



stochastic perturbation theory

a prequel to Reptation Quantum Monte Carlo

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disclaimer



variational Quantum Monte Carlo

the Langevin way:

$$\begin{aligned} E_0 &\lesssim \langle \Phi_0 | \hat{H} | \Phi_0 \rangle \\ &= \int \Phi_0(R) \hat{H} \Phi_0(R) dR \end{aligned}$$

where:

$$\begin{aligned} \hat{H} &= -\frac{\partial^2}{\partial R^2} + V(R), \\ R &= \{r_1, r_2, \dots, r_N\} \end{aligned}$$



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$$\approx \frac{1}{T} \int_0^T \mathcal{W}(R(t)) dt,$$

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$$\begin{aligned} \mathcal{W}(R) &= -\frac{1}{\Phi_0(R)} \frac{\partial^2 \Phi_0(R)}{\partial R^2} + V(R) \\ dR &= \mathcal{F}(R) dt + dW_t \end{aligned}$$

$$\begin{aligned} \langle dW_t^2 \rangle &= 2dt \\ \mathcal{F}(R) &= -\frac{\partial \mathcal{U}(R)}{\partial R} \\ \mathcal{U}(R) &= -\log(\Phi_0(R)^2) \\ \mathcal{F}(R) &= 2 \frac{1}{\Phi_0} \frac{\partial \Phi_0}{\partial X} \end{aligned}$$



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$$\approx \frac{1}{T} \int_0^T \mathcal{W}(R(t)) dt,$$

estimating the statistical error:

$$\Delta E_T^2 \approx \frac{1}{T^2} \left\langle \left[\int_0^T (\mathcal{W}(R(t)) - \langle \mathcal{W} \rangle) dt \right]^2 \right\rangle$$

where:

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where:

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Langevin dynamics

$$dR = \mathcal{F}(R)dt + dW_t$$

$$\langle dW_t^2 \rangle = 2dt$$
$$\mathcal{F}(R) = -\frac{\partial \mathcal{U}(R)}{\partial R}$$

$$\frac{\partial P(R, t)}{\partial t} = \frac{\partial^2 P(R, t)}{\partial R^2} - \frac{\partial}{\partial R} (\mathcal{F}(R)P(R, t))$$

$$P_0(R) \propto e^{-\mathcal{U}(R)}$$



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$$\hat{\mathcal{H}} = -\frac{\partial^2}{\partial R^2} + \mathcal{V}(R)$$

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$$\mathcal{V}(R) = \frac{1}{4}\mathcal{F}(R)^2 - \frac{1}{2}\Delta\mathcal{U}(R)$$
$$= \frac{1}{\Phi_0(R)} \frac{\partial^2 \Phi_0(R)}{\partial R^2}$$

$$\Phi(R, t) = P(R, t)/\Phi_0(R)$$
$$\Phi_0(R) = \sqrt{P_0(R)}$$



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$$\hat{\mathcal{H}}\Phi_0 = 0$$

$$\mathcal{E}_0 = 0; \quad \mathcal{E}_n > 0$$



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$$\begin{aligned} P(R, t) &= \int W(R, t|R', 0)P(R', 0)dR' \\ &= \Phi_0(R)\Phi(R, t) \\ &= \Phi_0(R) \underbrace{\int \mathcal{G}(R, t; R', 0)}_{\langle R | e^{-\hat{\mathcal{H}}(t-t')} | R' \rangle} \Phi(R', 0)dR' \end{aligned}$$



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$$W(R, t|R', 0) = \Phi_0(R)\mathcal{G}(R, t; R', 0)/\Phi_0(R')$$



computing the local-energy correlation time

summary:

$$dR = \mathcal{F}(R)dR + dW_t$$

$$\frac{\partial P(R, t)}{\partial t} = \frac{\partial^2 P(R, t)}{\partial R^2} - \frac{\partial}{\partial R}(\mathcal{F}(R)P(R, t))$$

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computing the local-energy correlation time

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$$dR = \mathcal{F}(R)dR + dW_t$$

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$$\begin{aligned}
 \langle \Delta \mathcal{W}(t) \Delta \mathcal{W}(0) \rangle &= \int \Delta \mathcal{W}(R) \Delta \mathcal{W}(R') P(R, t; R', 0) dR dR' \\
 &= \int \Delta \mathcal{W}(R) \Delta \mathcal{W}(R') W(R, t|R', 0) P_0(R) dR dR' \\
 &= \langle \Phi_0 | \Delta \widehat{\mathcal{W}} \widehat{\mathcal{G}}(t) \Delta \widehat{\mathcal{W}} | \Phi_0 \rangle
 \end{aligned}$$

$$\begin{aligned}
 \tau_{\mathcal{W}} \Delta \mathcal{W}^2 &\doteq \int_0^\infty \langle \Phi_0 | \Delta \widehat{\mathcal{W}} \widehat{\mathcal{G}}(t) \Delta \widehat{\mathcal{W}} | \Phi_0 \rangle dt \\
 &= \sum_{n>0} \frac{|\mathcal{W}_{0n}|^2}{\mathcal{E}_n}
 \end{aligned}$$



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$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$\begin{aligned} \langle \Delta \mathcal{W}(t) \Delta \mathcal{W}(0) \rangle &= \int \Delta \mathcal{W}(R) \Delta \mathcal{W}(R') P(R, t; R', 0) dR dR' \\ &= \int \Delta \mathcal{W}(R) \Delta \mathcal{W}(R') W(R, t|R', 0) P_0(R) dR dR' \\ &= \langle \Phi_0 | \Delta \widehat{\mathcal{W}} \widehat{\mathcal{G}}(t) \Delta \widehat{\mathcal{W}} | \Phi_0 \rangle \end{aligned}$$

$$\begin{aligned} \tau_{\mathcal{W}} \Delta \mathcal{W}^2 &\doteq \int_0^\infty \langle \Phi_0 | \Delta \widehat{\mathcal{W}} \widehat{\mathcal{G}}(t) \Delta \widehat{\mathcal{W}} | \Phi_0 \rangle dt \\ &= \sum_{n>0} \frac{|\mathcal{W}_{0n}|^2}{\mathcal{E}_n} \end{aligned}$$

$$\tau_{\mathcal{W}} = \frac{1}{\Delta \mathcal{W}^2} \sum_{n>0} \frac{|\mathcal{W}_{0n}|^2}{\mathcal{E}_n}$$



second-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$\begin{aligned}\tau_{\mathcal{W}} &= \frac{1}{\Delta\mathcal{W}^2} \sum_{n>0} \frac{|\mathcal{W}_{0n}|^2}{\mathcal{E}_n} \\ &= -\frac{E_0^{(2)}}{\Delta\mathcal{W}^2}\end{aligned}$$

$$\begin{aligned}E_0 &= \mathcal{E}_0 + \langle \Phi_0 | \widehat{\mathcal{W}} | \Phi_0 \rangle_{QM} - \sum_{n>0} \frac{|\langle \Phi_0 | \widehat{\mathcal{W}} | \Phi_n \rangle_{QM}|^2}{\mathcal{E}_n} + \dots \\ &\sim \underbrace{\frac{d}{dT} \left[\left\langle \int_0^T \mathcal{W}(t) dt \right\rangle_{RW} \right]}_{\mu_1[\mathcal{W}]} - \underbrace{\frac{1}{2} \left\langle \left(\int_0^T \Delta\mathcal{W}(t) dt \right)^2 \right\rangle_{RW} + \dots? \dots}_{\bar{\mu}_2[\mathcal{W}]}\end{aligned}$$



higher-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$\Phi_0 = c_0 \Psi_0 + c_1 \Psi_1 + \dots$$

