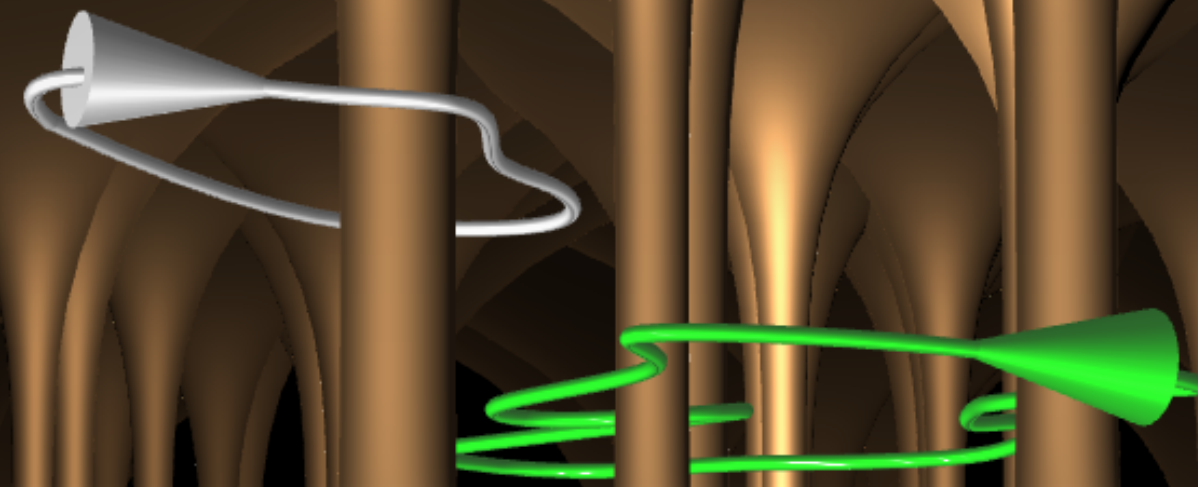




topological features of a quantum soup: electrons in broth, forty years on

Stefano Baroni
Scuola Internazionale Superiore di Studi Avanzati
Trieste — Italy



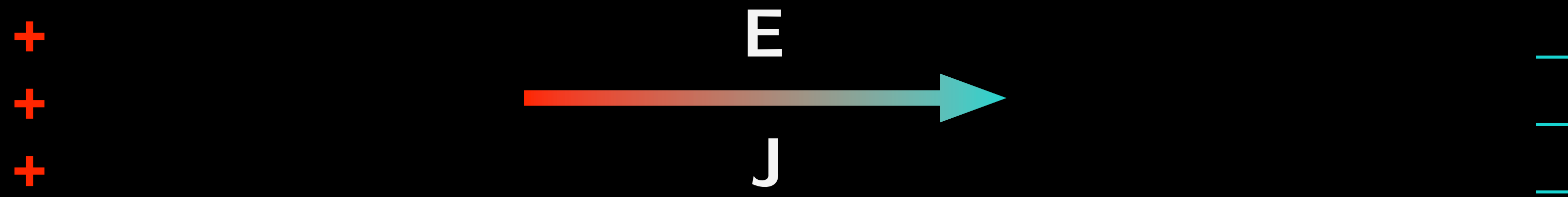
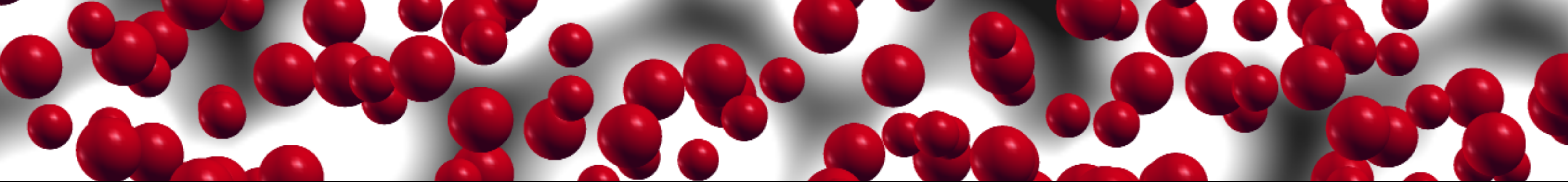
ionic transport

nothing flows as simply as it appears

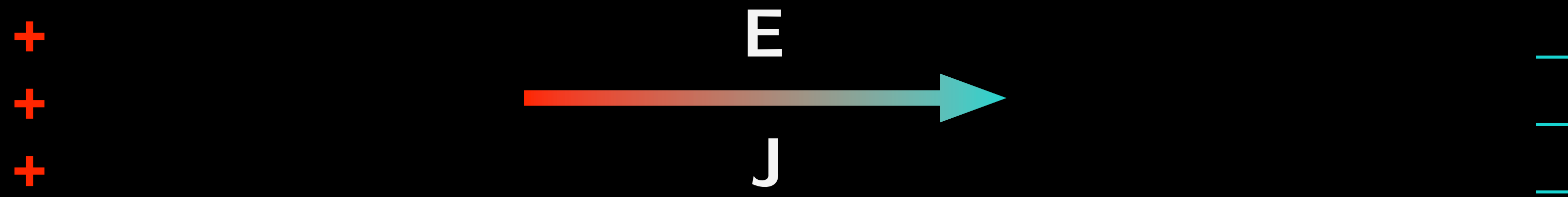
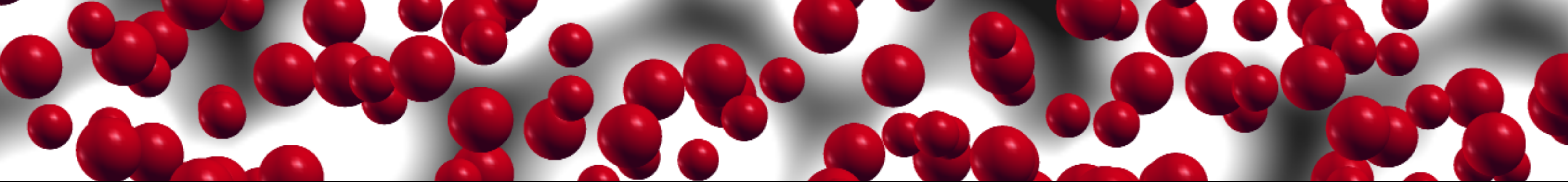


a prequel



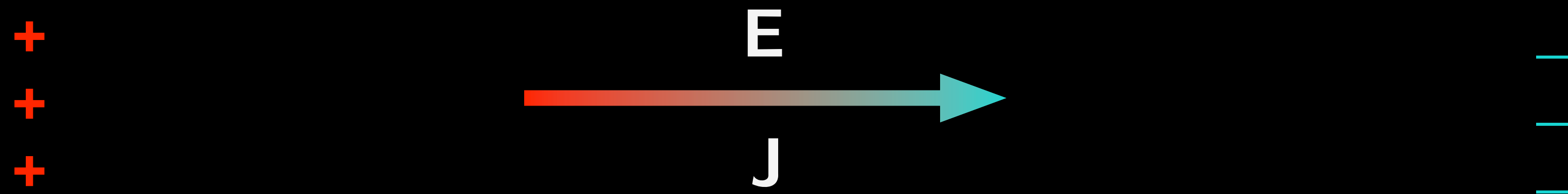
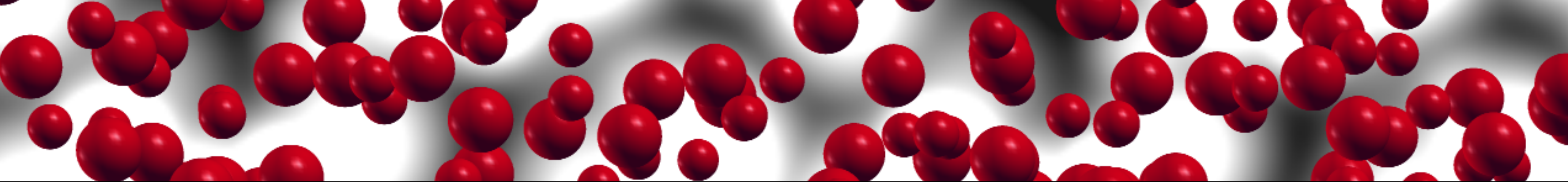


$$\mathbf{J} = \sigma \mathbf{E}$$



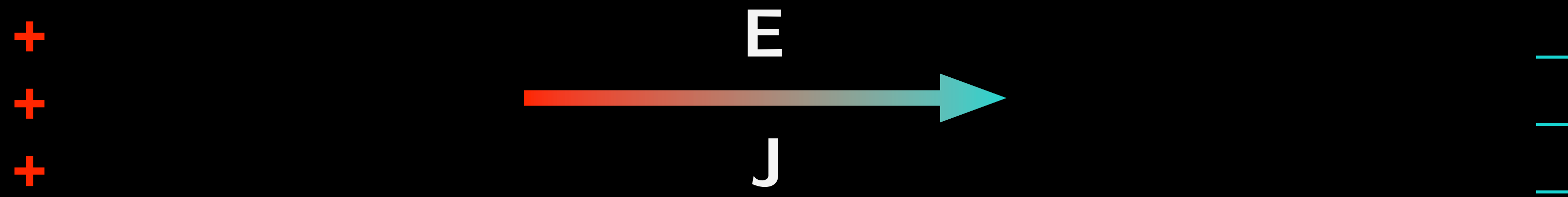
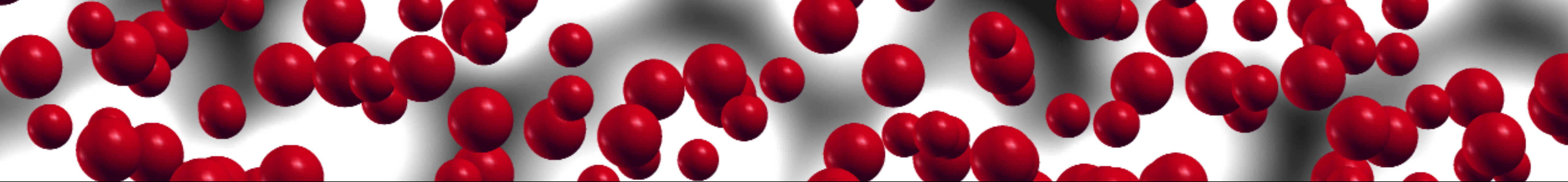
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$$\mathbf{J} = \frac{1}{\Omega} \dot{\boldsymbol{\mu}}$$



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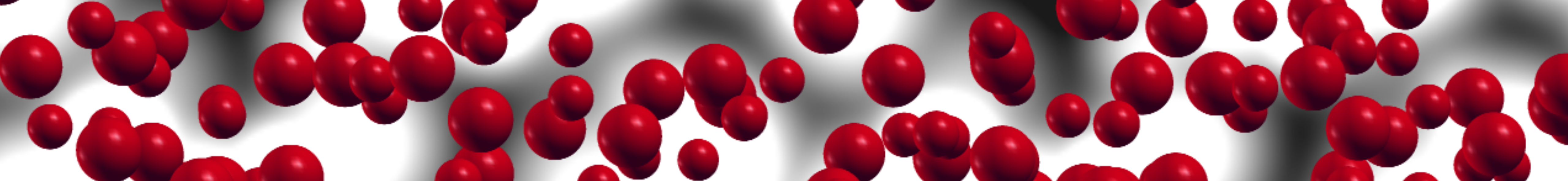
$$\begin{aligned} \mathbf{J} &= \frac{1}{\Omega} \dot{\boldsymbol{\mu}} \\ &= \frac{1}{\Omega} \sum_i \mathbf{z}_i^* \cdot \mathbf{v}_i \end{aligned}$$



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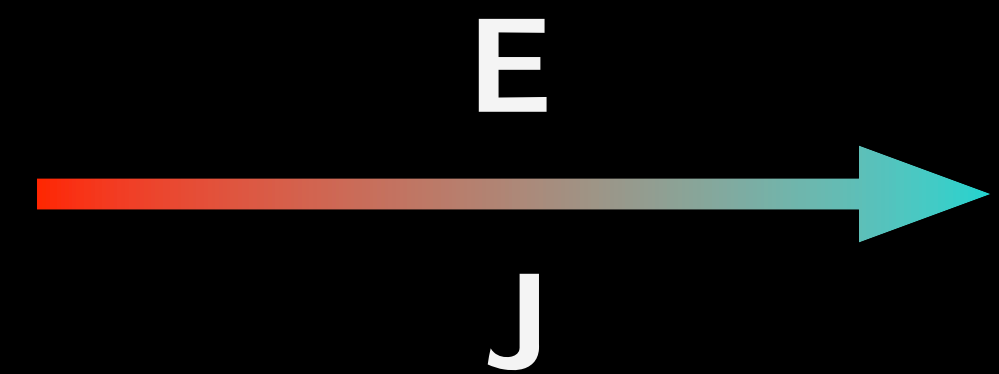
$$z_{i\alpha\beta}^* = \frac{\partial \mu_\alpha}{\partial u_{i\beta}}$$



