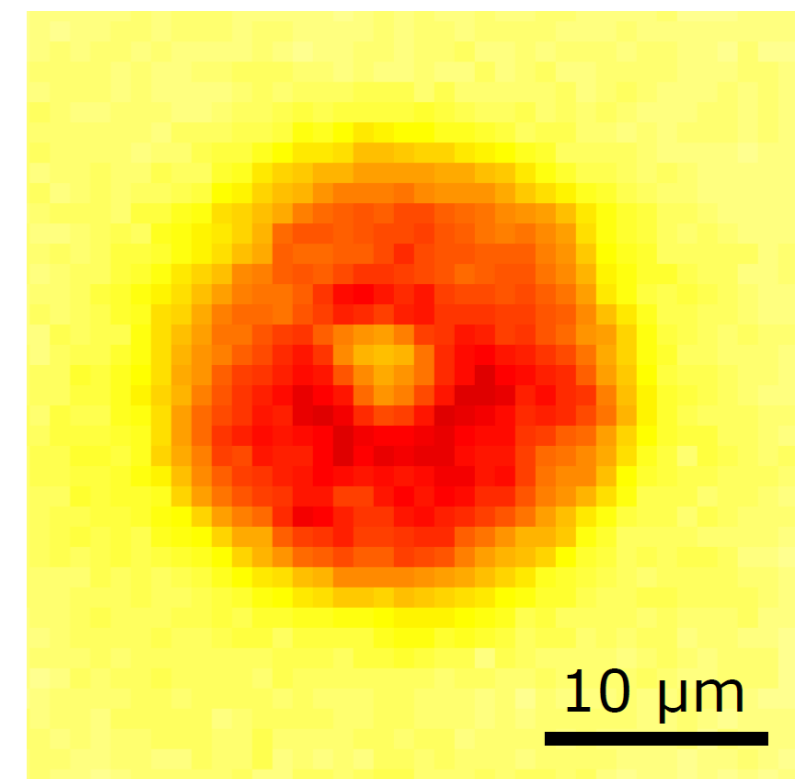
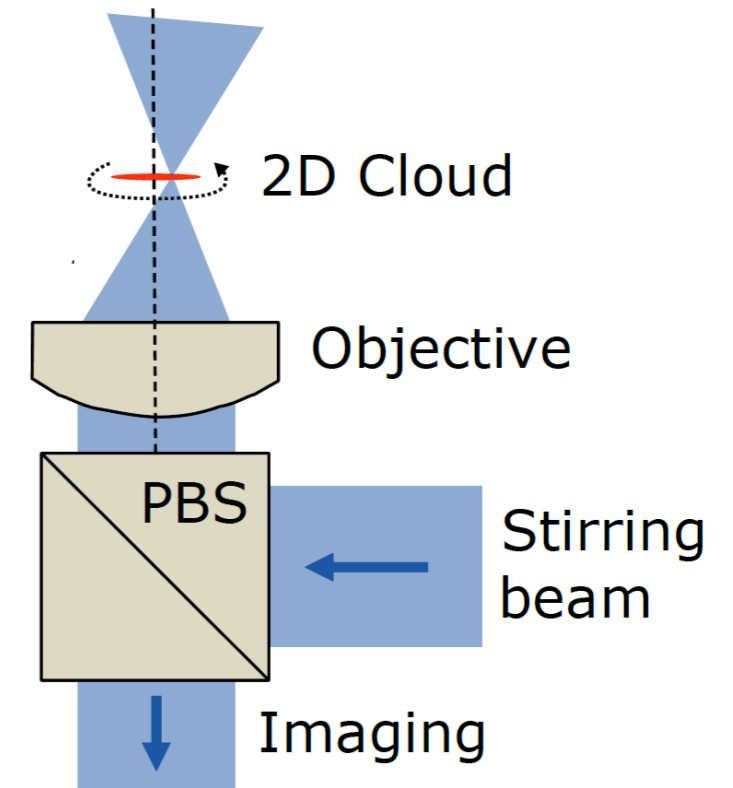
Quasi 2D regime:

$\hbar\omega_z = k_B \times 70 \text{ nK}$ comparable
to $k_B T$ and $U_{\text{int}} \sim k_B \times 40 \text{ nK}$.