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle = c_0^2 e^{-E_0 T} + c_1^2 e^{-E_1 T} + \dots$$



higher-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

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$$\log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle = -E_0 T + \log (c_0^2 + c_1^2 e^{-\Delta E_1 T} + \dots)$$



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$$E_0 = -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle + \mathcal{O}(e^{-\Delta E_1 T})$$



higher-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$E_0 = -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle$$

$$\begin{aligned} e^{-\hat{H}T} = & e^{-\hat{\mathcal{H}}T} \left(1 - \int_0^T dt_1 \widehat{\mathcal{W}}(t_1) dt_1 + \int_0^T dt_2 \widehat{\mathcal{W}}(t_2) \int_0^{t_2} \widehat{\mathcal{W}}(t_1) dt_1 + \right. \\ & \left. \dots \frac{(-)^n}{n!} \int_0^T \dots \int_0^T \mathcal{T} \left(\widehat{\mathcal{W}}(t_1) \dots \widehat{\mathcal{W}}(t_n) \right) dt_1 \dots dt_n + \dots \right) \end{aligned}$$



higher-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$E_0 = -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle$$

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$$\left. \dots \frac{(-)^n}{n!} \int_0^T \dots \int_0^T \mathcal{T} \left(\widehat{\mathcal{W}}(t_1) \dots \widehat{\mathcal{W}}(t_n) \right) dt_1 \dots dt_n + \dots \right)$$
$$\widehat{\mathcal{W}}(t) \doteq e^{\hat{\mathcal{H}}t} \widehat{\mathcal{W}} e^{-\hat{\mathcal{H}}t}$$



higher-order perturbation theory

$$\hat{H} = \hat{\mathcal{H}} + \mathcal{W}(R)$$

$$E_0 = -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle$$

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$$\widehat{\mathcal{W}}(t) \doteq e^{\hat{\mathcal{H}}t} \widehat{\mathcal{W}} e^{-\hat{\mathcal{H}}t}$$

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2} \mu_2 + \dots + \frac{(-)^n}{n!} \mu_n + \dots$$

$$\mu_n = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \dots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \dots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$



higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\mu_n = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$

$$\begin{aligned} \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM} &= \int \Phi_0(R_n) \mathcal{W}(R_n) \mathcal{G}(R_n, t_n, R_{n-1}, t_{n-1}) \\ &\quad \times \cdots \mathcal{W}(R_1) \mathcal{G}(R_2, t_2; R_2, t_1) \mathcal{W}(R_1) \Phi_0(R_1) \, dR_n \cdots dR_1 \end{aligned}$$



higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\mu_n = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$

$$\begin{aligned} \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM} &= \int \Phi_0(R_n) \mathcal{W}(R_n) \mathcal{G}(R_n, t_n, R_{n-1}, t_{n-1}) \\ &\quad \times \cdots \mathcal{W}(R_1) \mathcal{G}(R_2, t_2; R_2, t_1) \mathcal{W}(R_1) \Phi_0(R_1) \, dR_n \cdots dR_1 \\ &= \int W(R_n, t_n | R_{n-1}, t_{n-1}) \cdots W(R_2, t_2 | R_1, t_1) P_0(R_1) \\ &\quad \times \mathcal{W}(R_n) \cdots \mathcal{W}(R_1) \, dR_n \cdots dR_1 \end{aligned}$$



higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\mu_n = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$

$$\begin{aligned} \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM} &= \int \Phi_0(R_n) \mathcal{W}(R_n) \mathcal{G}(R_n, t_n, R_{n-1}, t_{n-1}) \\ &\quad \times \cdots \mathcal{W}(R_1) \mathcal{G}(R_2, t_2; R_2, t_1) \mathcal{W}(R_1) \Phi_0(R_1) \, dR_n \cdots dR_1 \\ &= \int W(R_n, t_n | R_{n-1}, t_{n-1}) \cdots W(R_2, t_2 | R_1, t_1) P_0(R_1) \\ &\quad \times \mathcal{W}(R_n) \cdots \mathcal{W}(R_1) \, dR_n \cdots dR_1 \\ &\doteq \langle \mathcal{W}(t_n) \mathcal{W}(t_{n-1}) \cdots \mathcal{W}(t_1) \rangle_{RW} \end{aligned}$$



higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\begin{aligned}\mu_n &= n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM} \\ &= \left\langle \left[\int_0^T \mathcal{W}(t) dt \right]^n \right\rangle_{RW}\end{aligned}$$

$$\begin{aligned}\langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM} &= \int \Phi_0(R_n) \mathcal{W}(R_n) \mathcal{G}(R_n, t_n, R_{n-1}, t_{n-1}) \\ &\quad \times \cdots \mathcal{W}(R_1) \mathcal{G}(R_2, t_2; R_2, t_1) \mathcal{W}(R_1) \Phi_0(R_1) dR_n \cdots dR_1 \\ &= \int W(R_n, t_n | R_{n-1}, t_{n-1}) \cdots W(R_2, t_2 | R_1, t_1) P_0(R_1) \\ &\quad \times \mathcal{W}(R_n) \cdots \mathcal{W}(R_1) dR_n \cdots dR_1 \\ &\doteq \langle \mathcal{W}(t_n) \mathcal{W}(t_{n-1}) \cdots \mathcal{W}(t_1) \rangle_{RW}\end{aligned}$$



infinite-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t) dt \right]^n \right\rangle_{RW}$$



infinite-order perturbation theory

$$\begin{aligned}\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots \\ &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\ &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}\end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$



infinite-order perturbation theory

$$\begin{aligned}\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots + \frac{(-)^n}{n!}\mu_n + \cdots \\ &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\ &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}\end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$

$$E_0 \sim \frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM}$$



infinite-order perturbation theory

$$\begin{aligned}\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots + \frac{(-)^n}{n!}\mu_n + \cdots \\ &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\ &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}\end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$

$$\begin{aligned}E_0 &\sim \frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} \\ &= \frac{\langle \mathcal{W}(T) e^{-\mathcal{S}(T)} \rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}}\end{aligned}$$



infinite-order perturbation theory

$$\begin{aligned}\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots + \frac{(-)^n}{n!}\mu_n + \cdots \\ &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\ &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}\end{aligned}$$