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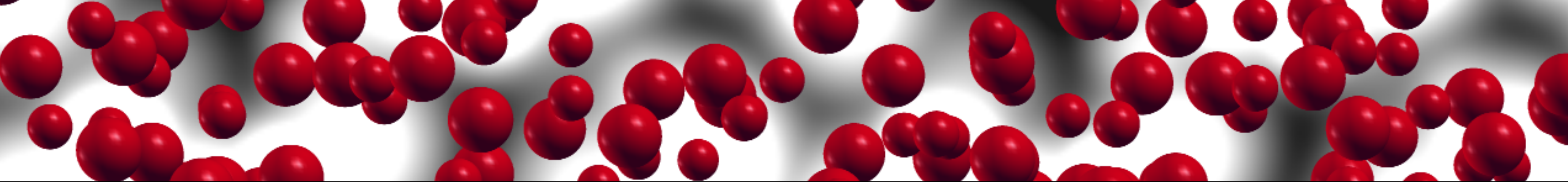
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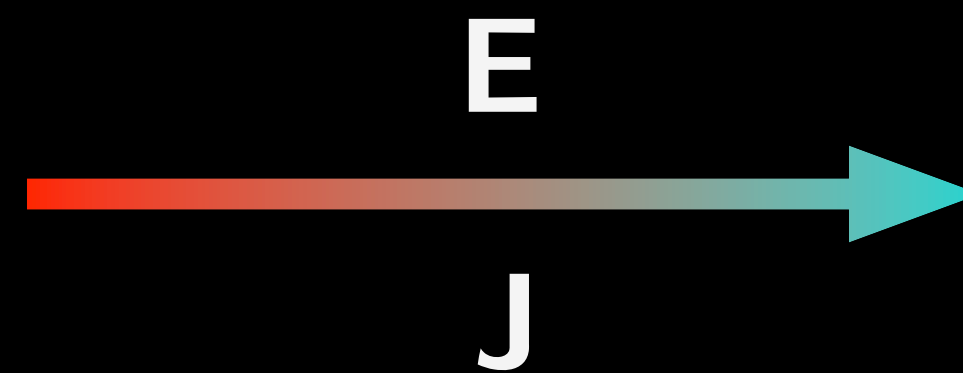
$$\sigma = \frac{\Omega}{3k_B T} \int_0^\infty \langle \mathbf{J}(t) \cdot \mathbf{J}(0) \rangle dt$$



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$$= \frac{\Omega}{3k_B T} \lim_{t \rightarrow \infty} \frac{1}{2t} \left\langle \left| \int_0^t \mathbf{J}(t') dt' \right|^2 \right\rangle$$

keep in mind!



a conundrum in transport theory

PRL **107**, 185901 (2011)

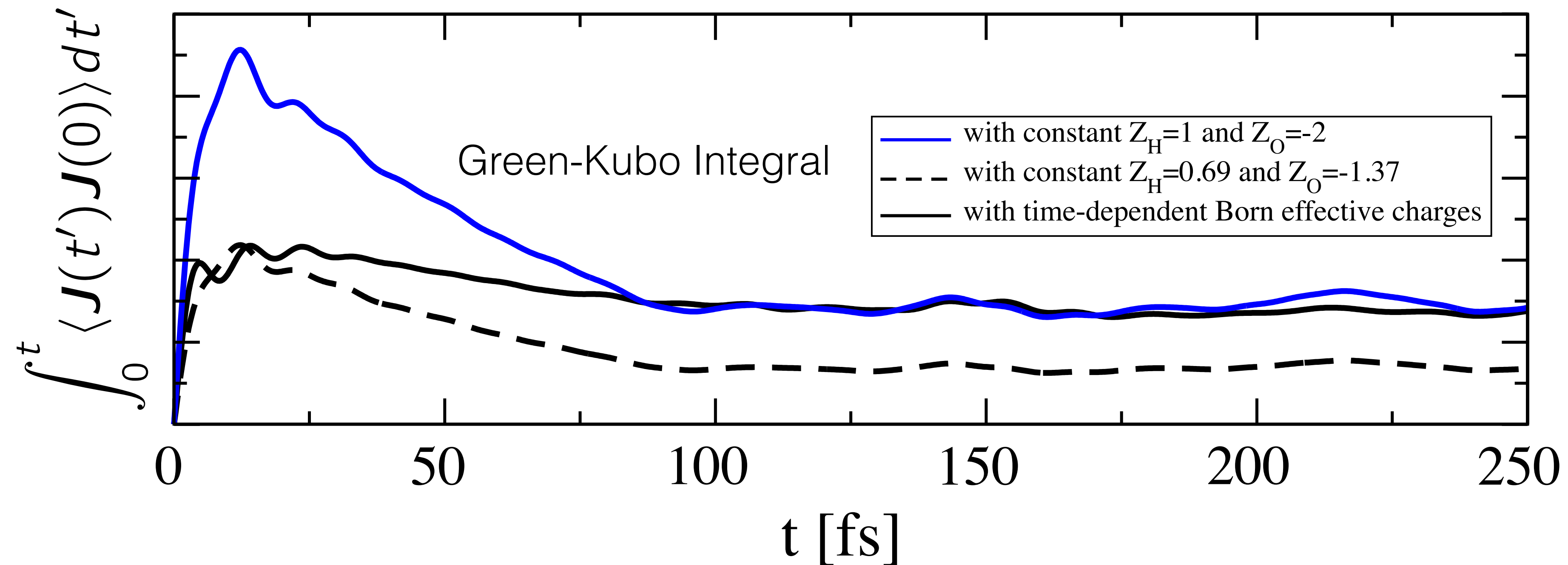
PHYSICAL REVIEW LETTERS

week ending
28 OCTOBER 2011

Dynamical Screening and Ionic Conductivity in Water from *Ab Initio* Simulations

Martin French,¹ Sebastien Hamel,² and Ronald Redmer¹

partially dissociated water



a conundrum in transport theory

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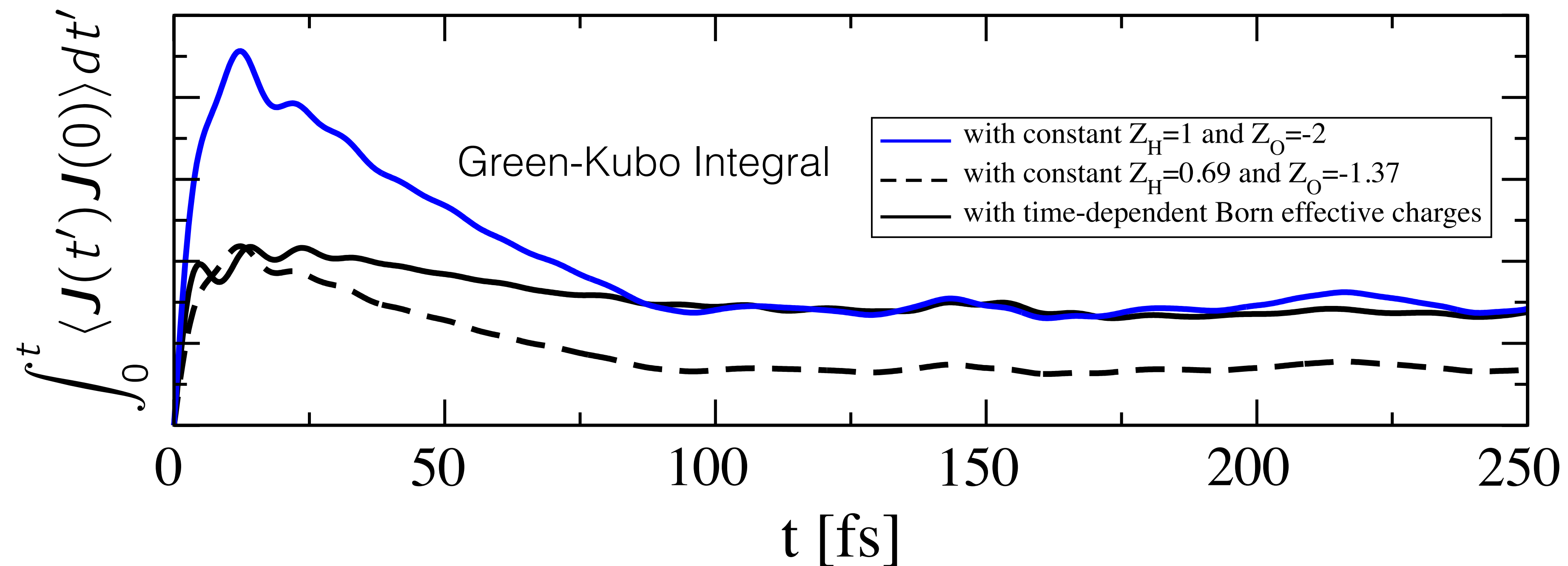
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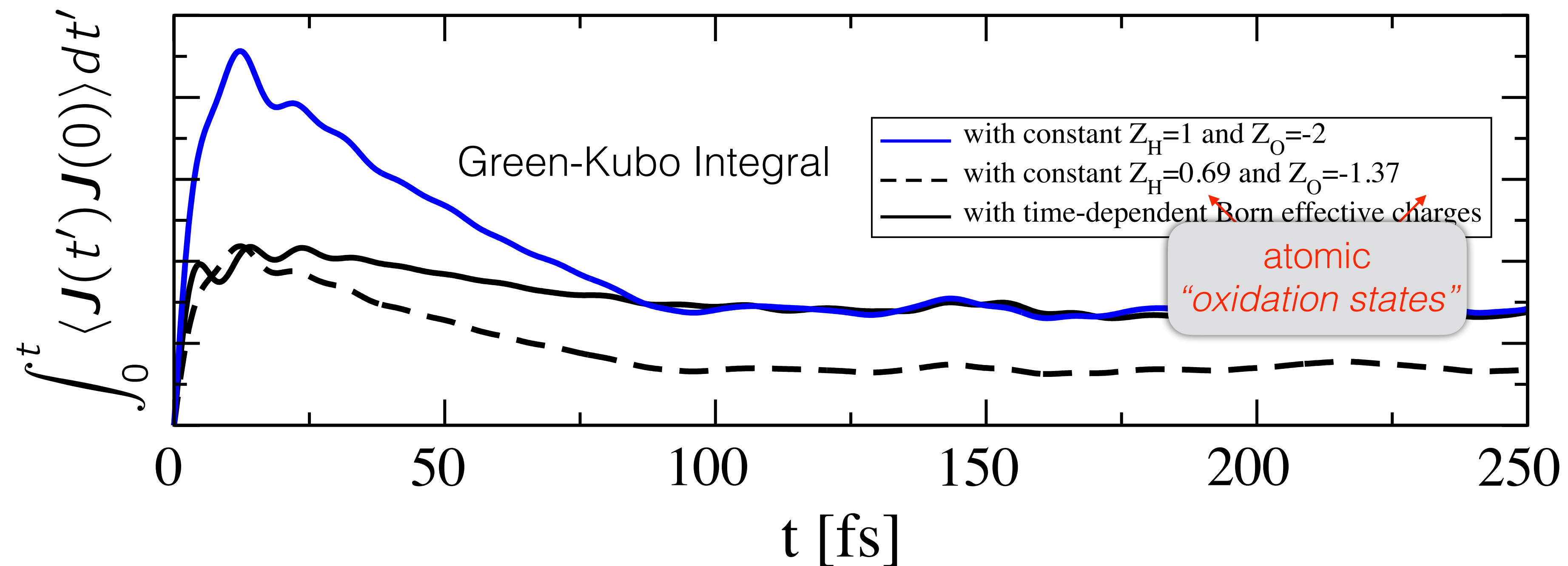
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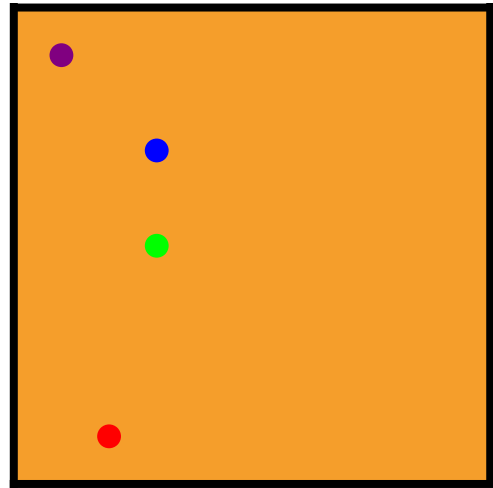


how come?

and what are oxidation states, in the first place?