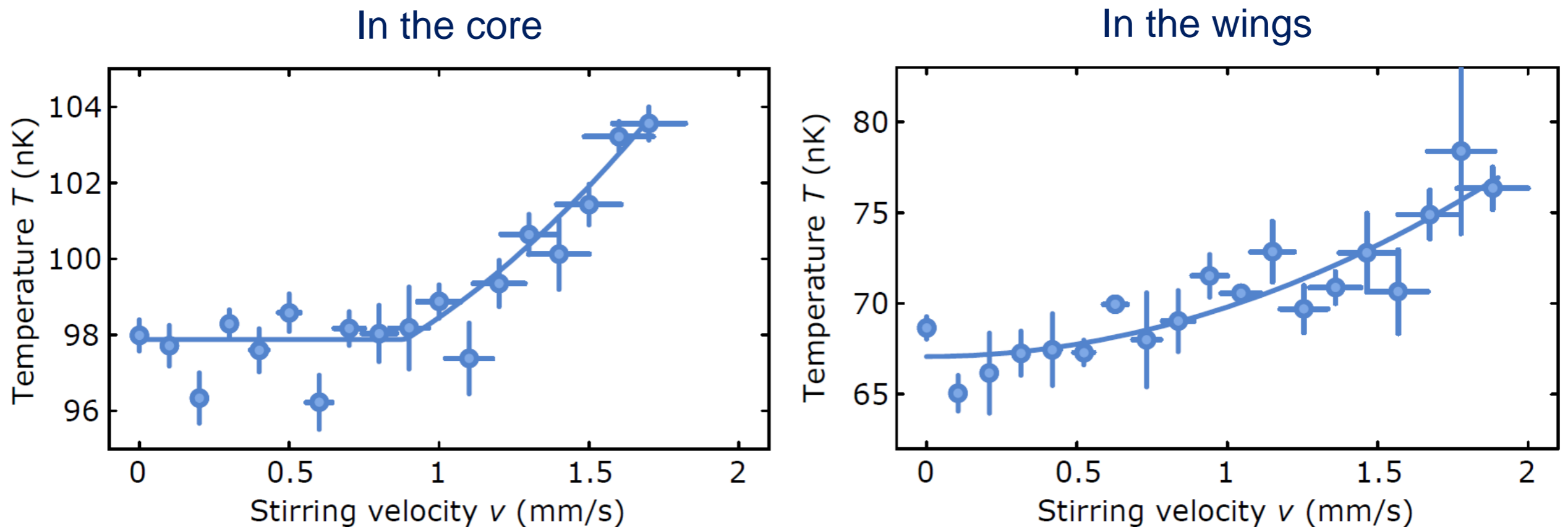
Stirring the cloud



- Use the imaging microscope objective to focus a blue-detuned laser beam (waist $w_0 = 2 \mu\text{m}$, trap depth $V = k_B \times 80 \text{ nK}$)
- Stir in circles at constant velocity v for $t_{\text{stir}} = 200 \text{ ms}$.
- Let the cloud equilibrate and measure temperature T (average over 10 shots).