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$$\begin{aligned}E_0 &\sim \frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} \\ &= \frac{\langle \mathcal{W}(T) e^{-\mathcal{S}(T)} \rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \\ &= \frac{\left\langle \frac{\mathcal{S}(T)}{T} e^{-\mathcal{S}(T)} \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}}\end{aligned}$$



infinite-order perturbation theory

$$\begin{aligned}\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots + \frac{(-)^n}{n!}\mu_n + \cdots \\ &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\ &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}\end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$

$$\widehat{H}(\boldsymbol{\lambda}) = \widehat{H} + \sum_i \lambda_i \widehat{A}_i$$

$$\begin{aligned}E_0 &\sim \frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} \\ &= \frac{\langle \mathcal{W}(T) e^{-\mathcal{S}(T)} \rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \\ &= \frac{\left\langle \frac{\mathcal{S}(T)}{T} e^{-\mathcal{S}(T)} \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}}\end{aligned}$$



infinite-order perturbation theory

$$\begin{aligned}
 \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots \\
 &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\
 &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}
 \end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$

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$$\begin{aligned}
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 &= \frac{\langle \mathcal{W}(T) e^{-\mathcal{S}(T)} \rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \\
 &= \frac{\left\langle \frac{\mathcal{S}(T)}{T} e^{-\mathcal{S}(T)} \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}}
 \end{aligned}$$

$$\begin{aligned}
 \langle \widehat{A}_i \rangle &= \left. \frac{\partial E_0(\boldsymbol{\lambda})}{\partial \lambda_i} \right|_{\boldsymbol{\lambda}=0} \\
 &= \frac{\left\langle e^{-\mathcal{S}(T)} \frac{1}{T} \int_0^T \mathcal{A}(t)dt \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \doteq \left\langle \left\langle \frac{1}{T} \int_0^T \mathcal{A}(t)dt \right\rangle \right\rangle
 \end{aligned}$$



infinite-order perturbation theory

$$\begin{aligned}
 \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} &= 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots \\
 &= \left\langle e^{-\int_0^T \mathcal{W}(t)dt} \right\rangle_{RW} \\
 &\doteq \left\langle e^{-\mathcal{S}(T)} \right\rangle_{RW}
 \end{aligned}$$

$$\mu_n = \left\langle \left[\int_0^T \mathcal{W}(t)dt \right]^n \right\rangle_{RW}$$

$$\widehat{H}(\boldsymbol{\lambda}) = \widehat{H} + \sum_i \lambda_i \widehat{A}_i$$

$$\begin{aligned}
 E_0 &\sim \frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} \\
 &= \frac{\langle \mathcal{W}(T) e^{-\mathcal{S}(T)} \rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \\
 &= \frac{\left\langle \frac{\mathcal{S}(T)}{T} e^{-\mathcal{S}(T)} \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}}
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$$\begin{aligned}
 \langle \widehat{A}_i \rangle &= \frac{\partial E_0(\boldsymbol{\lambda})}{\partial \lambda_i} \Big|_{\boldsymbol{\lambda}=0} \\
 &= \frac{\left\langle e^{-\mathcal{S}(T)} \frac{1}{T} \int_0^T \mathcal{A}(t)dt \right\rangle_{RW}}{\langle e^{-\mathcal{S}(T)} \rangle_{RW}} \doteq \left\langle \left\langle \frac{1}{T} \int_0^T \mathcal{A}(t)dt \right\rangle \right\rangle
 \end{aligned}$$

$$\frac{\partial \langle \widehat{A}_i \rangle}{\partial \lambda_j} = 2 \int_0^T C_{ij}(t) dt$$

$$\begin{aligned}
 C_{ij}(t) &= \langle\langle \Delta \mathcal{A}_i(t) \Delta \mathcal{A}_j(0) \rangle\rangle \\
 &= \langle \Delta \widehat{A}_i(-it) \Delta \widehat{A}_j(0) \rangle_{QM}
 \end{aligned}$$



Reptation Quantum Monte Carlo

$$\langle \hat{A} \rangle \sim \left\langle e^{-\mathcal{S}} \frac{1}{T} \int_0^T \mathcal{A}(t) dt \right\rangle_{RW}$$

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle \sim \left\langle e^{-\int_0^T \mathcal{E}(t) dt} \right\rangle_{RW}$$

$$\langle A(-it)A(0) \rangle = \left\langle e^{-\mathcal{S}} \frac{1}{T} \int_0^T A(t' + t)A(t') dt \right\rangle_{RW}$$

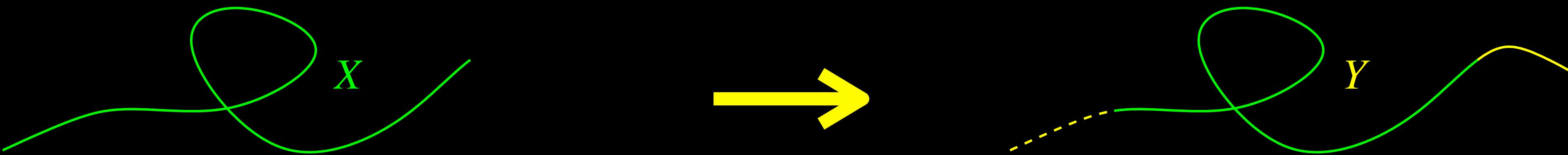


Reptation Quantum Monte Carlo

$$\langle \hat{A} \rangle \sim \left\langle e^{-\mathcal{S}} \frac{1}{T} \int_0^T \mathcal{A}(t) dt \right\rangle_{RW}$$

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle \sim \left\langle e^{-\int_0^T \mathcal{E}(t) dt} \right\rangle_{RW}$$

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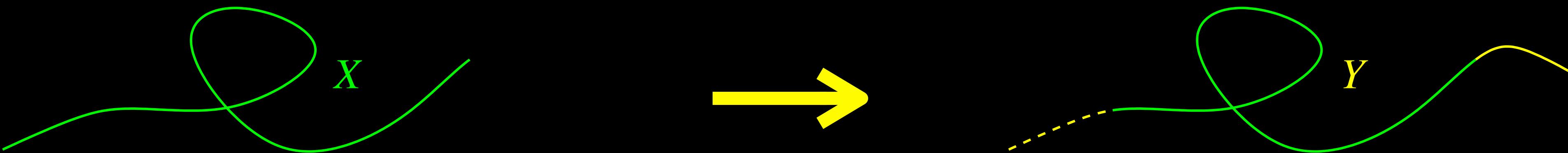