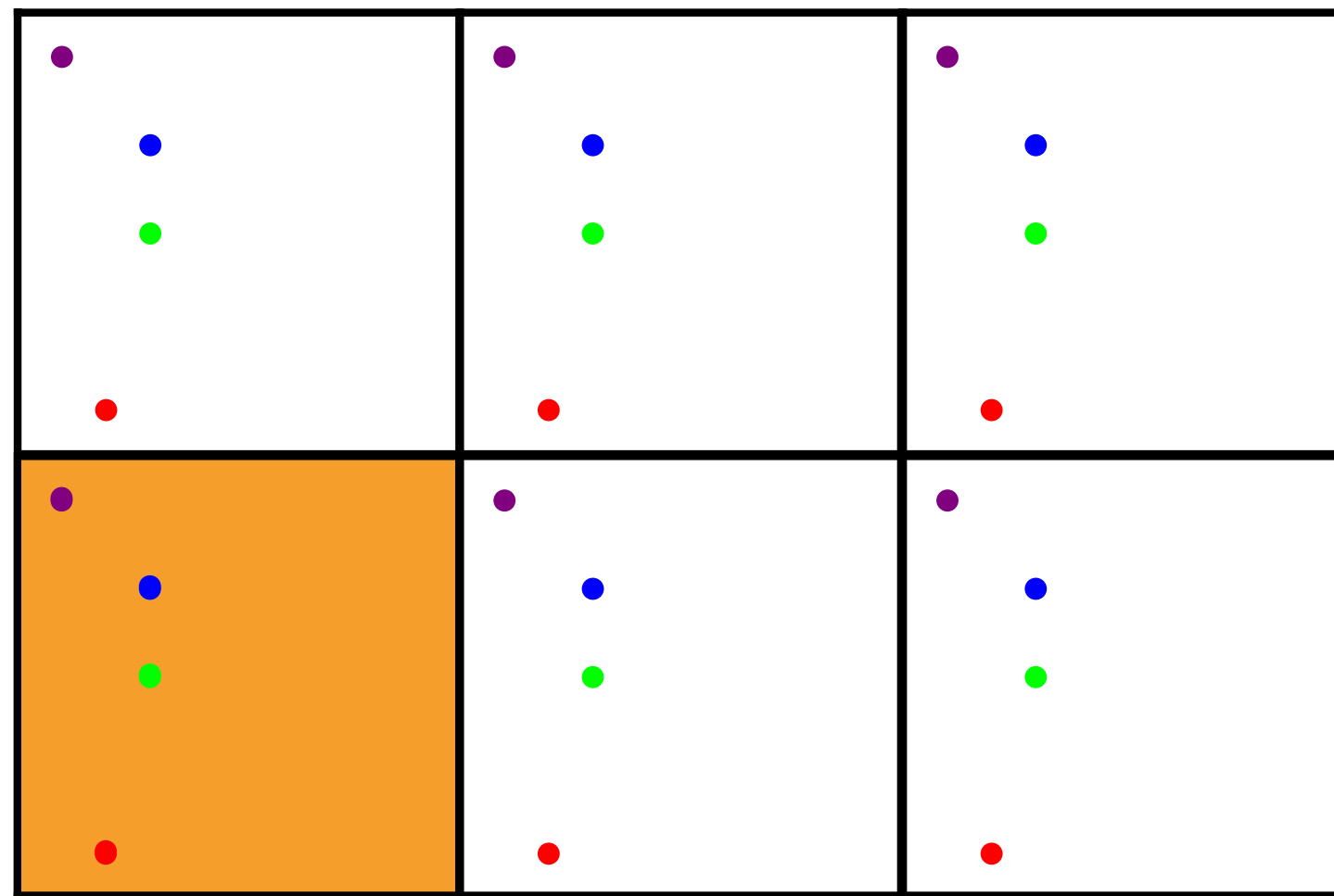
what are oxidation states, in the first place?

$$[0, L]^{3N}$$



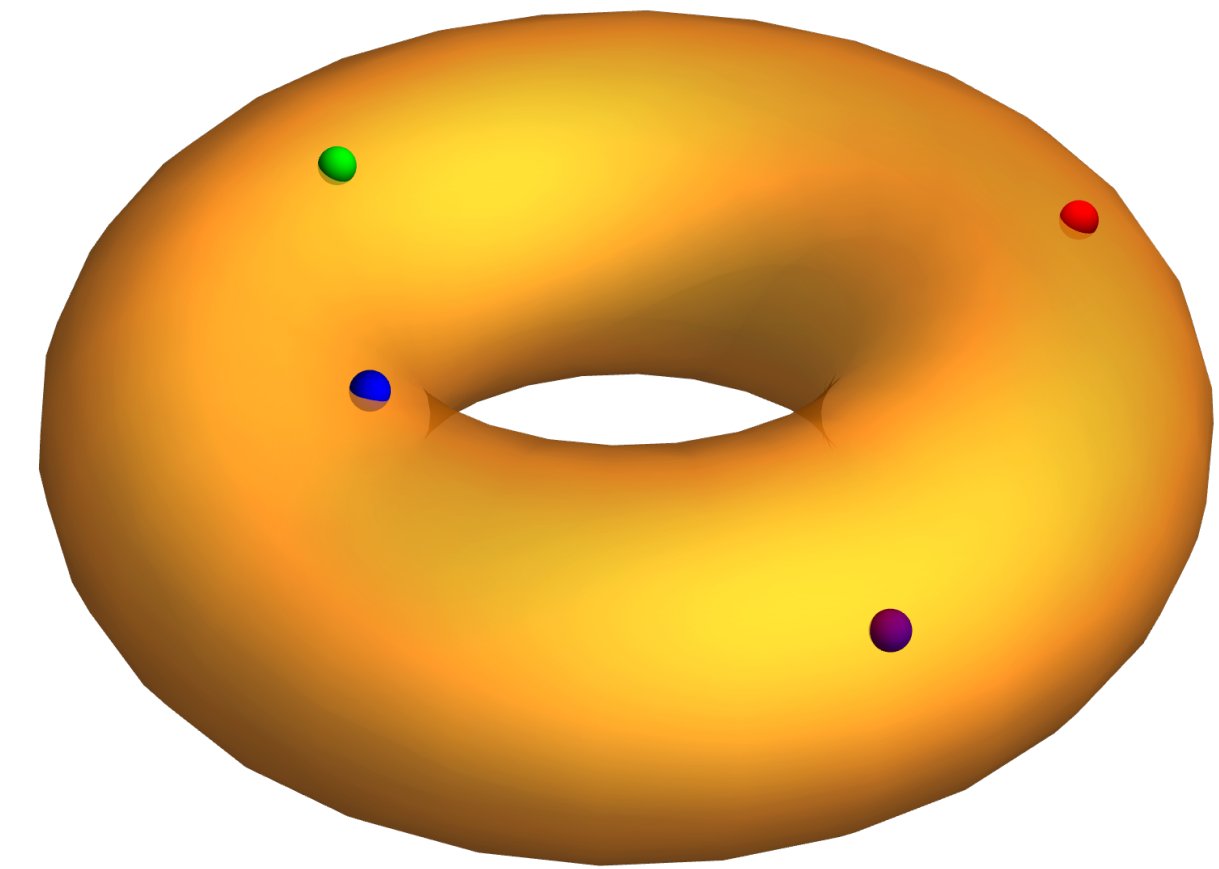
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PBC
→

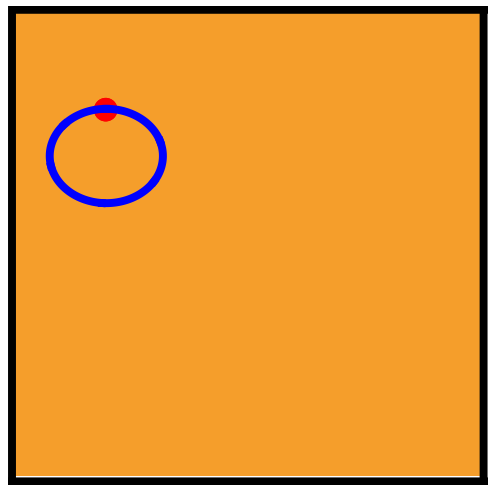
$$\mathbb{T}^{3N}$$



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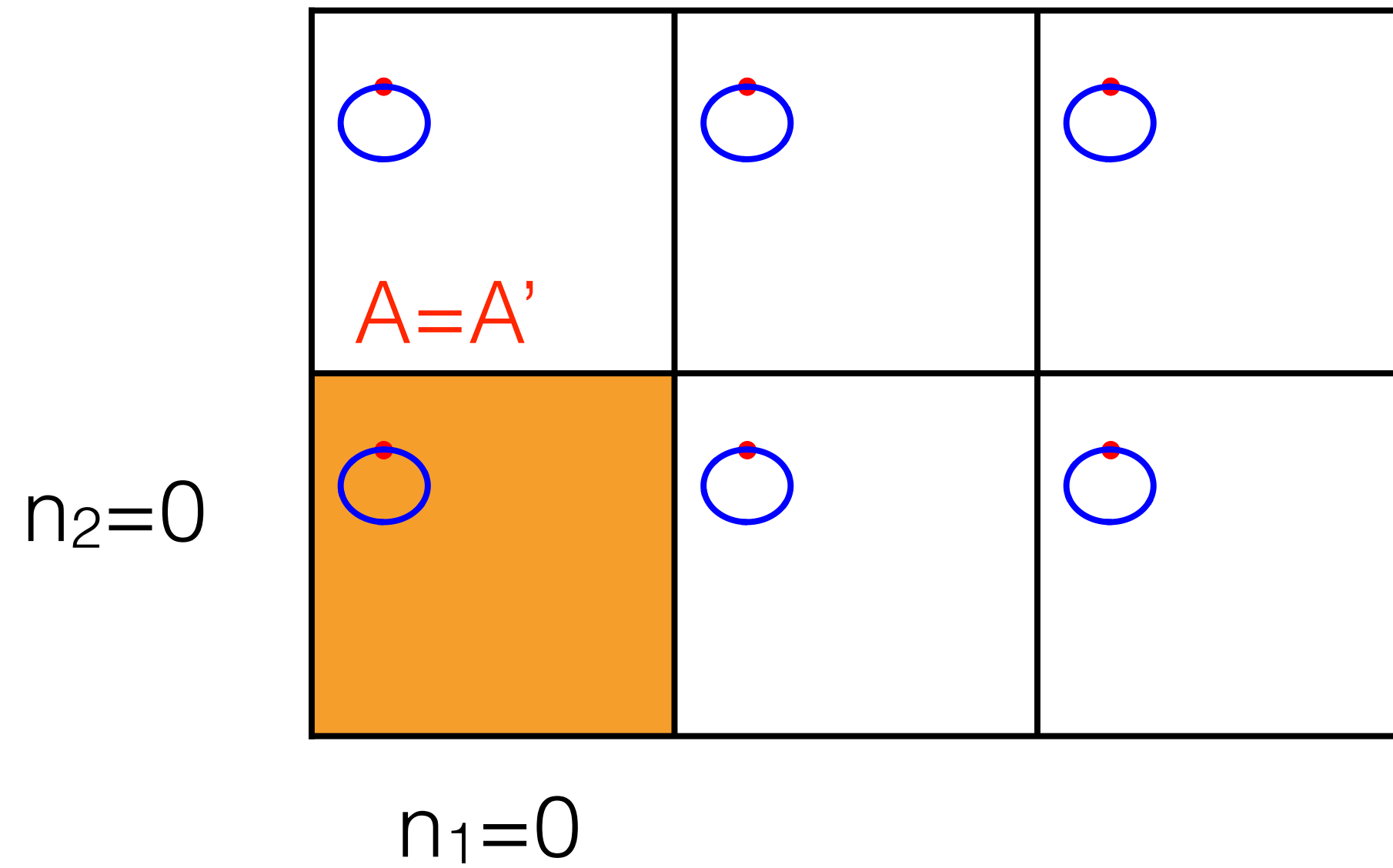
$$A=A'$$