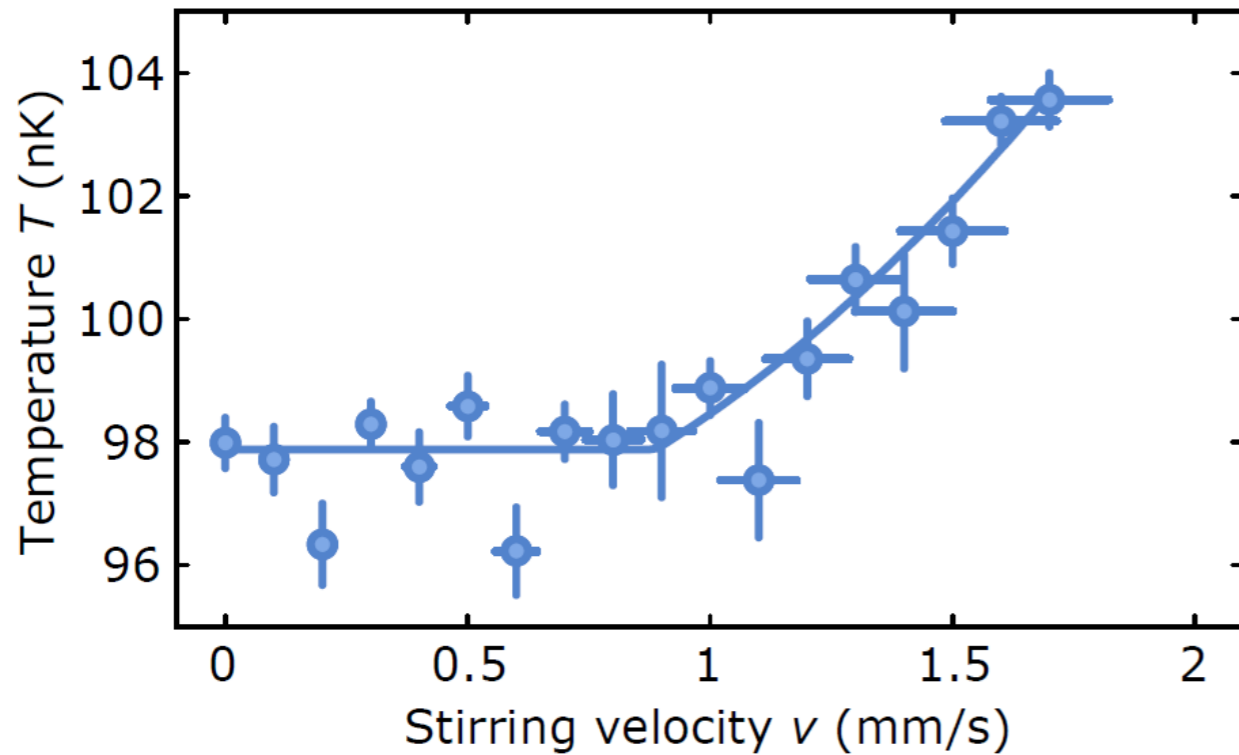
Evidence for a critical velocity



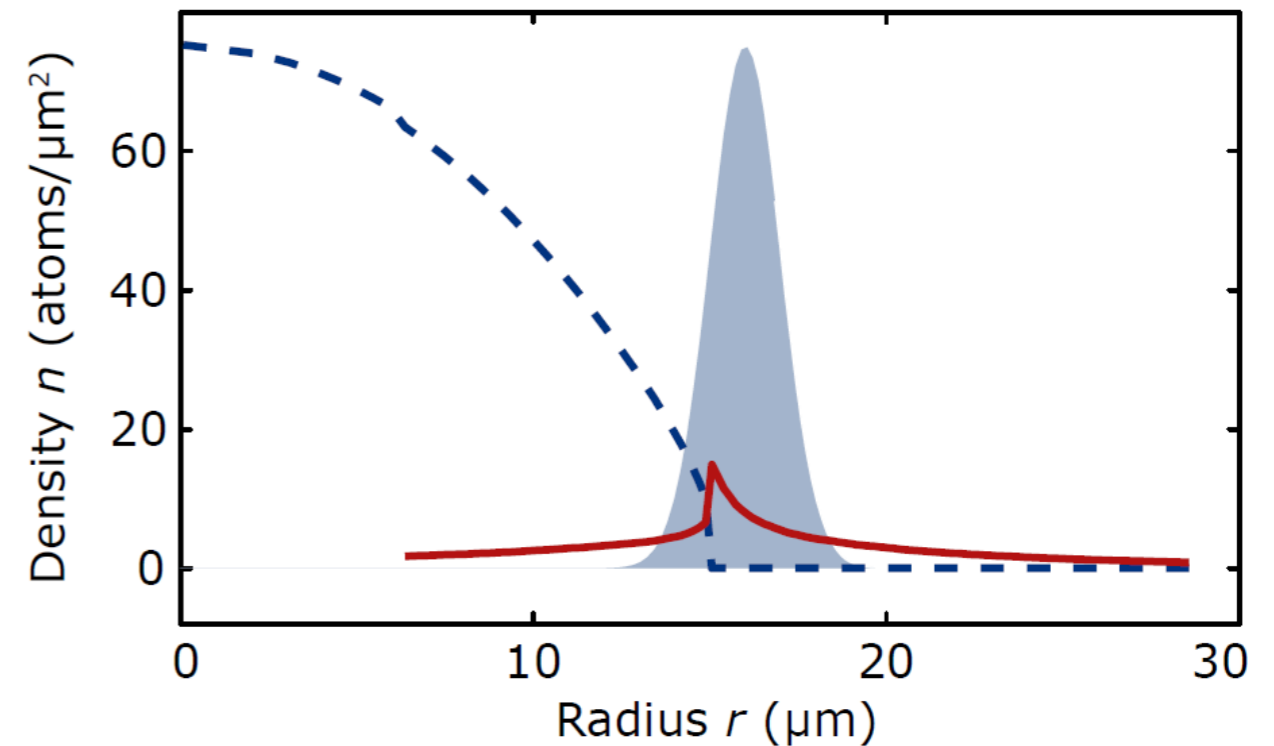
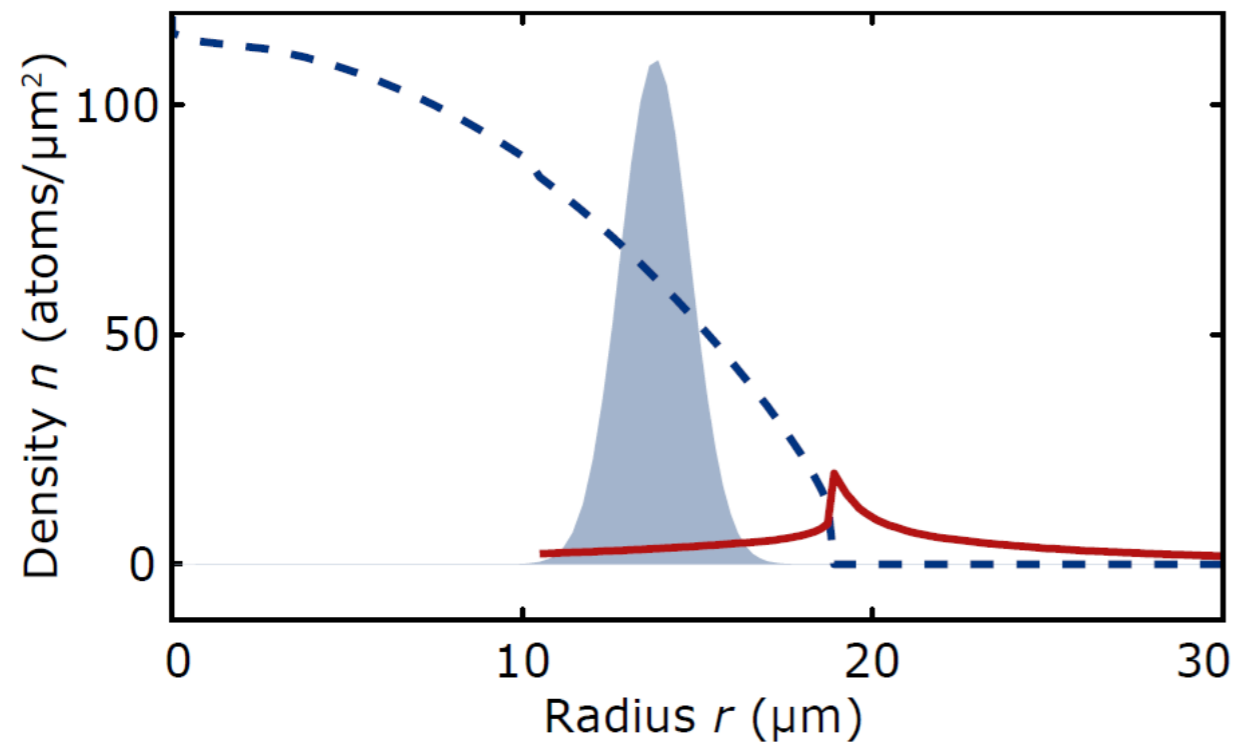
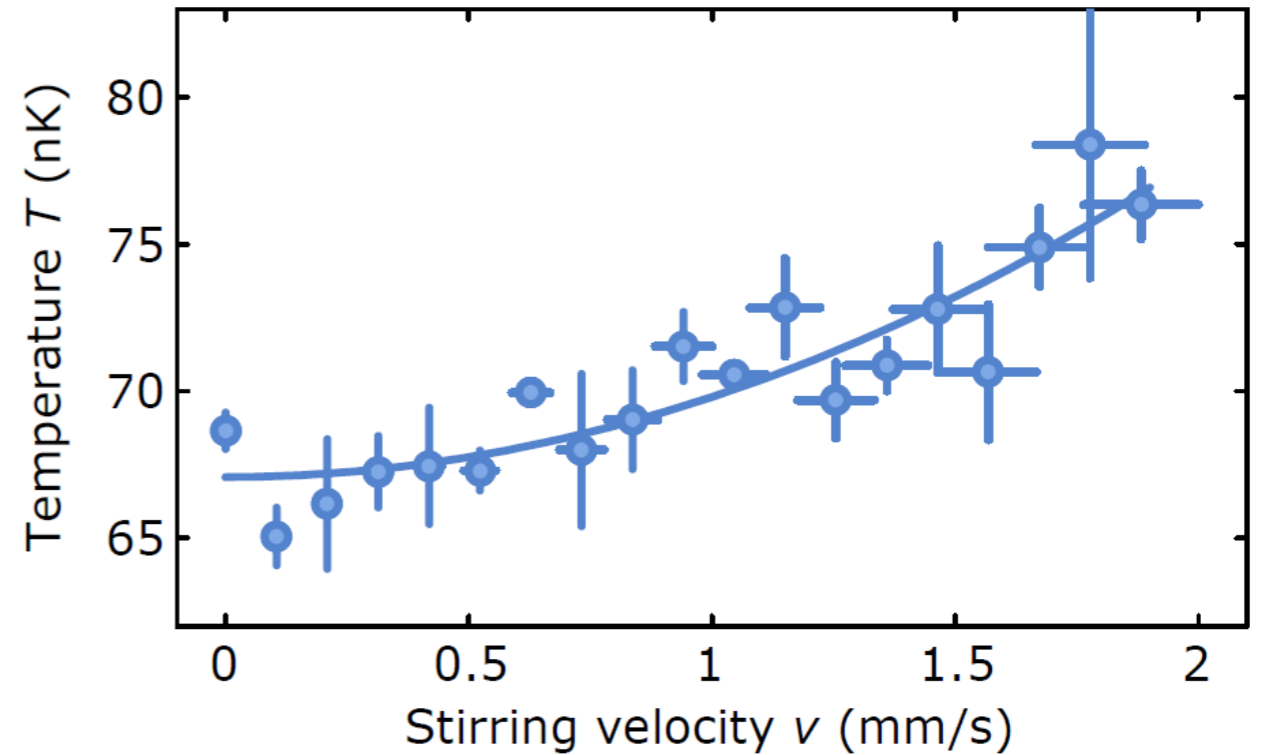
- In the core (superfluid state): clear threshold behaviour with no discernable dissipation below a critical velocity.
- In the wings (normal state): quadratic heating without threshold. Due to linear scaling of the drag force.
- Model with fit function : $T = T_{v=0} + \kappa \cdot t_{\text{stir}} \cdot \max[v^2 - v_c^2, 0]$
- Record the fit parameters κ and v_c for a given experimental configuration (N, T, r)