Reptation Quantum Monte Carlo

$$\langle \hat{A} \rangle \sim \left\langle e^{-\mathcal{S}} \frac{1}{T} \int_0^T \mathcal{A}(t) dt \right\rangle_{RW}$$

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle \sim \left\langle e^{-\int_0^T \mathcal{E}(t) dt} \right\rangle_{RW}$$

$$\langle A(-it)A(0) \rangle = \left\langle e^{-\mathcal{S}} \frac{1}{T} \int_0^T A(t' + t)A(t') dt \right\rangle_{RW}$$



$$\text{accept}[X \rightarrow Y] = \min(1, e^{-\Delta S})$$

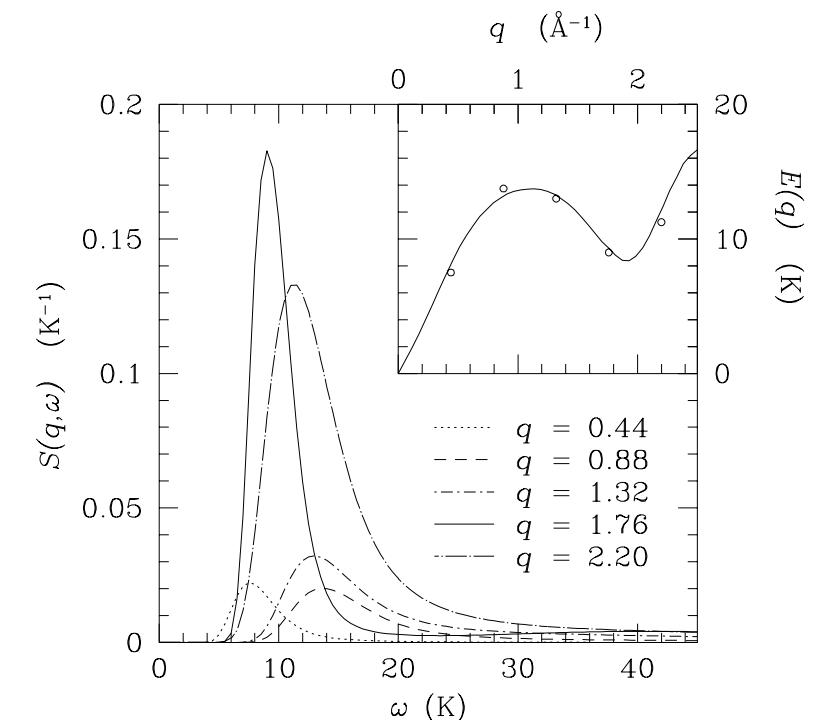
a few highlights

REPTATION QUANTUM MONTE CARLO

a round-trip tour from classical diffusion to quantum mechanics

STEFANO BARONI* AND SAVERIO MORONI†

Proceedings of the NATO ASI Quantum Monte Carlo Methods in Physics
and Chemistry, Cornell University, July, 12-25, 1998



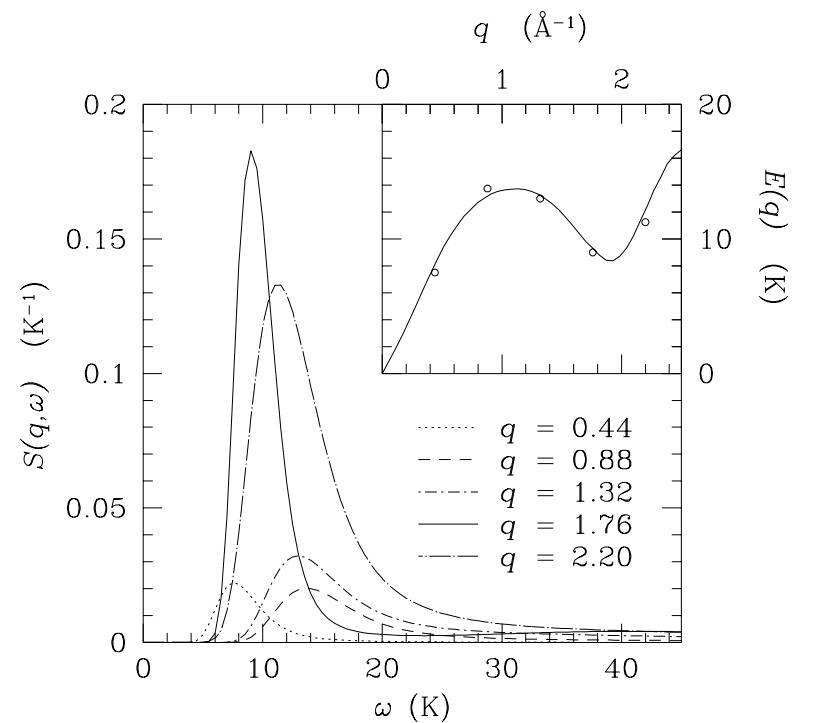
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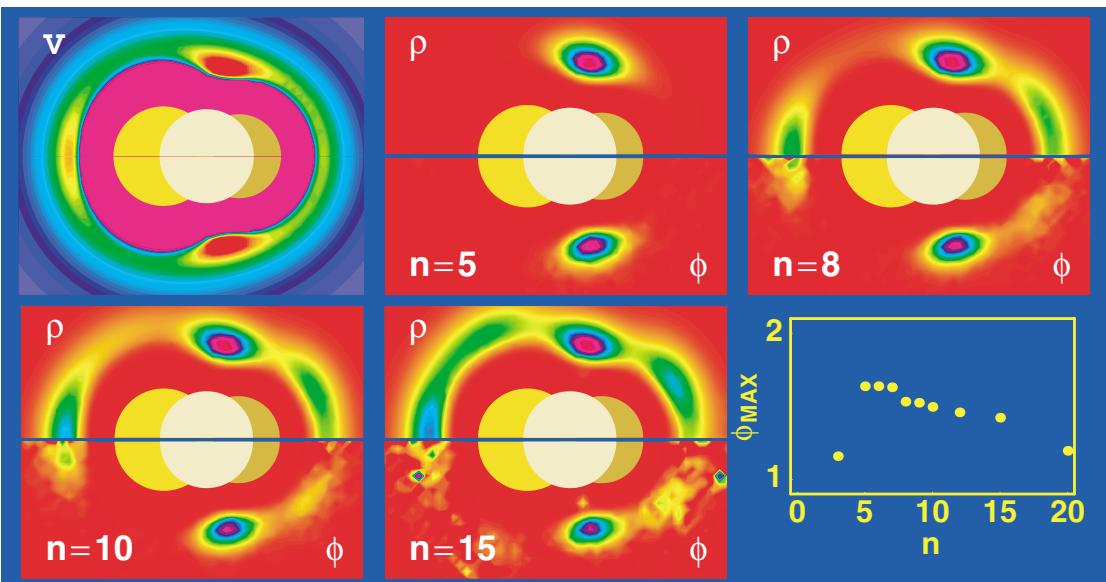
VOLUME 90, NUMBER 14