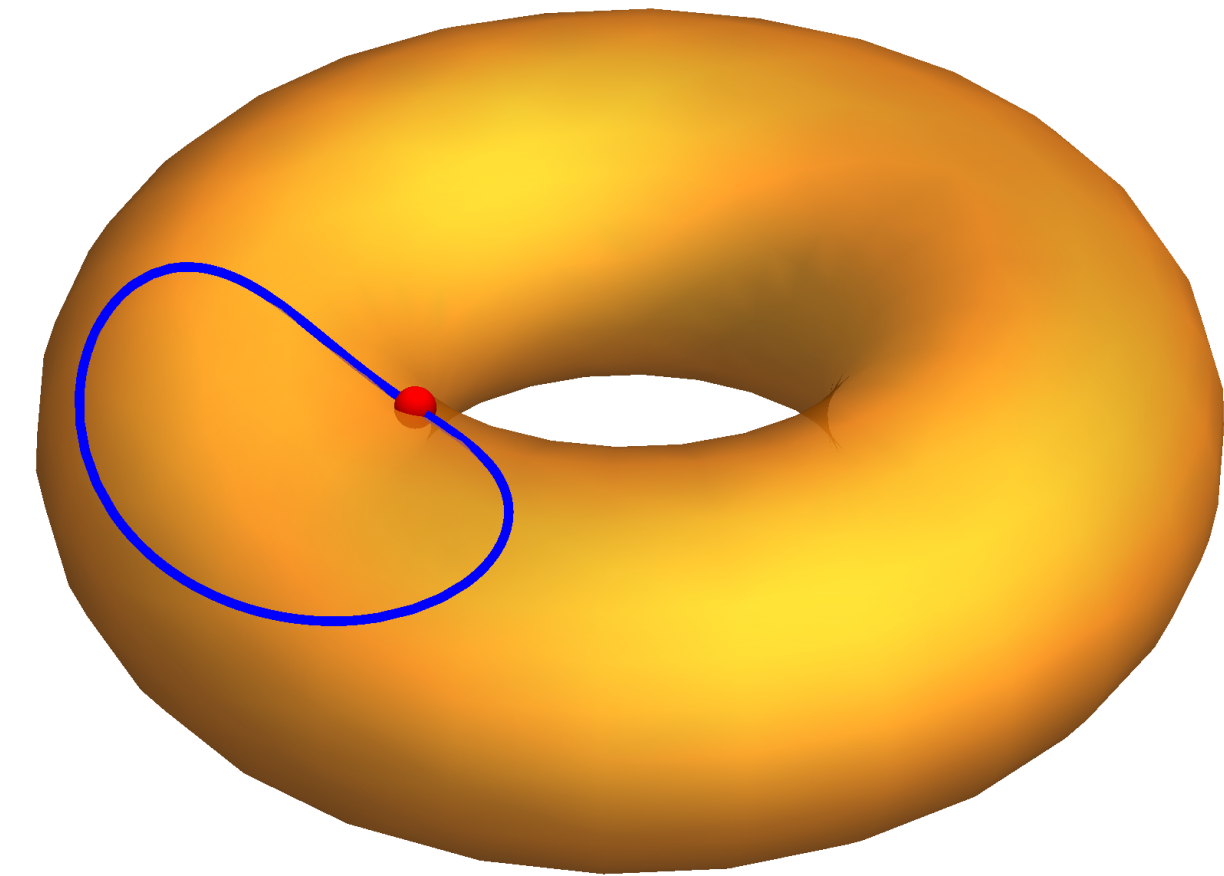
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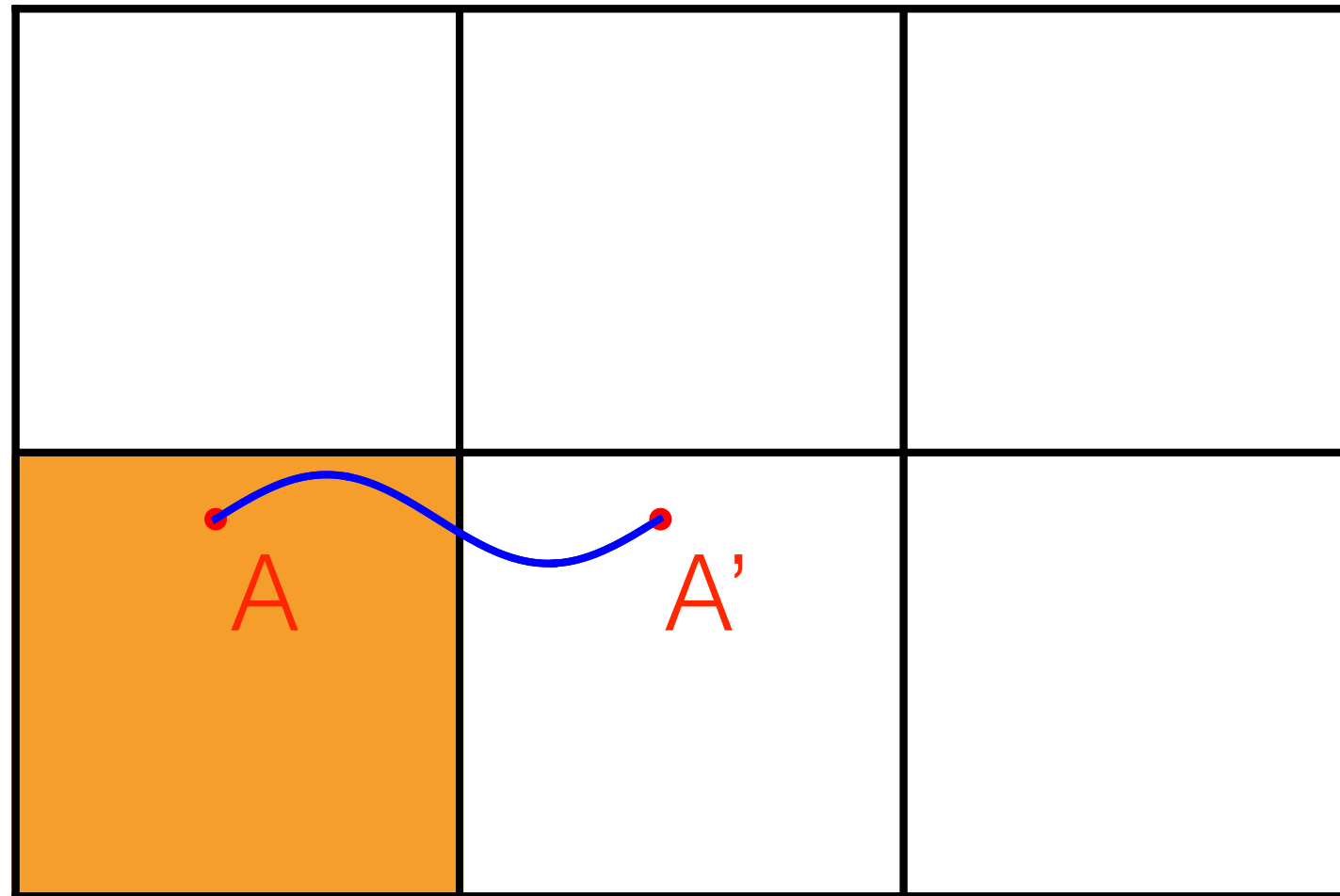


PBC
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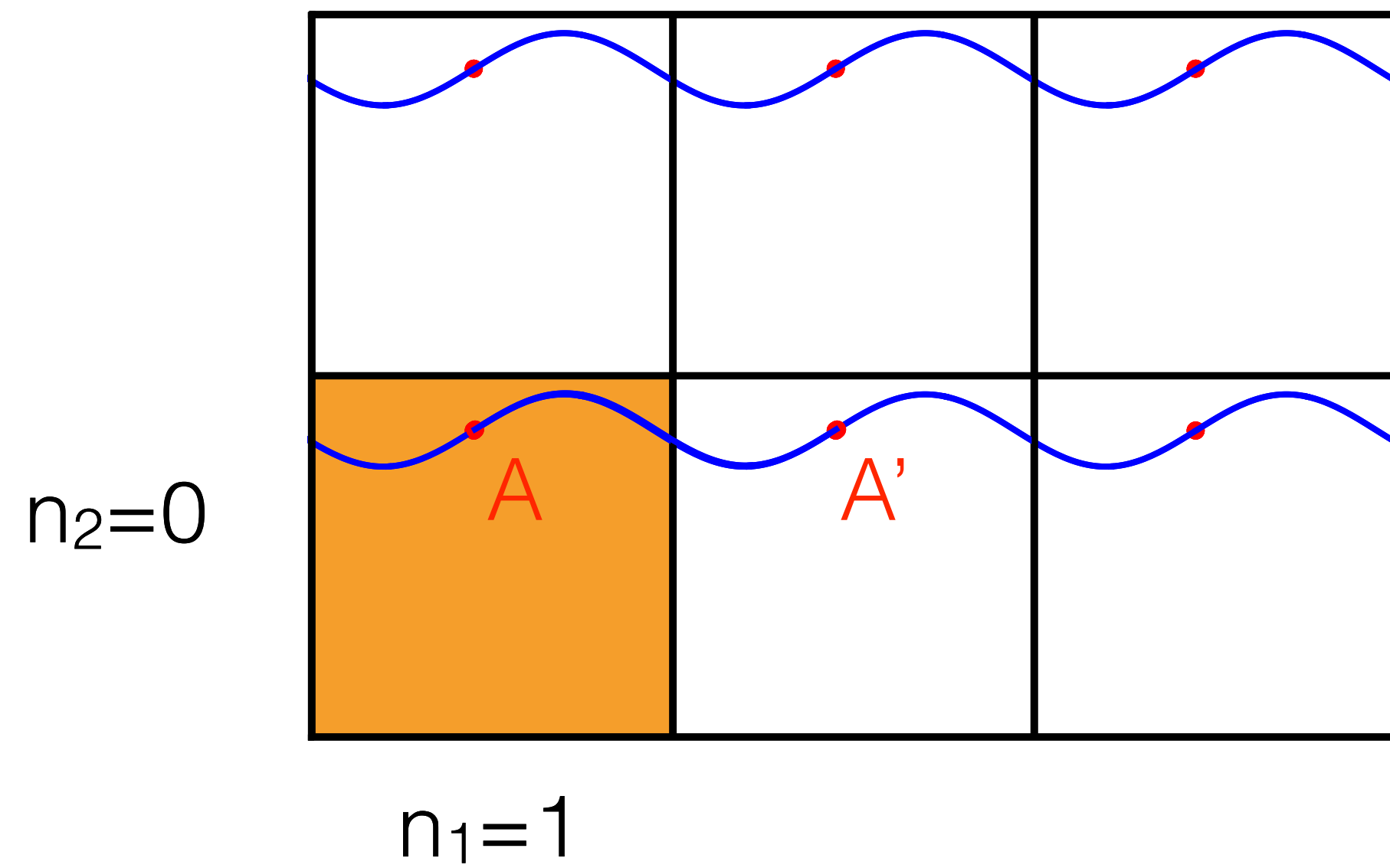
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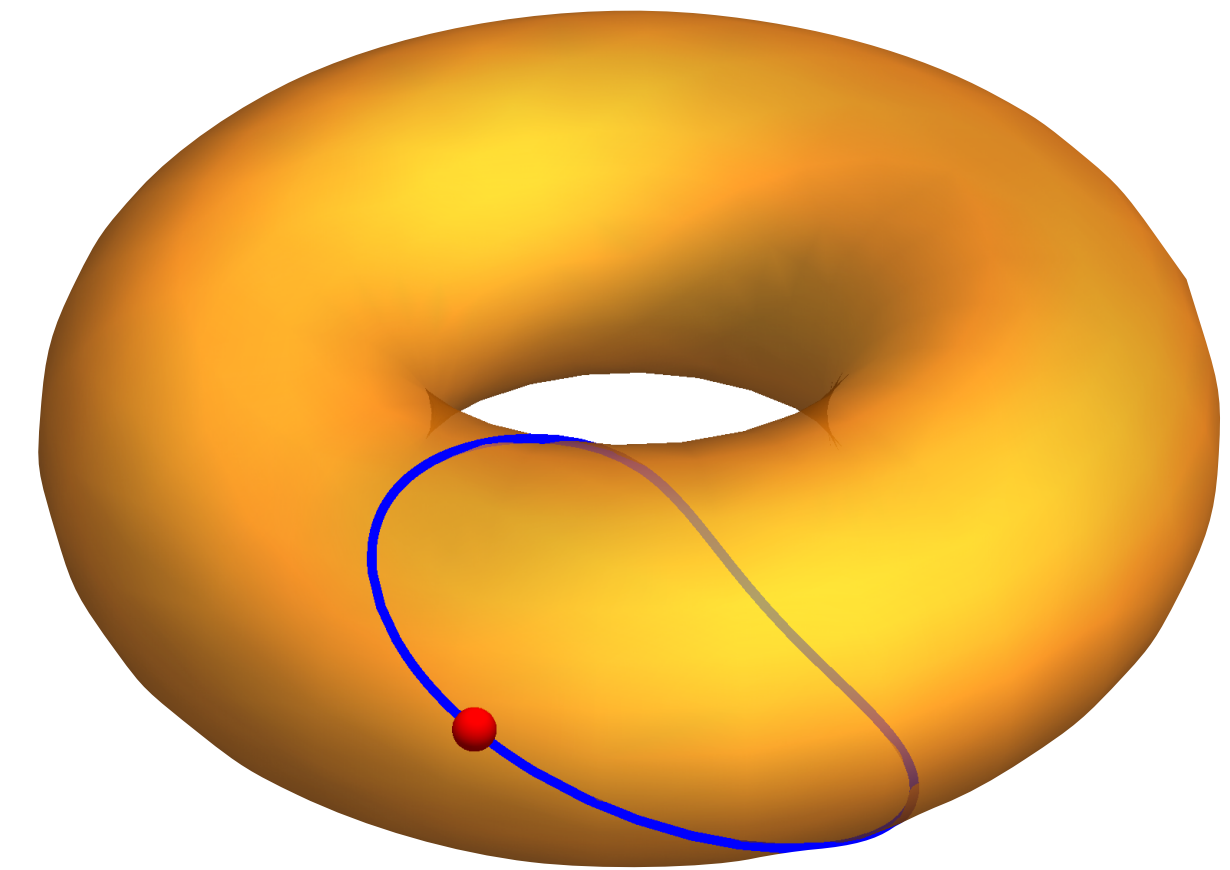
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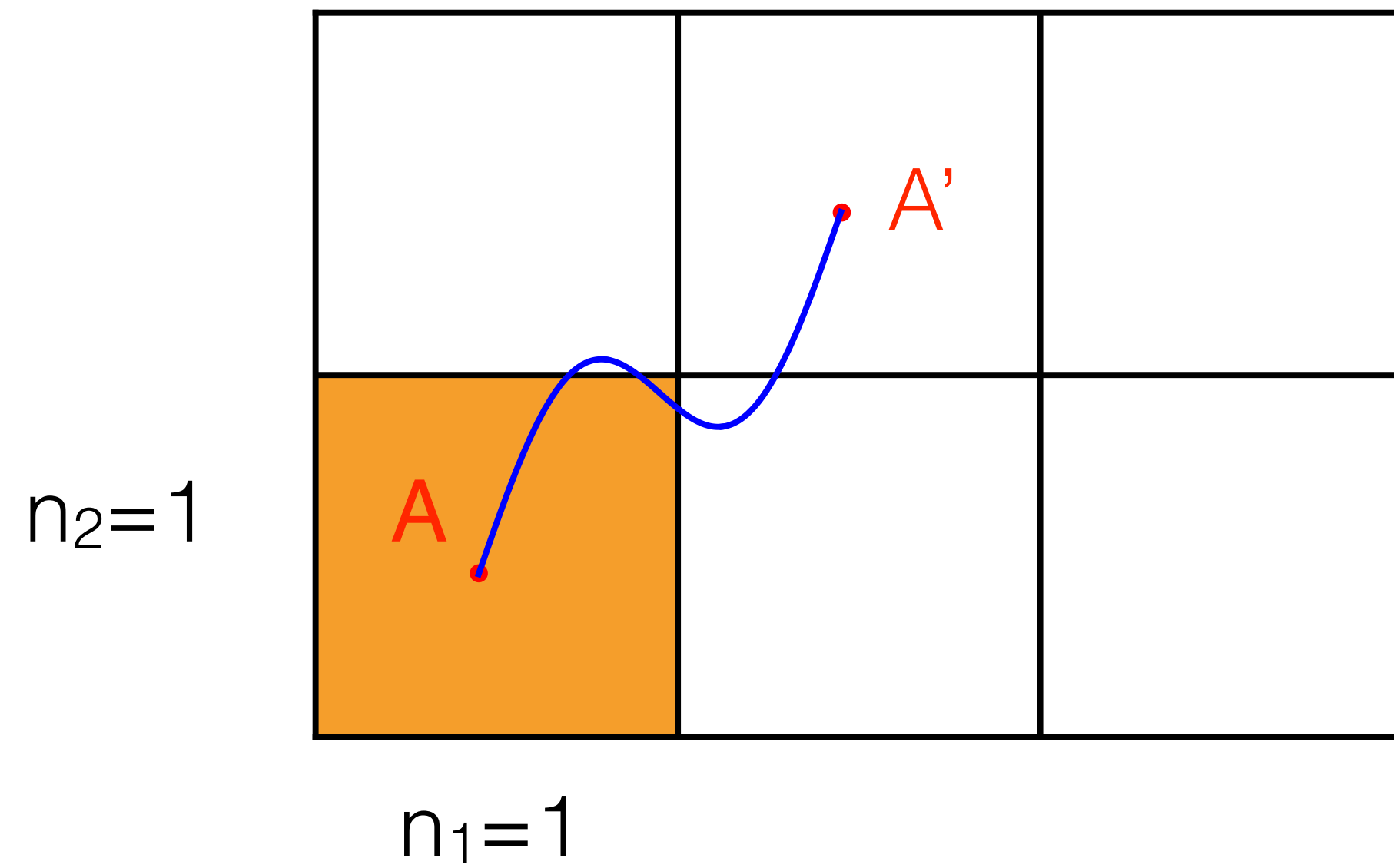