Evidence for a critical velocity

In the core



In the wings



The trapped 2D Bose gas

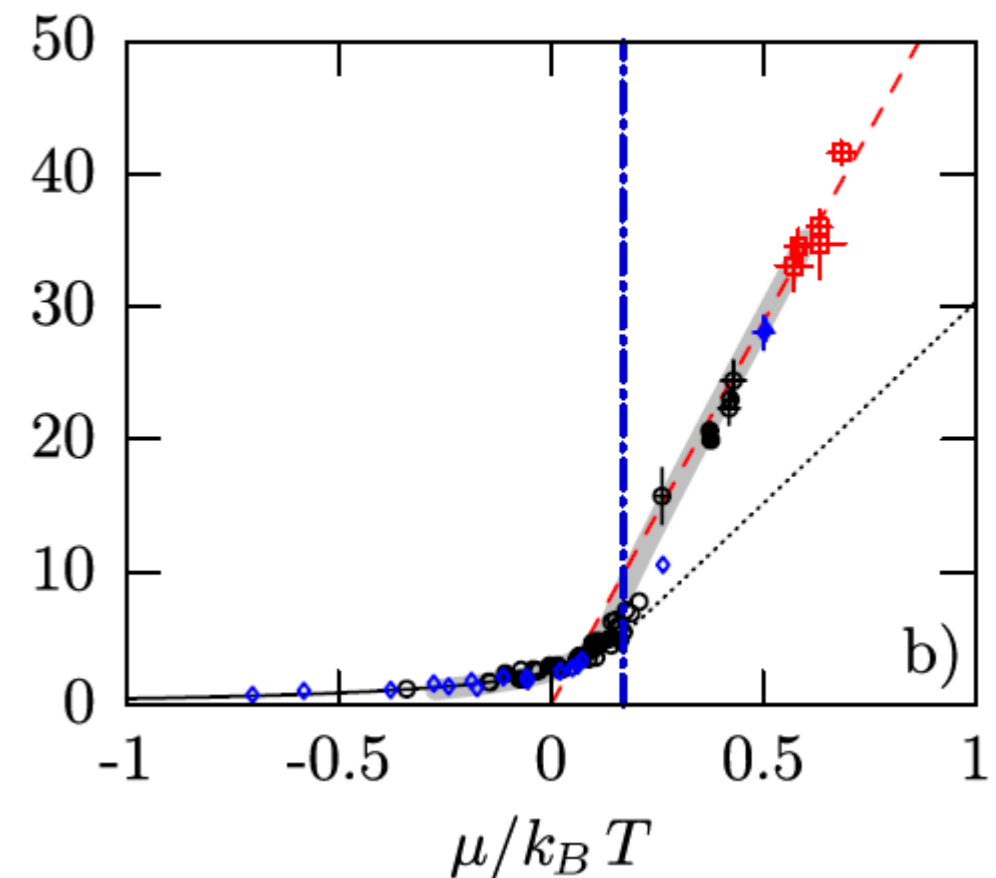
Local Density Approximation (LDA):