PHYSICAL REVIEW LETTERS

week ending
11 APRIL 2003

Structure, Rotational Dynamics, and Superfluidity of Small OCS-Doped He Clusters

Saverio Moroni,¹ Antonio Sarsa,^{2,*} Stefano Fantoni,² Kevin E. Schmidt,^{2,†} and Stefano Baroni^{2,3,‡}



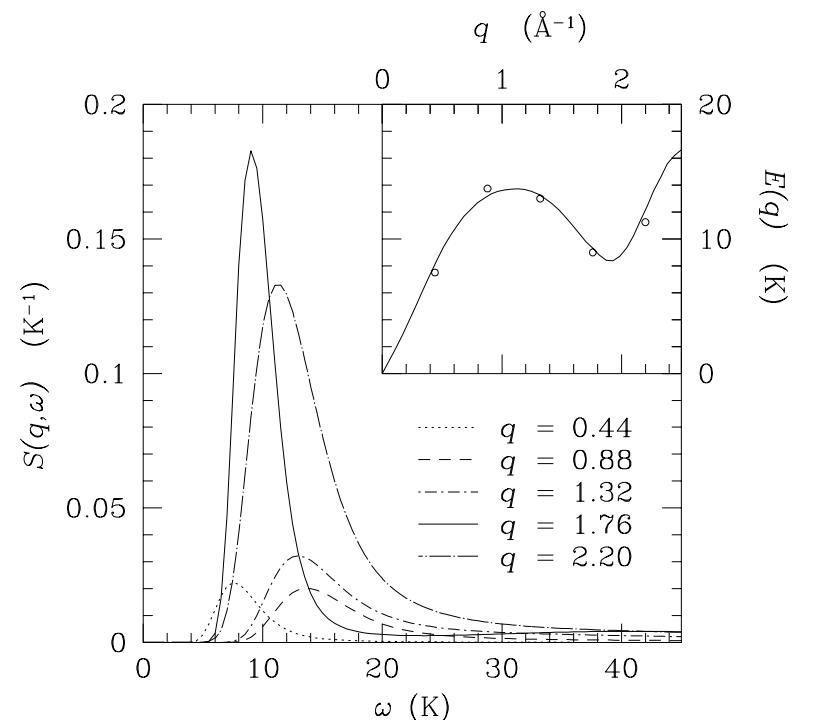
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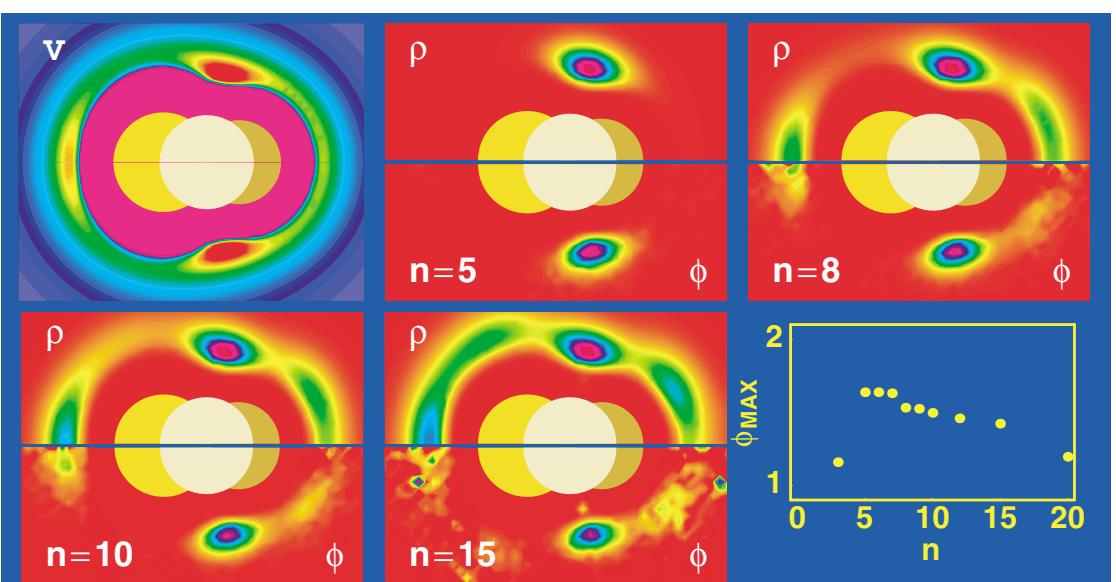
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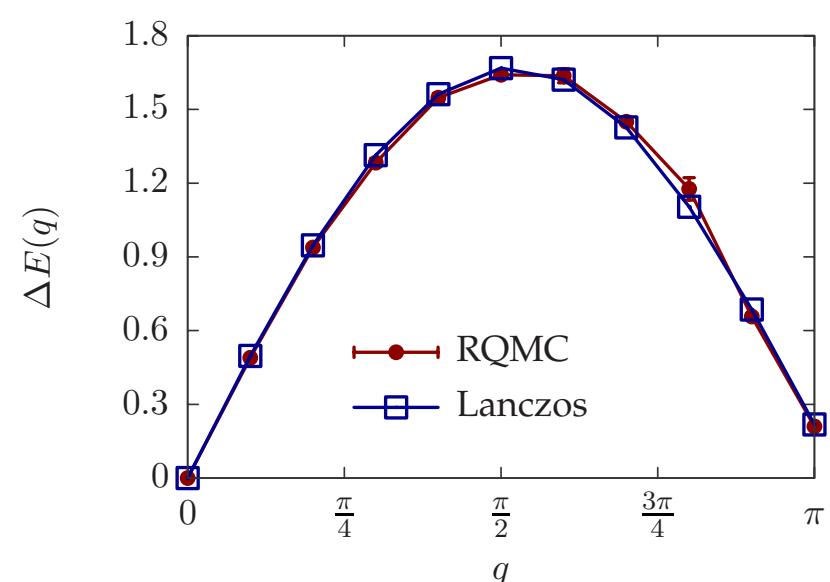
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PHYSICAL REVIEW E 82, 046710 (2010)

Reptation quantum Monte Carlo algorithm for lattice Hamiltonians with a directed-update scheme