PBC



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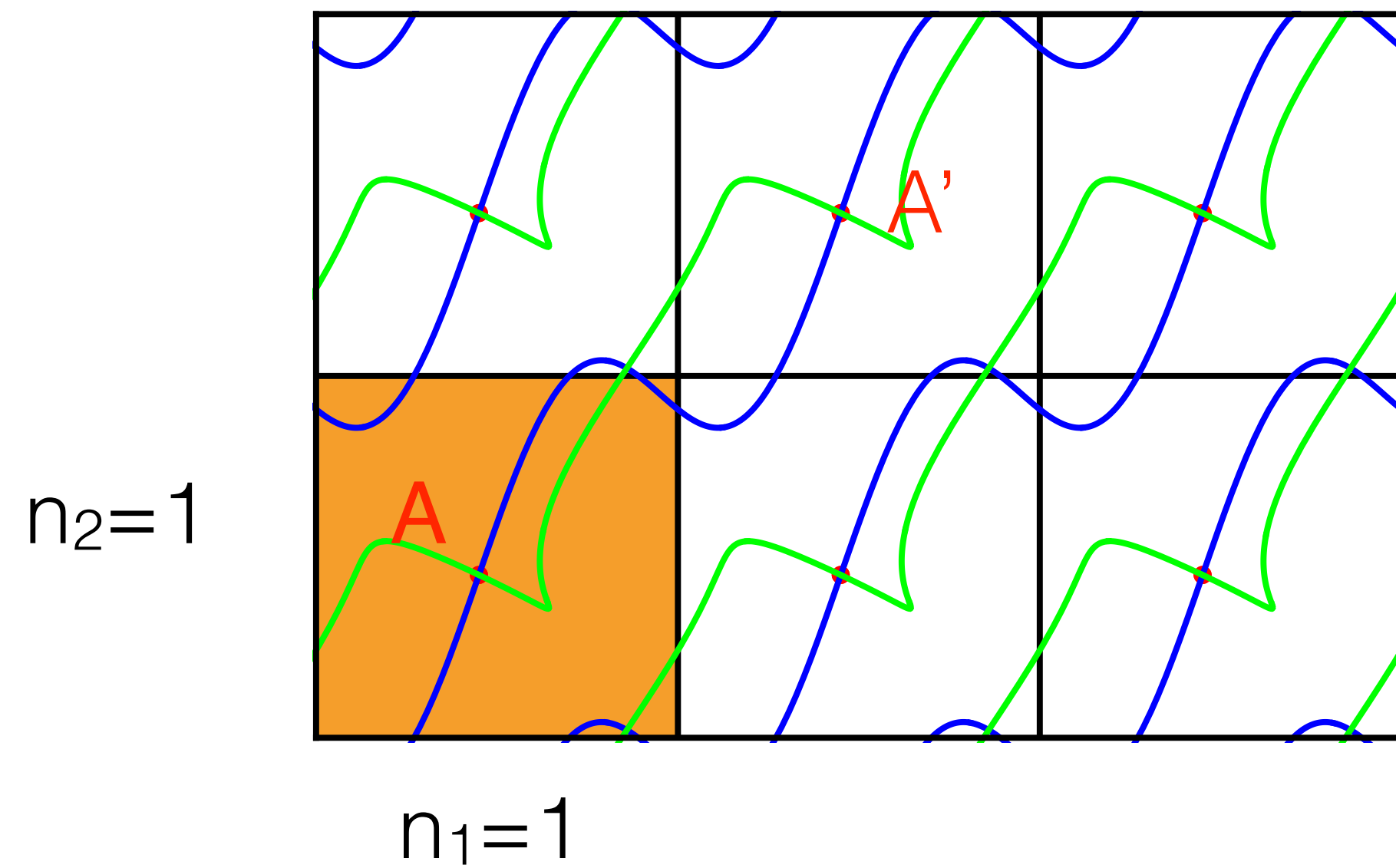
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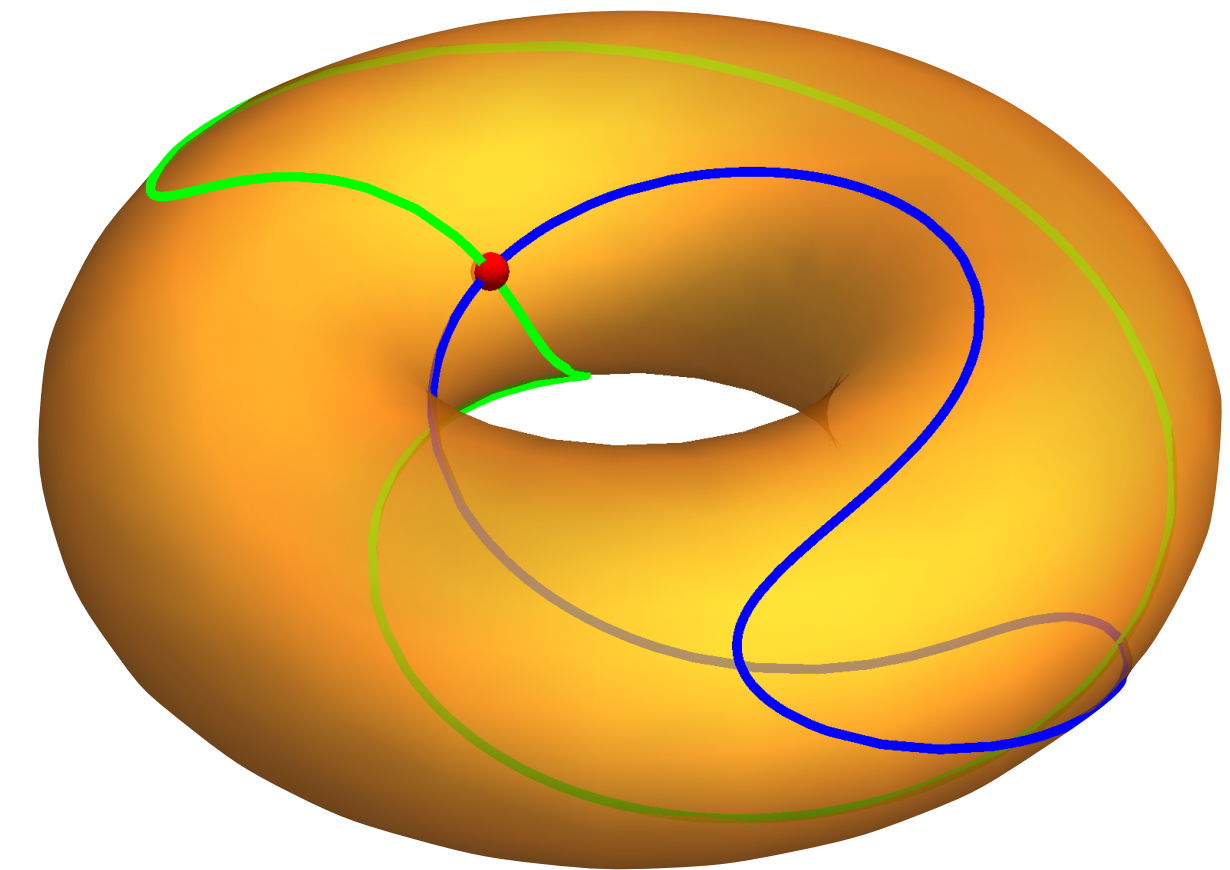
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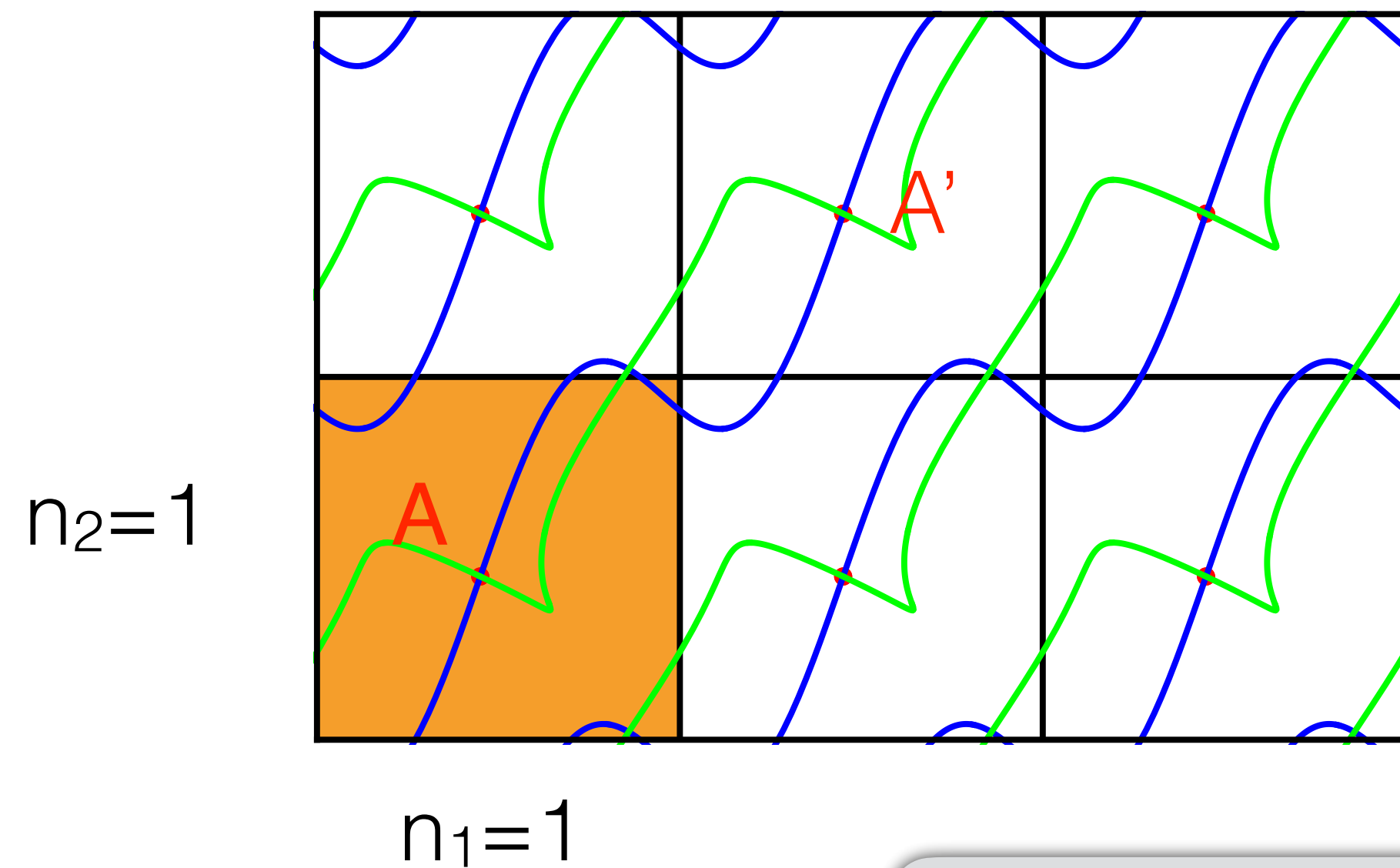
PBC
→