- For each configuration (N, T, r) , apply LDA to determine $\mu_{\text{loc}}(r) = \mu_0 - V(r)$
- relate to the uniform gas with T and μ_{loc}

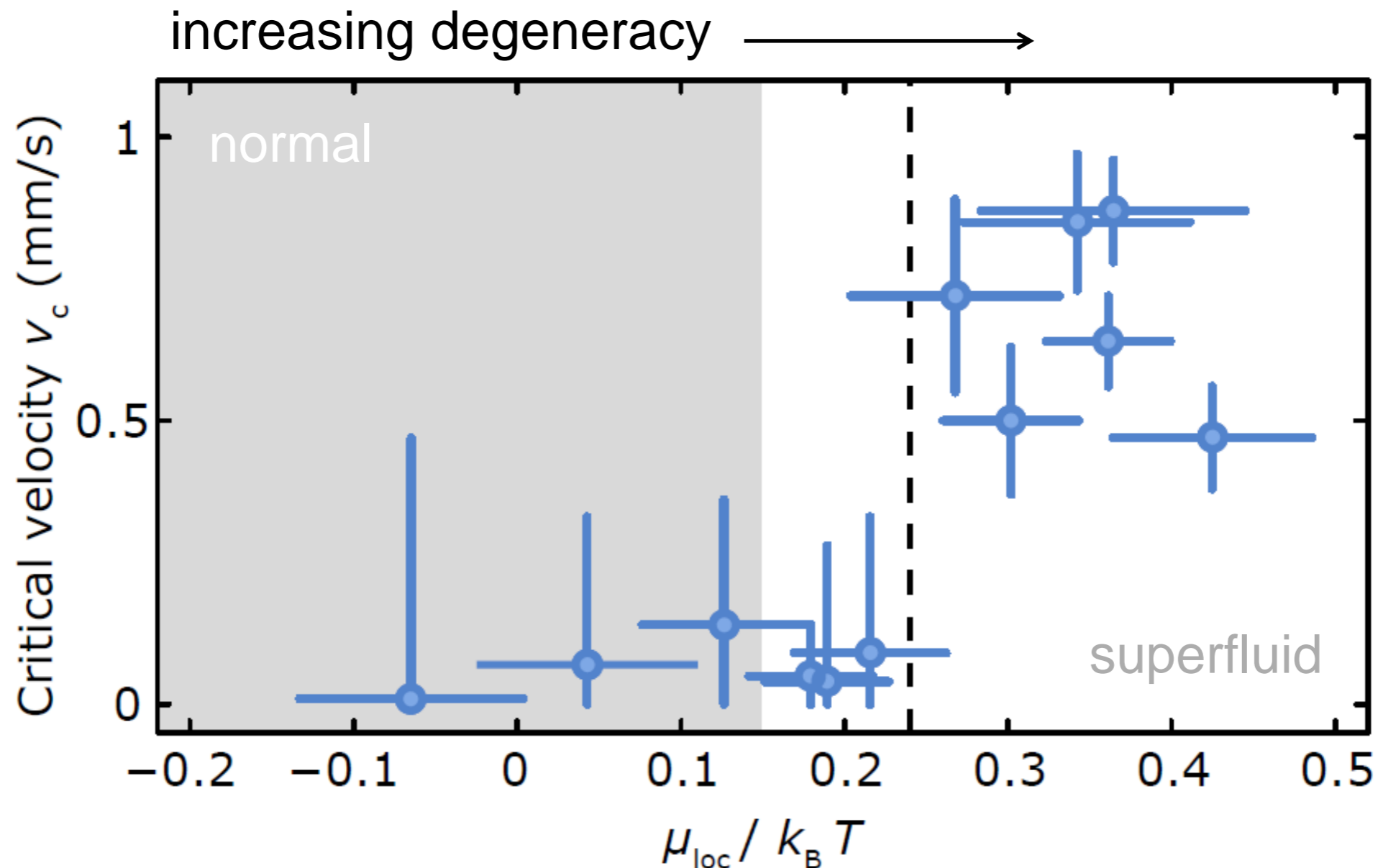
Scale invariance:

- Plot v_c against $\mu_{\text{loc}}/k_B T$.
- Due to the scale invariance of the 2D Bose gas, dimensionless quantities depend only on this ratio.
- It is univocally related to PSD and characterizes the degree of degeneracy of the cloud.