Giuseppe Carleo, Federico Becca, Saverio Moroni, and Stefano Baroni



a few highlights

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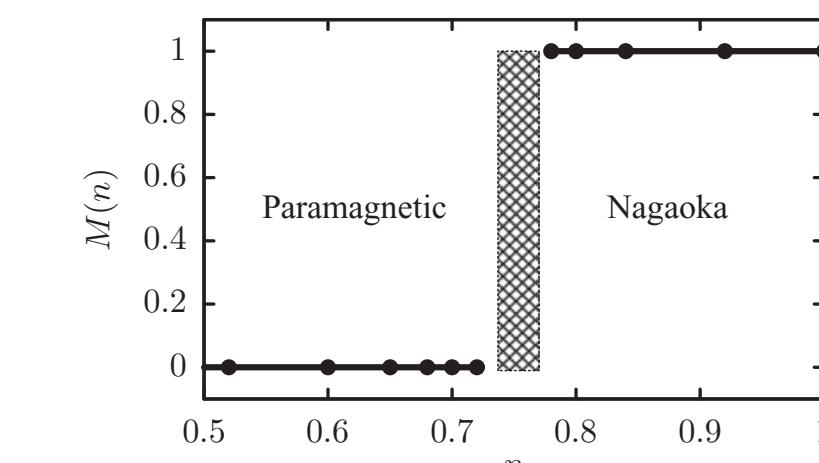
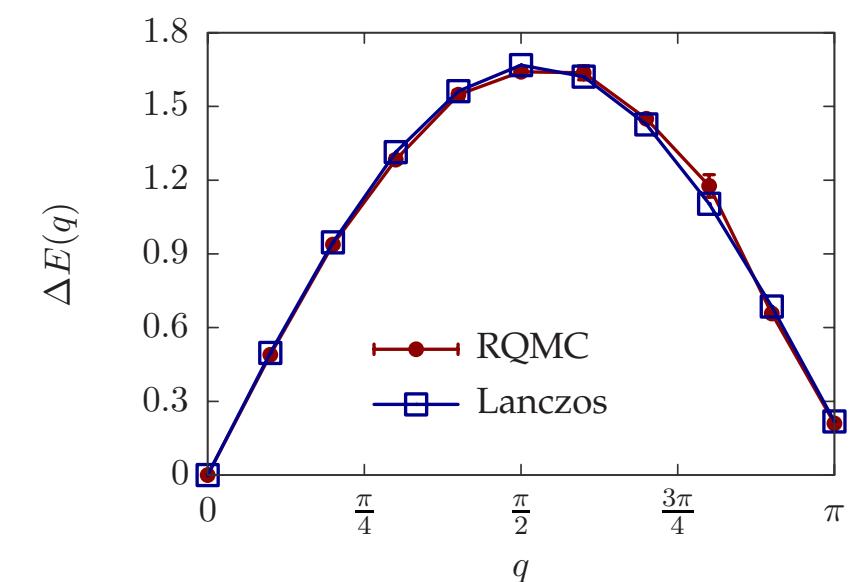
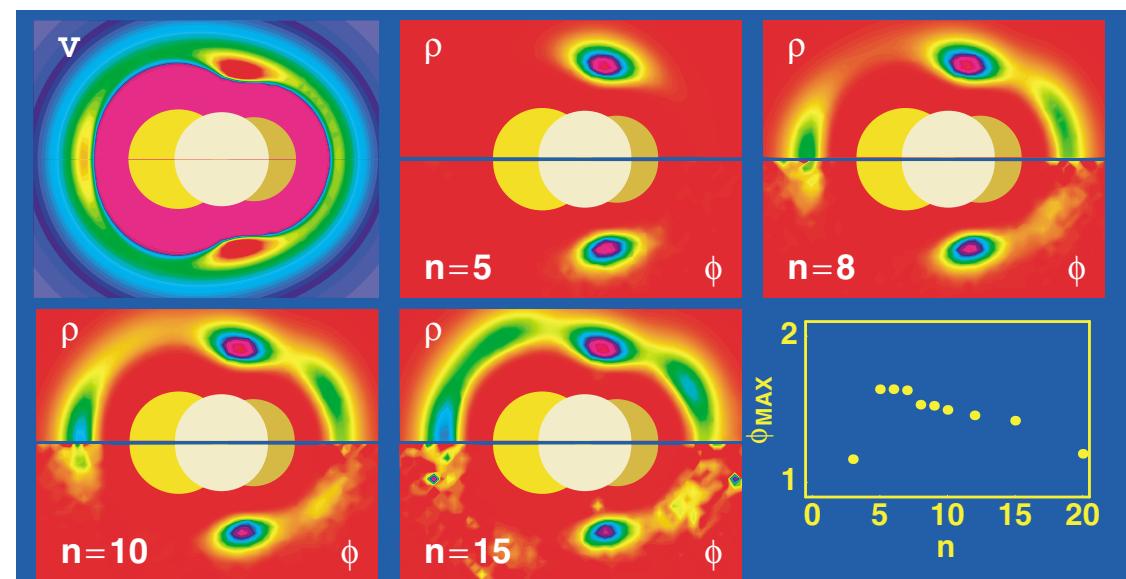
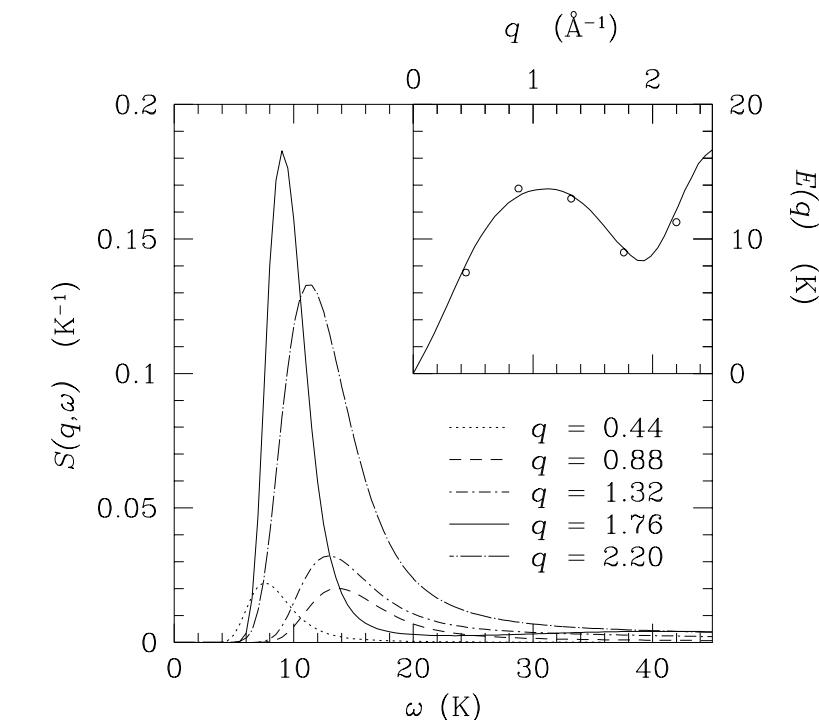
Giuseppe Carleo, Federico Becca, Saverio Moroni, and Stefano Baroni