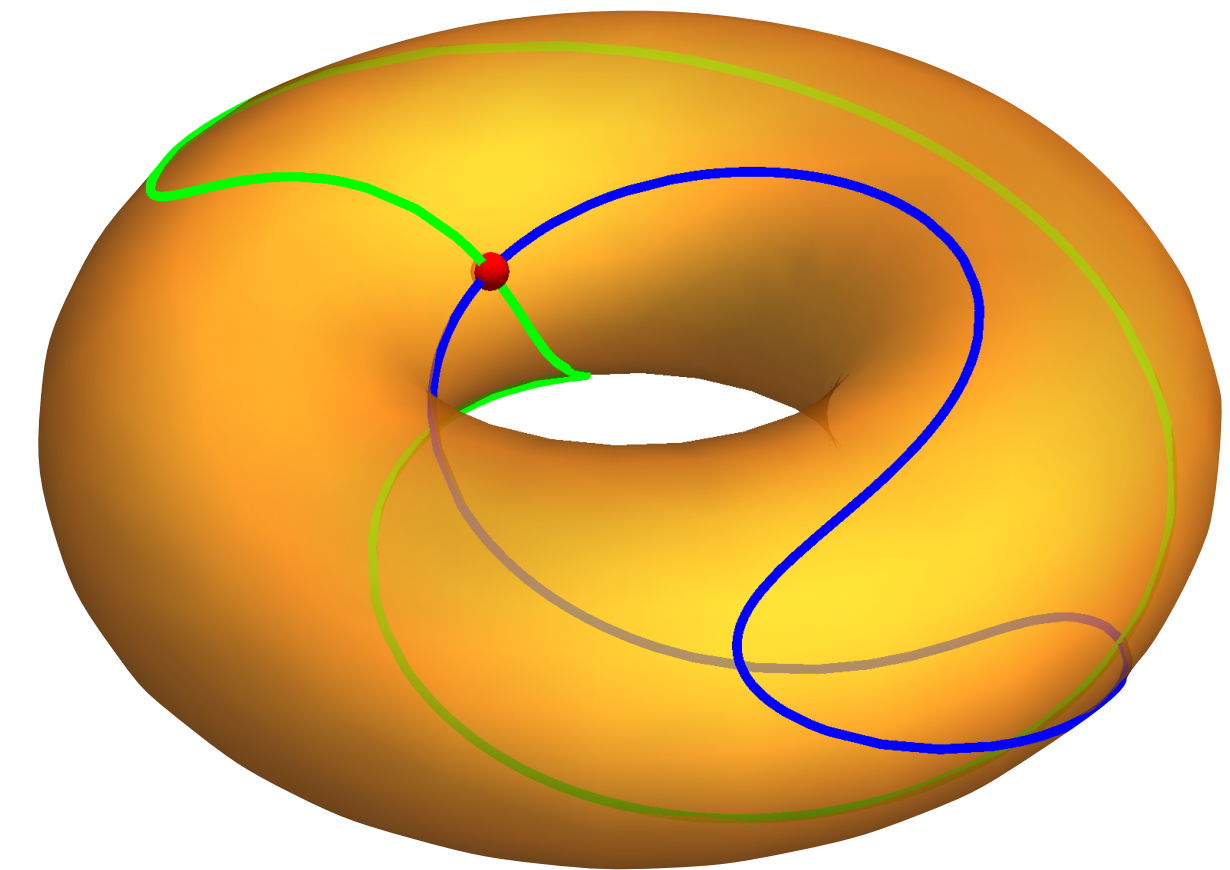
what are oxidation states, in the first place?

$$[0, L]^{3N}$$

$$\mathbb{T}^{3N}$$



PBC $\rightarrow$



$$\frac{L^2}{e} \int_0^T J_\alpha(t) dt = \frac{1}{Le} \int d\mu_\alpha[X] = Q_\alpha \in \mathbb{Z}$$

D.J. Thouless, 1983
R.Resta, 1990s

$$Q_\alpha(AA') = Q_\alpha(AA') = Q_\alpha[n_1 = 1, n_2 = 1]$$

what are oxidation states, in the first place?

$$Q_{\alpha}[\mathcal{C}] = \frac{1}{\ell} \mu_{\alpha}[\mathcal{C}]$$



what are oxidation states, in the first place?

$$\begin{aligned} Q_{\alpha}[\mathcal{C}] &= \frac{1}{\ell} \mu_{\alpha}[\mathcal{C}] \\ &= Q_{\alpha}(n_{1x}, n_{1y}, n_{1z}, \dots, n_{Nz}) \end{aligned}$$



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- All loops can be shrunk to a point without closing the gap (*strong adiabaticity*);
- Any two like atoms can be swapped without closing the gap



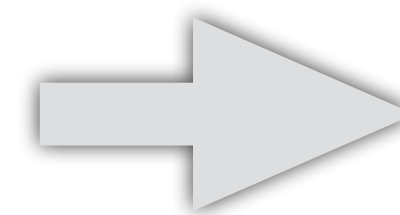
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$$q_{i\alpha\beta} = q_{S(i)} \delta_{\alpha\beta}$$

atomic oxidation state

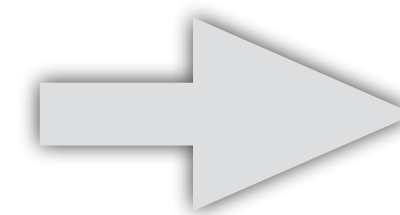
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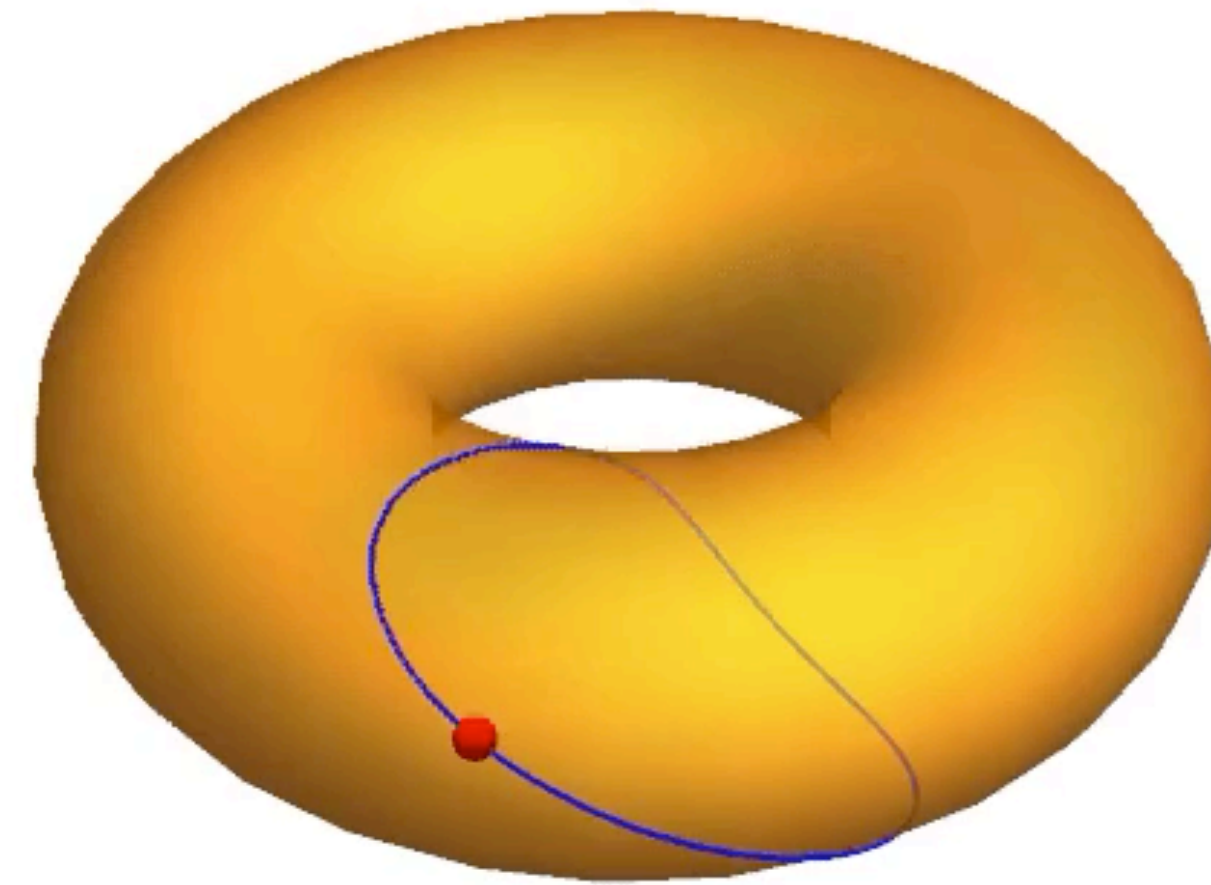
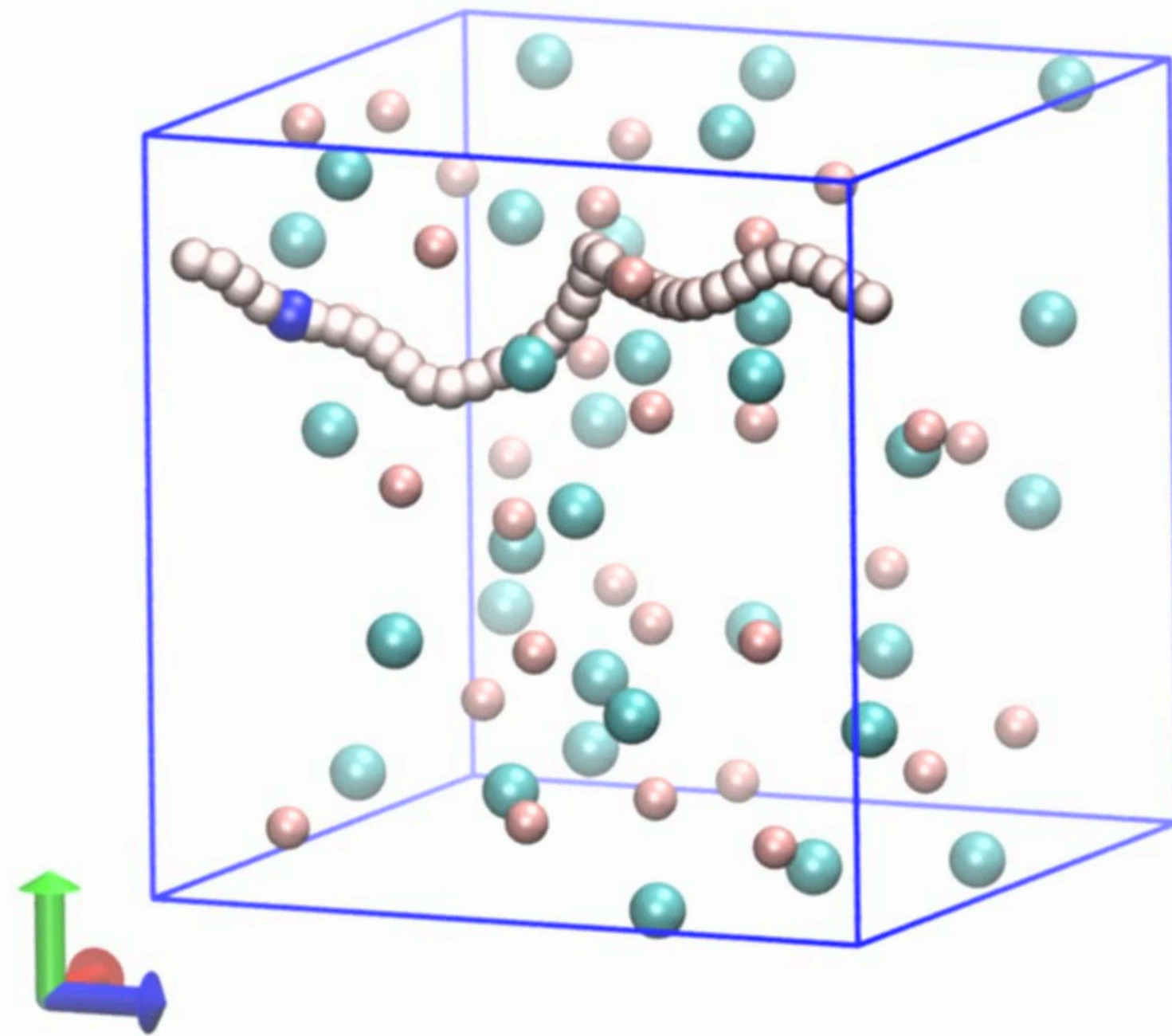