Phase space density \mathcal{D}

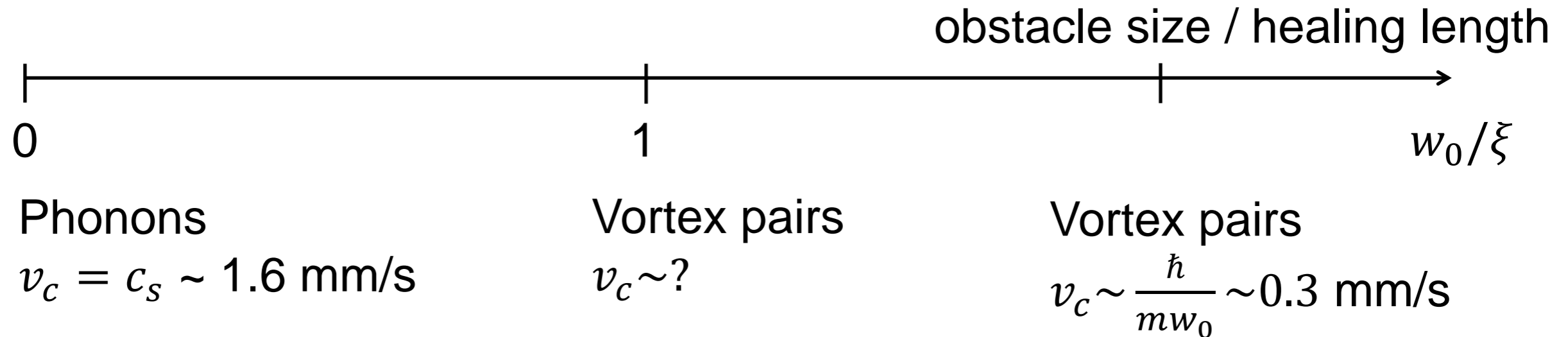


Across the BKT transition



- Two regimes: $v_c = 0$ and $v_c > 0$
- x error bars represent the finite beam size and the “background heating”
- BKT transition expected at $\mu_{\text{loc}}/k_B T = 0.15$ (for homogeneous system and $\tilde{g} = 0.093$) but shifted to 0.24, possibly due to the size of the beam.

Dissipation mechanisms



- What can we tell from the value of the critical velocity?
We find $v_c = 0.5 - 1.0 \text{ mm/s}$.
- This means $v_c/c_s = 0.3 - 0.7$ (compare to 0.1 in 3D experiments)
- We are in intermediate regime, cannot determine dominant mechanism from value of v_c .
- Simulations for similar systems suggest vortices (one should observe vortices directly)

Langer and Fischer, PRL **19**, 560 (1967)

Frisch, Pomeau, Rica, PRL **69**, 1644 (1992)

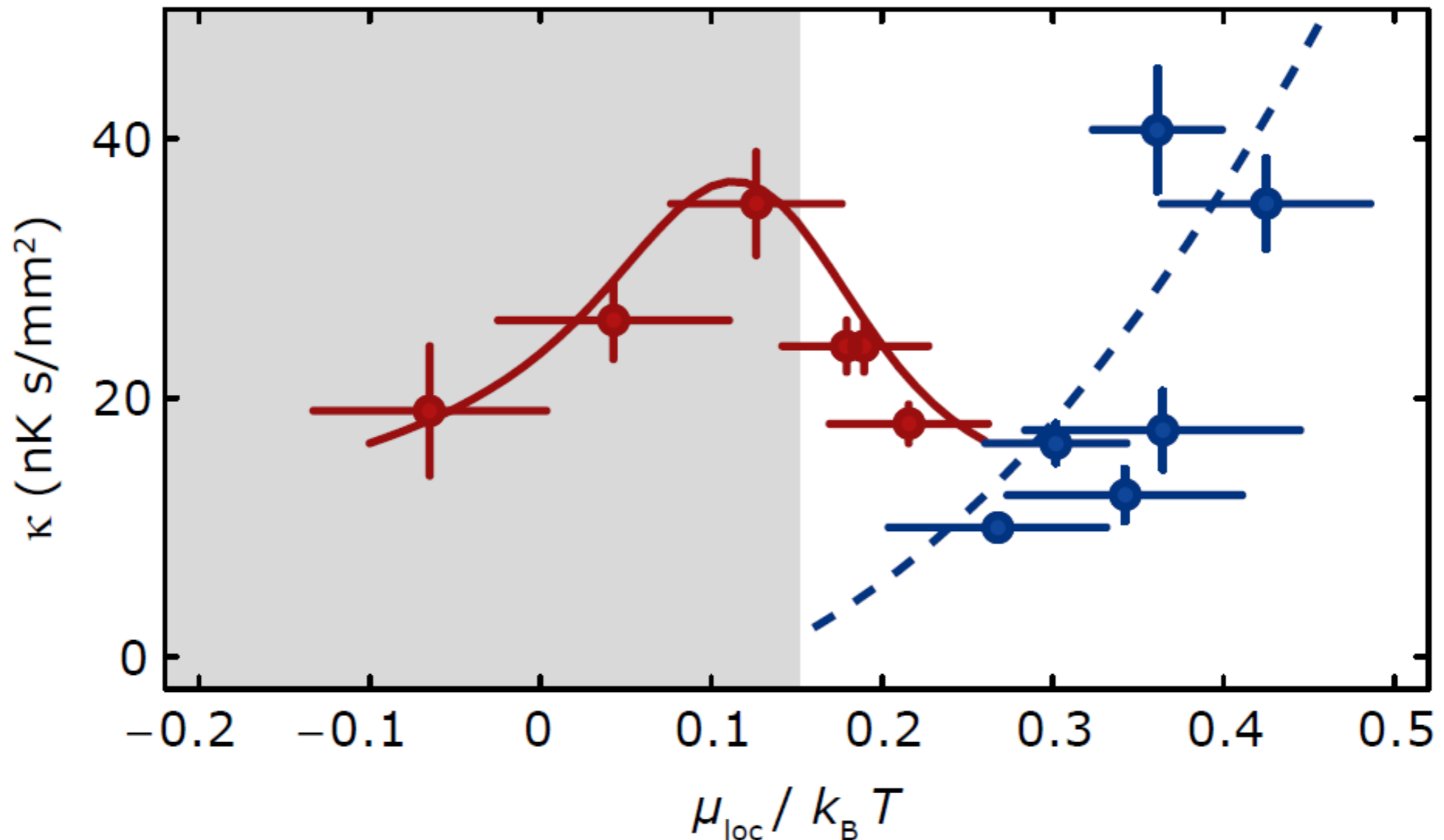
Winiiecki, McCann, Adams, PRL **82**, 5186 (1999)