RAPID COMMUNICATIONS

PHYSICAL REVIEW B 83, 060411(R) (2011)

Itinerant ferromagnetic phase of the Hubbard model

Giuseppe Carleo, Saverio Moroni, Federico Becca, and Stefano Baroni



back to higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = -\kappa_1 + \frac{1}{2}\kappa_2 + \cdots \frac{(-)^n}{n!}\kappa_n + \cdots$$

$$\begin{aligned} \kappa_1 &= \mu_1 \\ \kappa_2 &= \mu_2 - \mu_1^2 &= \bar{\mu}_2 \\ \kappa_3 &= \mu_3 - 3\mu_2\mu_1 + 2\mu_1^3 &= \bar{\mu}_3 \\ \kappa_4 &= \mu_4 - 4\mu_3^3\mu_1 - \cdots &= \bar{\mu}_4 - 3\bar{\mu}_2^2 \\ &&\dots \\ \kappa_n &= \mu_n - \sum_{m=1}^{n-1} \binom{n-1}{m} \kappa_{n-m} \mu_m \\ &= \bar{\mu}_n - \sum_{m=1}^{n-2} \binom{n-1}{m} \kappa_{n-m} \bar{\mu}_m \end{aligned}$$

$$\mu_n(T) = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$



back to higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = -\kappa_1 + \frac{1}{2}\kappa_2 + \cdots \frac{(-)^n}{n!}\kappa_n + \cdots$$

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$$\mu_n(T) = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$

$$\begin{aligned} E_0 &= -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle \\ &= \dot{\kappa}_1 - \frac{1}{2}\dot{\kappa}_2 + \cdots \frac{1}{n!}\dot{\kappa}_n + \cdots \\ &= \mathcal{E}_0 + E_0^{(1)} + E_0^{(2)} + \cdots E_0^{(n)} + \cdots \end{aligned}$$



back to higher-order perturbation theory

$$\langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = 1 - \mu_1 + \frac{1}{2}\mu_2 + \cdots \frac{(-)^n}{n!}\mu_n + \cdots$$

$$\log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle_{QM} = -\kappa_1 + \frac{1}{2}\kappa_2 + \cdots \frac{(-)^n}{n!}\kappa_n + \cdots$$

$$\begin{aligned} \kappa_1 &= \mu_1 \\ \kappa_2 &= \mu_2 - \mu_1^2 &= \bar{\mu}_2 \\ \kappa_3 &= \mu_3 - 3\mu_2\mu_1 + 2\mu_1^3 &= \bar{\mu}_3 \\ \kappa_4 &= \mu_4 - 4\mu_3^3\mu_1 - \cdots &= \bar{\mu}_4 - 3\bar{\mu}_2^2 \\ &&\dots \\ \kappa_n &= \mu_n - \sum_{m=1}^{n-1} \binom{n-1}{m} \kappa_{n-m} \mu_m \\ &= \bar{\mu}_n - \sum_{m=1}^{n-2} \binom{n-1}{m} \kappa_{n-m} \bar{\mu}_m \end{aligned}$$

$$\mu_n(T) = n! \int_0^T dt_n \int_0^{t_n} dt_{n-1} \cdots \int_0^{t_2} dt_1 \langle \Phi_0 | \widehat{\mathcal{W}}(t_n) \widehat{\mathcal{W}}(t_{n-1}) \cdots \widehat{\mathcal{W}}(t_1) | \Phi_0 \rangle_{QM}$$

$$\begin{aligned} E_0 &= -\frac{d}{dT} \log \langle \Phi_0 | e^{-\hat{H}T} | \Phi_0 \rangle \\ &= \dot{\kappa}_1 - \frac{1}{2}\dot{\kappa}_2 + \cdots \frac{1}{n!}\dot{\kappa}_n + \cdots \\ &= \mathcal{E}_0 + E_0^{(1)} + E_0^{(2)} + \cdots E_0^{(n)} + \cdots \end{aligned}$$

$$E_0^{(n)} = \frac{(-)^{n+1}}{n!} \dot{\kappa}_n$$



back to higher-order perturbation theory

$$E_0^{(n)} = \frac{(-)^{n+1}}{n!} \dot{\kappa}_n$$

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$$\dot{\kappa}_n \sim \dot{\mu}_n^\circ - \sum_{m=1}^{n-1} \binom{n-1}{m} (\dot{\kappa}_{n-m} \mu_m^\circ + \kappa_{n-m}^\circ \dot{\mu}_m^\circ)$$

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$$\dot{\kappa}_n \sim \dot{\mu}_n^\circ - \sum_{m=1}^{n-1} \binom{n-1}{m} (\dot{\kappa}_{n-m} \mu_m^\circ + \kappa_{n-m}^\circ \dot{\mu}_m^\circ)$$

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$$\mathcal{L}[\mu_n](z) = n! \sum_{k_1 \cdots k_{n-1}} \frac{\mathcal{W}_{0k_1} \mathcal{W}_{k_1 k_2} \cdots \mathcal{W}_{10}}{z^2(z + \mathcal{E}_{k_{n-1}}) \cdots (z + \mathcal{E}_{k_1})}$$



back to higher-order perturbation theory

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$$\mu_n^\circ \sim \text{zero-th order of } \frac{1}{(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_n)} = \frac{1}{z^k(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_{n-k})}$$

$$= \frac{1}{k!} \frac{d^k}{dz^k} \left(\frac{1}{(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_{n-k})} \right)_{z=0} = \frac{1}{k!} \left(\frac{\partial}{\partial \mathcal{E}_1} + \frac{\partial}{\partial \mathcal{E}_{n-k}} \right)^k \frac{1}{\mathcal{E}_1 \cdots \mathcal{E}_{n-k}}$$



back to higher-order perturbation theory

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$$\dot{\kappa}_n \sim \dot{\mu}_n^\circ - \sum_{m=1}^{n-1} \binom{n-1}{m} (\dot{\kappa}_{n-m} \mu_m^\circ + \kappa_{n-m}^\circ \dot{\mu}_m^\circ)$$

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$$\mu_n^\circ \sim \text{zero-th order of } \frac{1}{(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_n)} = \frac{1}{z^k(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_{n-k})} \quad \text{and similarly for } \dot{\mu}_0^n$$

$$= \frac{1}{k!} \frac{d^k}{dz^k} \left(\frac{1}{(z + \mathcal{E}_1) \cdots (z + \mathcal{E}_{n-k})} \right)_{z=0} = \frac{1}{k!} \left(\frac{\partial}{\partial \mathcal{E}_1} + \frac{\partial}{\partial \mathcal{E}_{n-k}} \right)^k \frac{1}{\mathcal{E}_1 \cdots \mathcal{E}_{n-k}}$$



an example: third-order perturbation theory

$$E_0 = \dot{\kappa}_1 - \frac{1}{2} \dot{\kappa}_2 + \frac{1}{6} \dot{\kappa}_3 + \dots$$
$$E_0^{(1)} \quad E_0^{(2)} \quad E_0^{(3)}$$