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atomic oxidation state

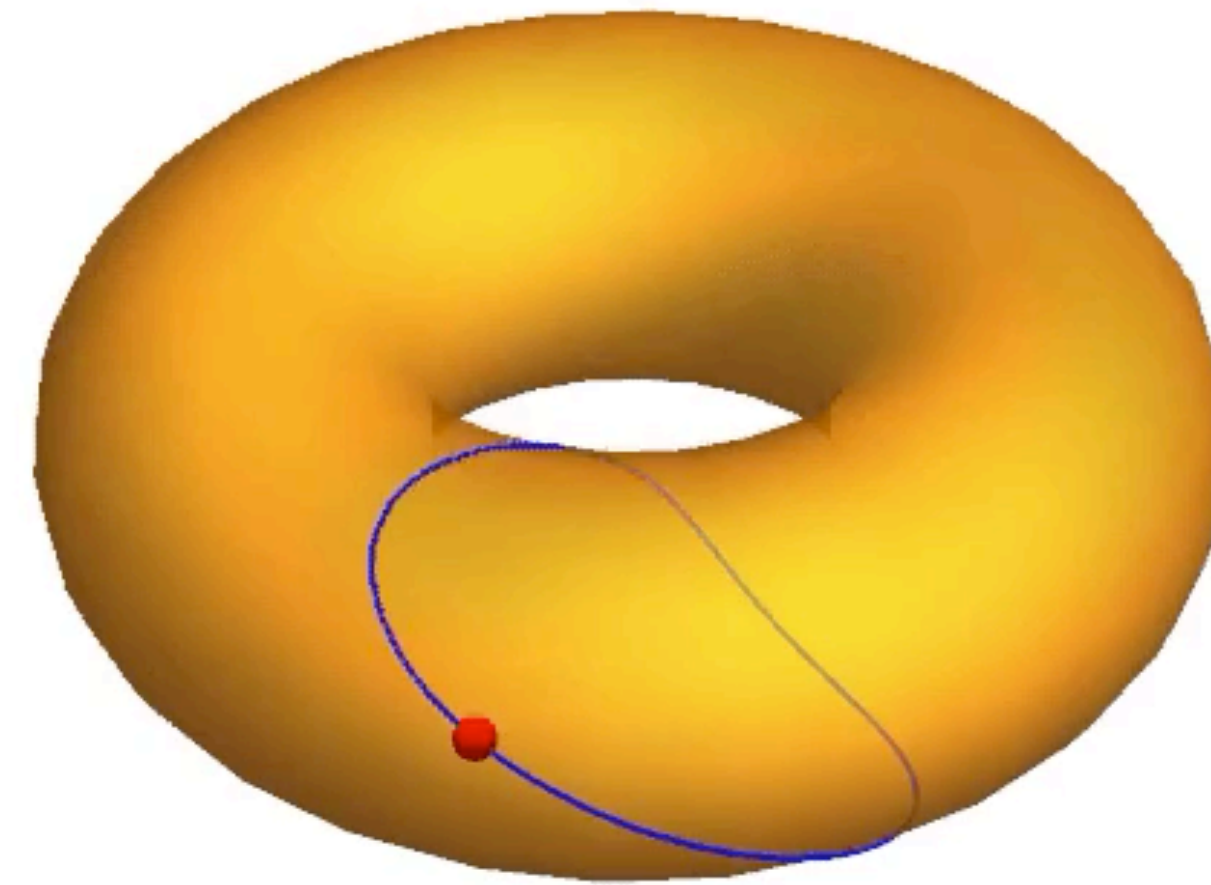
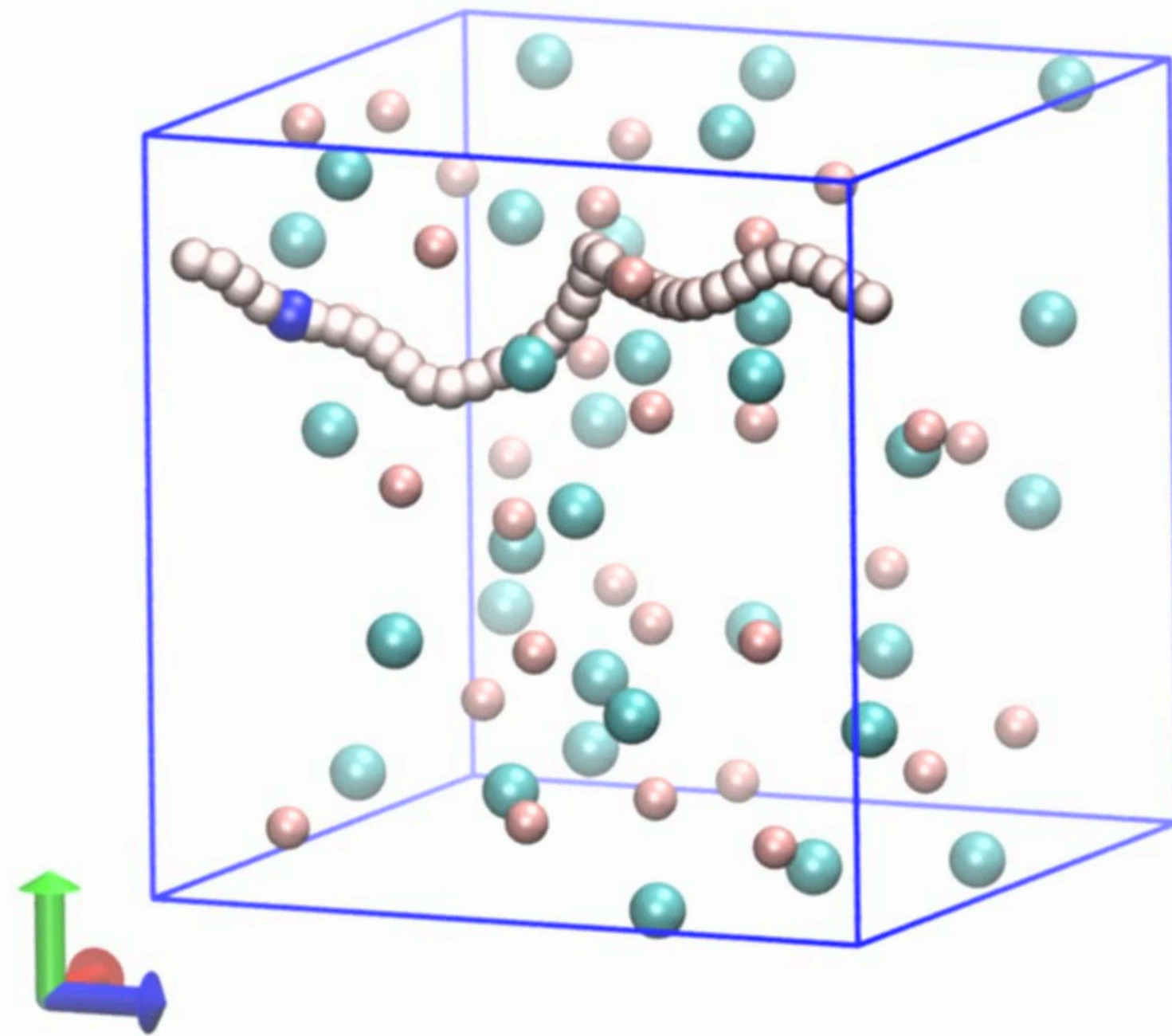
... they are topological invariants!

a numerical experiment on molten KCl



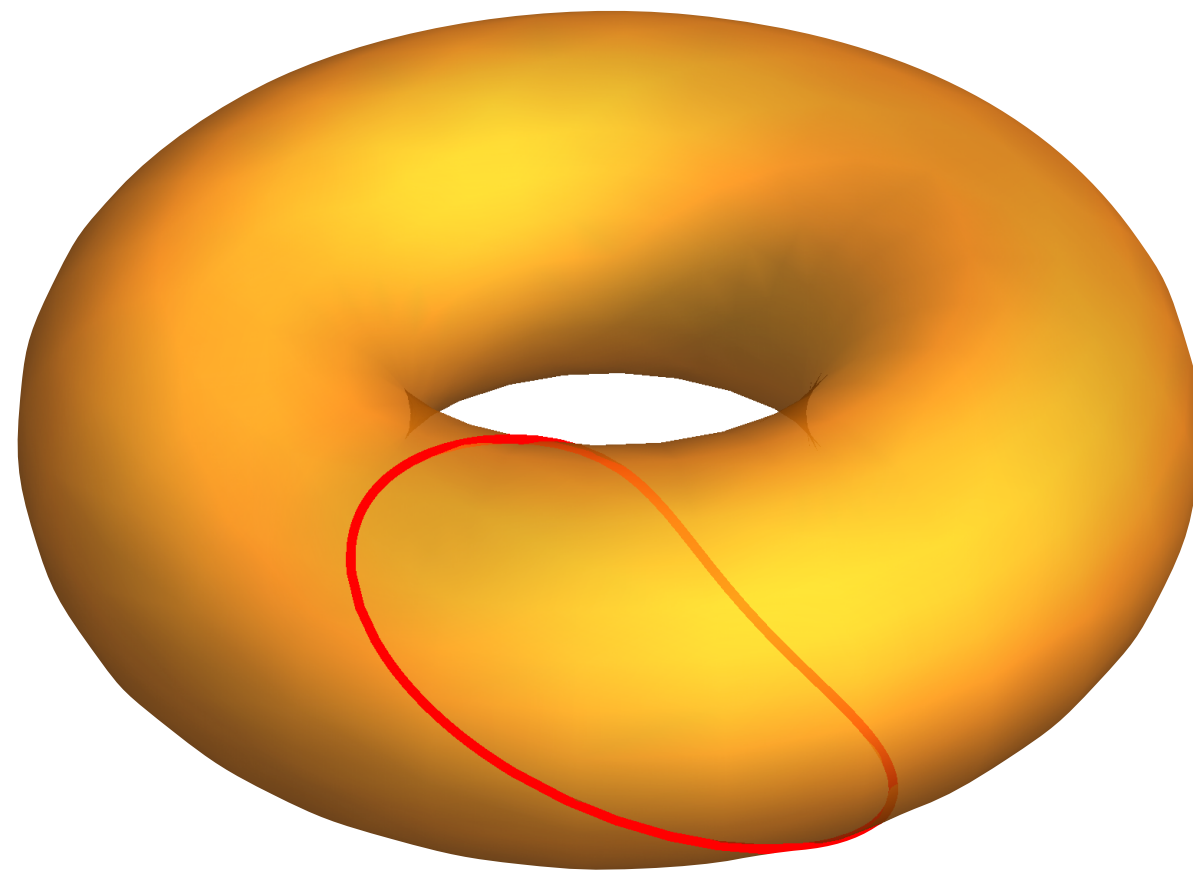
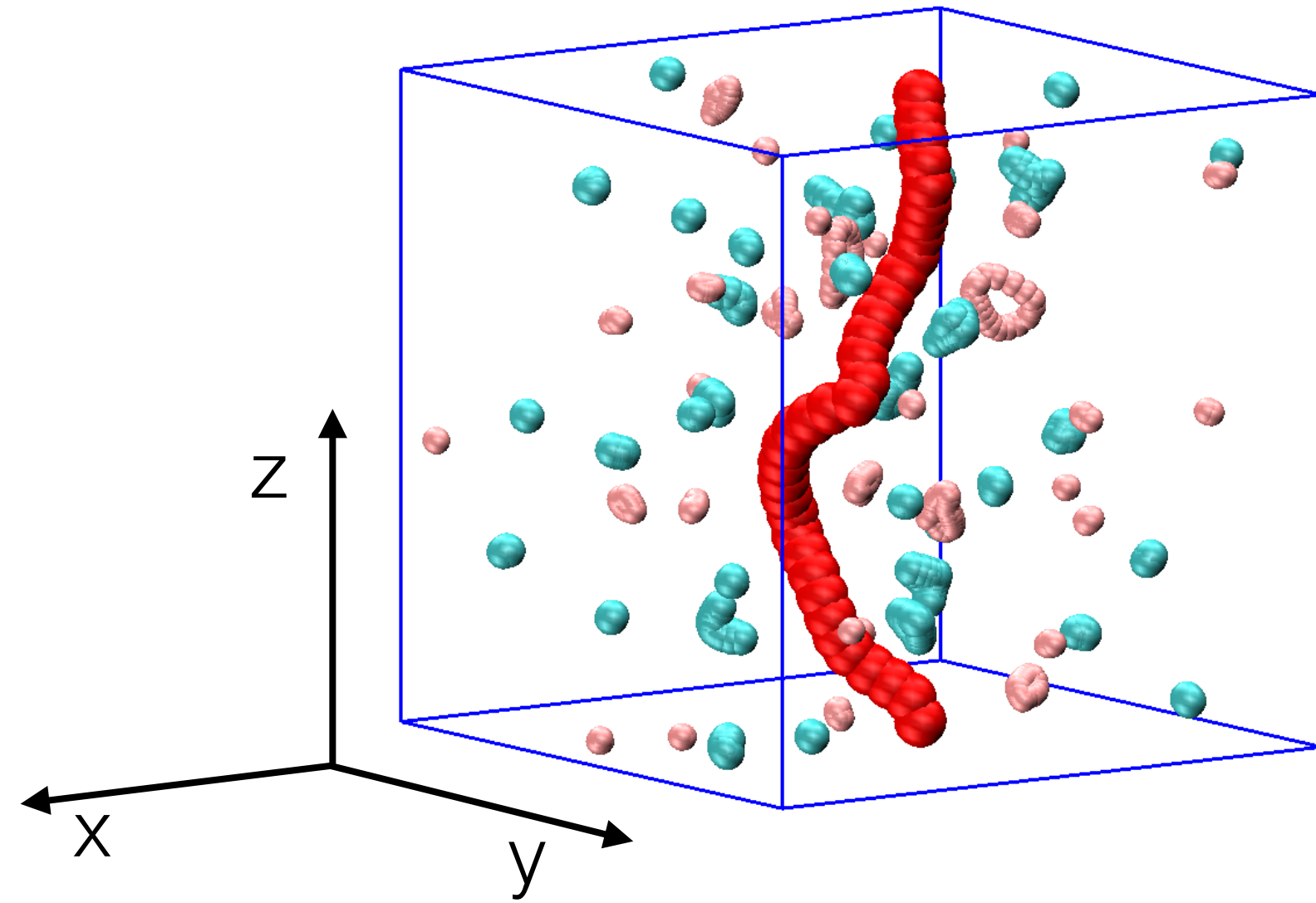
a topologically non-trivial minimum-energy path
connecting two identical configurations of a ionic melt