Stießberger and Zwerger, PRA **62**, 061601R (2000)

Crescimanno et al., PRA **62**, 063612 (2000)

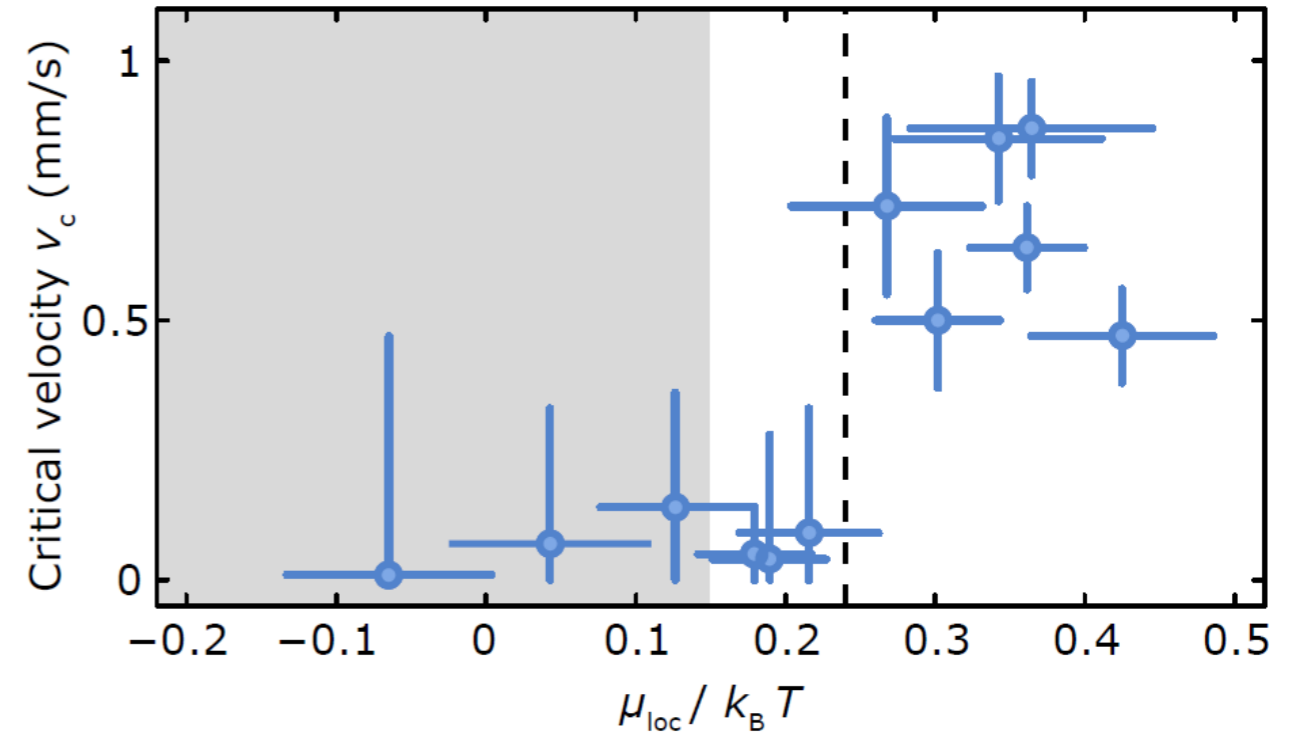
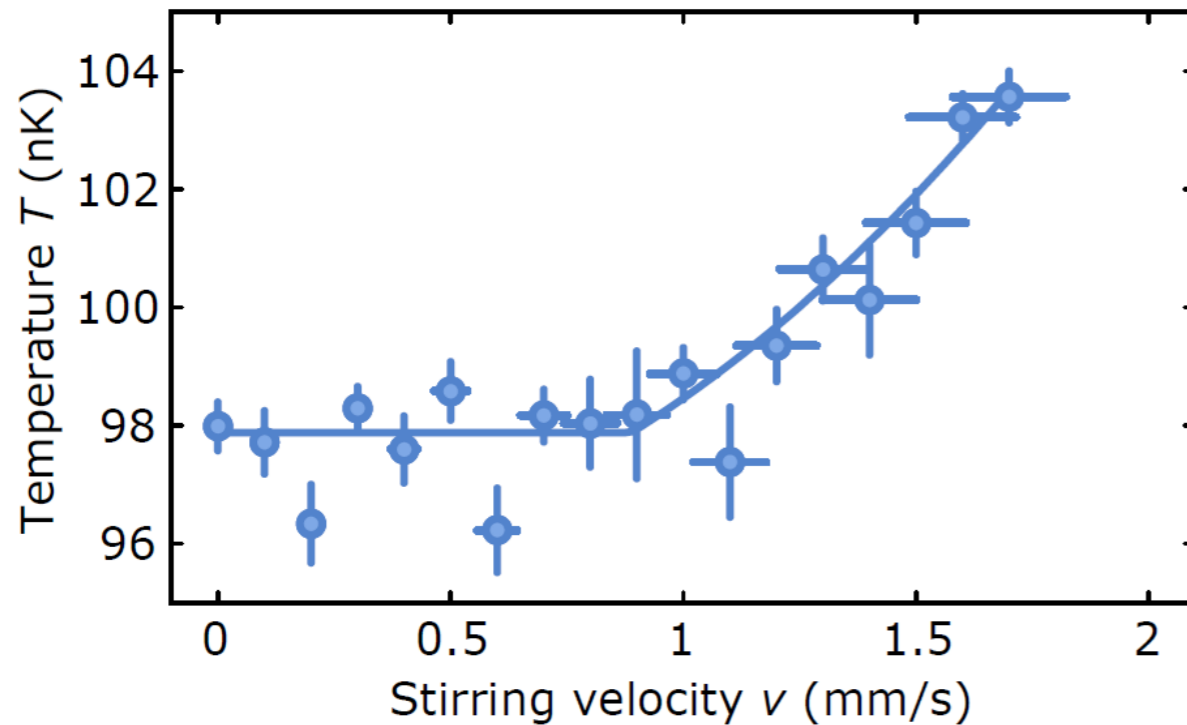
$$c_s = \hbar/m \sqrt{\tilde{g}n} \text{ (for uniform gas, } T = 0)$$
$$c_s \sim 1.6 \text{ mm/s (for } n = 50 \text{ atoms}/\mu\text{m}^2)$$