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$$E_0 = \frac{\dot{\kappa}_1}{E_0^{(1)}} - \frac{1}{2} \frac{\dot{\kappa}_2}{E_0^{(2)}} + \frac{1}{6} \frac{\dot{\kappa}_3}{E_0^{(3)}} + \dots$$

$$\begin{aligned}\kappa_3 &= \mu_3 - 3\mu_2\mu_1 + 2\mu_1^2 & \dot{\kappa}_3 &= \dot{\mu}_3 - 3\dot{\mu}_2\mu_1 + 3\mu_2\dot{\mu}_1 + 6\dot{\mu}_1\mu_1^2 \\ & & & \sim \ddot{\mu}_3^\circ + 3\ddot{\mu}_2^\circ\dot{\mu}_1\end{aligned}$$



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$$\mathcal{L}[\dot{\mu}_3](z) = \frac{6}{z} \sum_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l + z)(\mathcal{E}_m + z)}$$



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an example: third-order perturbation theory

$$E_0 = \frac{\dot{\kappa}_1}{E_0^{(1)}} - \frac{1}{2} \frac{\dot{\kappa}_2}{E_0^{(2)}} + \frac{1}{6} \frac{\dot{\kappa}_3}{E_0^{(3)}} + \dots$$

$$\kappa_3 = \mu_3 - 3\mu_2\mu_1 + 2\mu_1^2$$

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$$\begin{aligned}\mathcal{L}[\dot{\mu}_3](z) &= \frac{6}{z} \sum_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} \\ &= \frac{6}{z} \sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} + \frac{12\mathcal{W}_{00}}{z} \sum'_l \frac{\mathcal{W}_{l0}^2}{(\mathcal{E}_l+z)} \\ &= \frac{1}{z} \color{blue}{6} \left[\sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{\mathcal{E}_l\mathcal{E}_m} - 2\mathcal{W}_{00} \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2} \right] + \mathcal{O}(z^{-2})\end{aligned}$$

$$\mathcal{L}[\mu_2](z) = \frac{2}{z^2} \sum_l \frac{\mathcal{W}_{0l}^2}{\mathcal{E}_l + z}$$



an example: third-order perturbation theory

$$E_0 = \frac{\dot{\kappa}_1}{E_0^{(1)}} - \frac{1}{2} \frac{\dot{\kappa}_2}{E_0^{(2)}} + \frac{1}{6} \frac{\dot{\kappa}_3}{E_0^{(3)}} + \dots$$

$$\kappa_3 = \mu_3 - 3\mu_2\mu_1 + 2\mu_1^2$$

$$\begin{aligned}\dot{\kappa}_3 &= \dot{\mu}_3 - 3\dot{\mu}_2\mu_1 + 3\mu_2\dot{\mu}_1 + 6\dot{\mu}_1\mu_1^2 \\ &\sim \color{blue}{\dot{\mu}_3^\circ} + \color{red}{3\mu_2^\circ\dot{\mu}_1}\end{aligned}$$

$$\begin{aligned}\mathcal{L}[\dot{\mu}_3](z) &= \frac{6}{z} \sum_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} \\ &= \frac{6}{z} \sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} + \frac{12\mathcal{W}_{00}}{z} \sum'_l \frac{\mathcal{W}_{l0}^2}{(\mathcal{E}_l+z)} \\ &= \frac{1}{z} \color{blue}{6} \left[\sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{\mathcal{E}_l\mathcal{E}_m} - 2\mathcal{W}_{00} \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2} \right] + \mathcal{O}(z^{-2})\end{aligned}$$

$$\begin{aligned}\mathcal{L}[\mu_2](z) &= \frac{2}{z^2} \sum_l \frac{\mathcal{W}_{0l}^2}{\mathcal{E}_l+z} \\ &= \frac{1}{z} \left[-2 \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2} \right] + \mathcal{O}(z^{-2})\end{aligned}$$



an example: third-order perturbation theory

$$E_0 = \frac{\dot{\kappa}_1}{E_0^{(1)}} - \frac{1}{2} \frac{\dot{\kappa}_2}{E_0^{(2)}} + \frac{1}{6} \frac{\dot{\kappa}_3}{E_0^{(3)}} + \dots$$

$$\kappa_3 = \mu_3 - 3\mu_2\mu_1 + 2\mu_1^2$$

$$\begin{aligned}\dot{\kappa}_3 &= \dot{\mu}_3 - 3\dot{\mu}_2\mu_1 + 3\mu_2\dot{\mu}_1 + 6\dot{\mu}_1\mu_1^2 \\ &\sim \dot{\mu}_3^\circ + 3\mu_2^\circ\dot{\mu}_1\end{aligned}$$

$$\begin{aligned}\mathcal{L}[\dot{\mu}_3](z) &= \frac{6}{z} \sum_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} \\ &= \frac{6}{z} \sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{(\mathcal{E}_l+z)(\mathcal{E}_m+z)} + \frac{12\mathcal{W}_{00}}{z} \sum'_l \frac{\mathcal{W}_{l0}^2}{(\mathcal{E}_l+z)} \\ &= \frac{1}{z} \textcolor{blue}{6} \left[\sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{\mathcal{E}_l\mathcal{E}_m} - 2\mathcal{W}_{00} \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2} \right] + \mathcal{O}(z^{-2})\end{aligned}$$

$$E_0^{(3)} = \sum'_{lm} \frac{\mathcal{W}_{0l}\mathcal{W}_{lm}\mathcal{W}_{m0}}{\mathcal{E}_l\mathcal{E}_m} - \mathcal{W}_{00} \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2}$$

$$\begin{aligned}\mathcal{L}[\mu_2](z) &= \frac{2}{z^2} \sum_l \frac{\mathcal{W}_{0l}^2}{\mathcal{E}_l+z} \\ &= \frac{1}{z} \left[-2 \sum'_l \frac{\mathcal{W}_{l0}^2}{\mathcal{E}_l^2} \right] + \mathcal{O}(z^{-2})\end{aligned}$$





A man with light brown hair and a beard is lying on his back on a red and black patterned rug. He is wearing a dark green zip-up hoodie. A small white toy airplane is positioned near his head. The background is a dark, textured surface.

keep dreaming and flying
high for another 60 years,
dude!

thanks



thanks

- to Saverio for being such a nice dude, a life-long friend, and an esteemed colleague



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- to Giovanni, for prodding me to (almost) carry out the algebra that I started 25 years ago, but never dared to complete



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- to Giovanni, for prodding me to (almost) carry out the algebra that I started 25 years ago, but never dared to complete
- to all of you, for bearing with me today

these slides available soon at <https://talks.baroni.me>