Heating coefficient



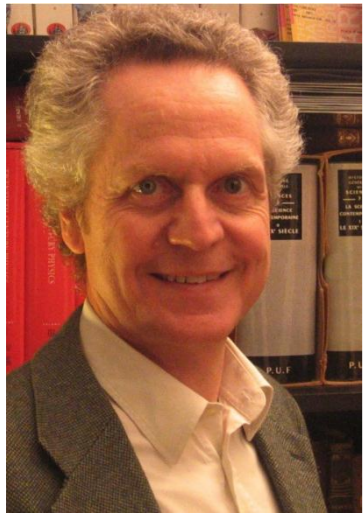
- Normal region (red): heating coefficient scales linear with the normal density (averaged over the beam size).
- Maximum of normal density around the expected superfluid transition.
- Compatible with model of single-particles colliding with beam.
- Superfluid region (blue): stronger scaling with density (here quadratic).

Summary

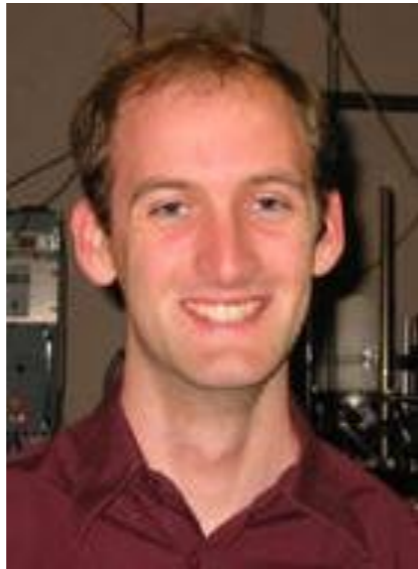


- We measured the superfluid behaviour of a 2D Bose gas via the dissipationless response to a moving obstacle.
- We probed a trapped system at fixed density, and directly observed the transition between the normal and the superfluid phase.

The Team at LKB



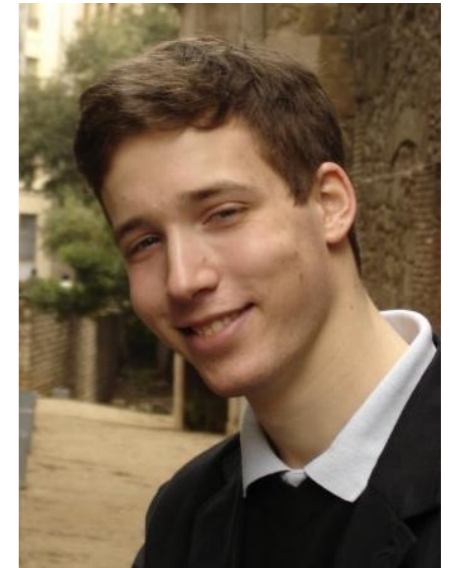
Jean Dalibard



Jérôme Beugnon



Tarik Yefsah



Rémi Desbuquois



Sylvain Nascimbène



Christof Weitenberg



Lauriane Chomaz