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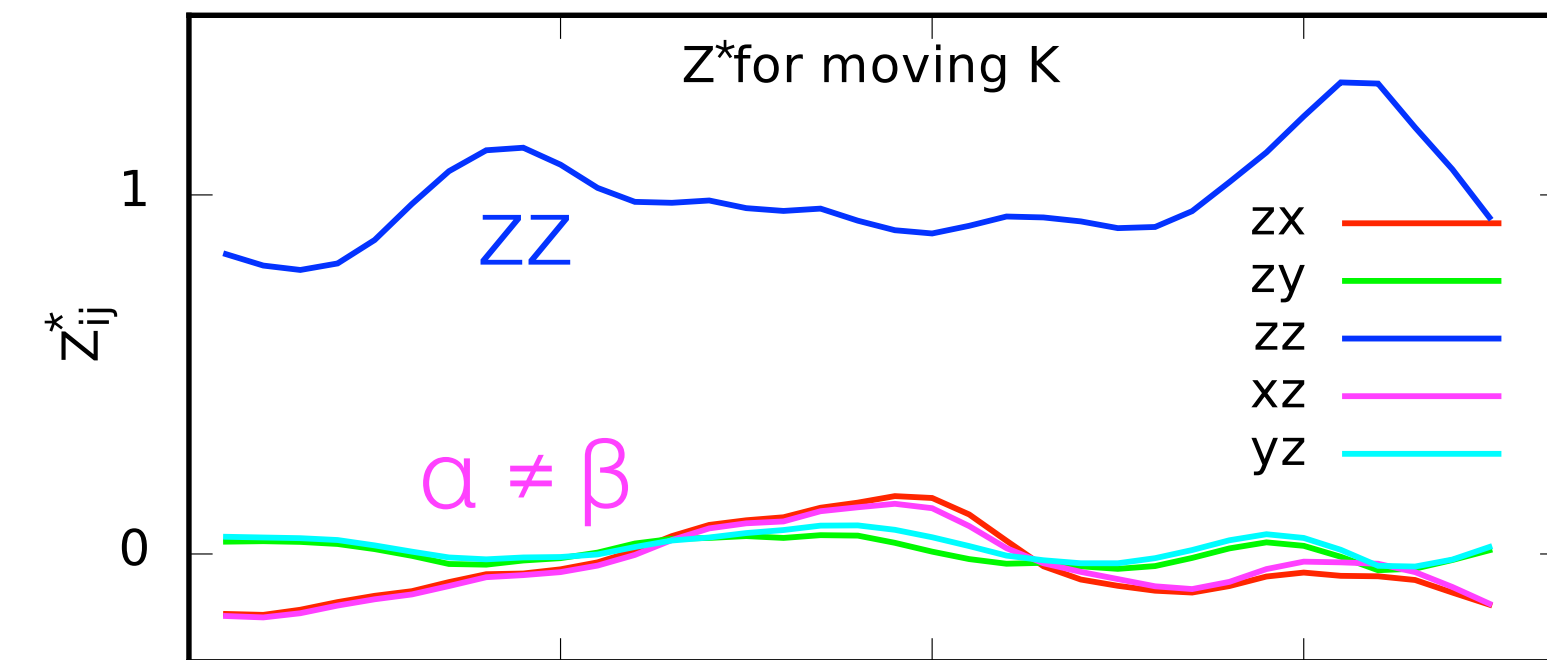
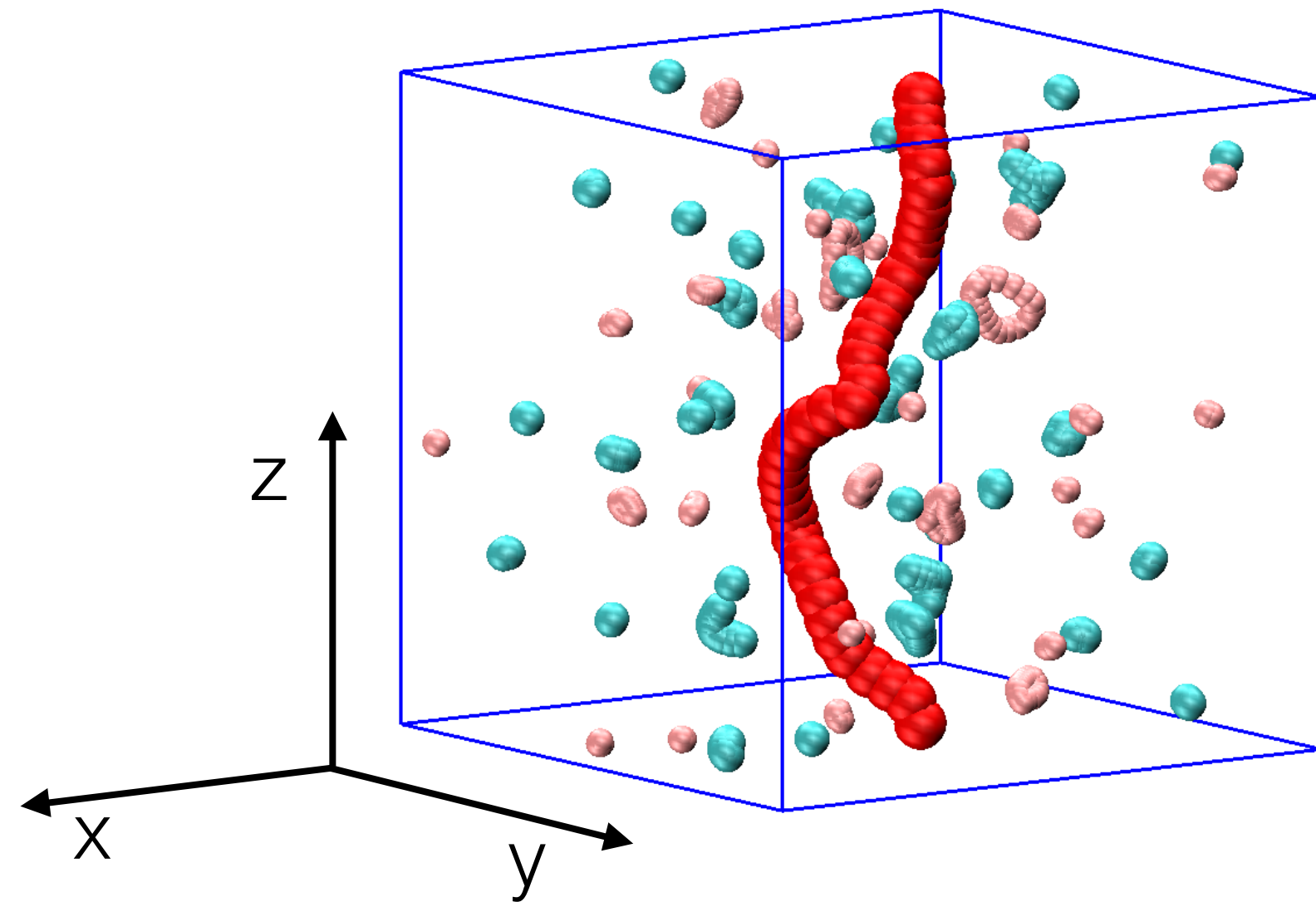


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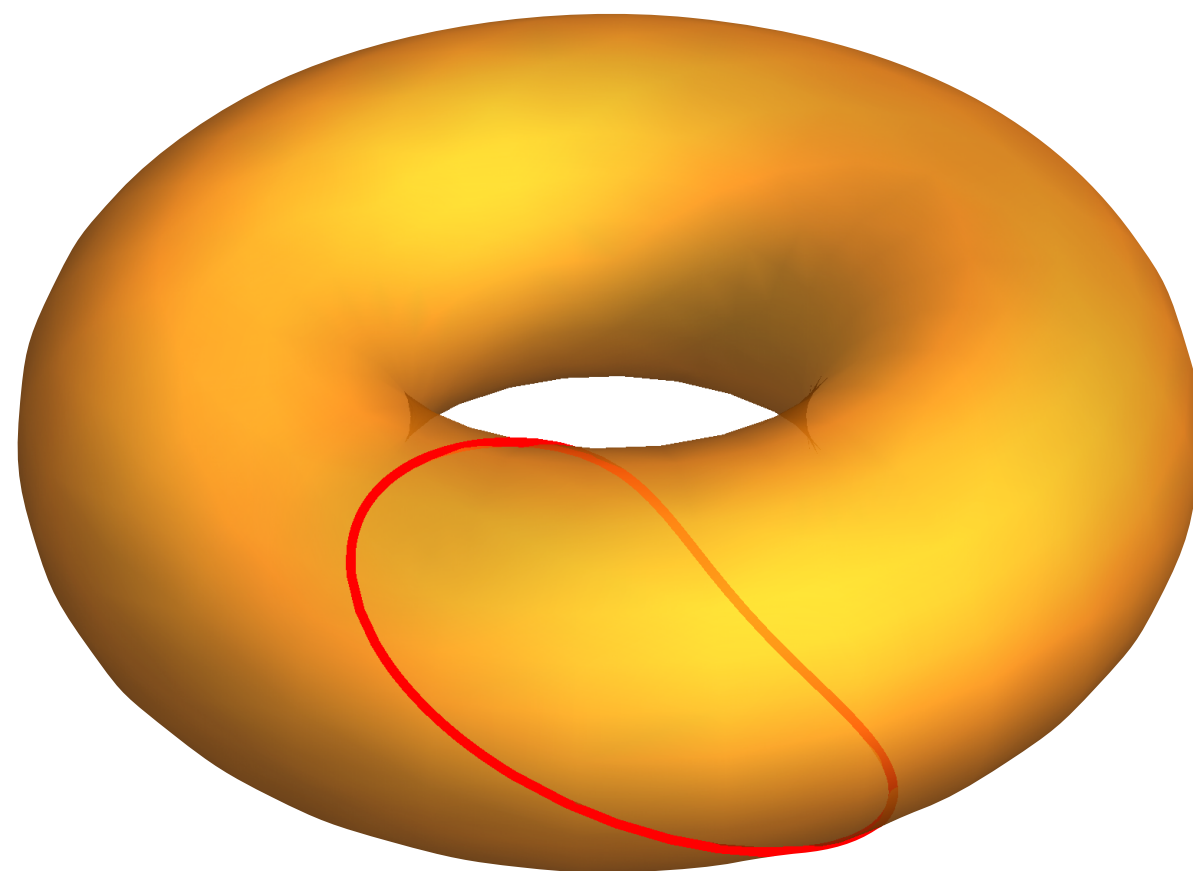
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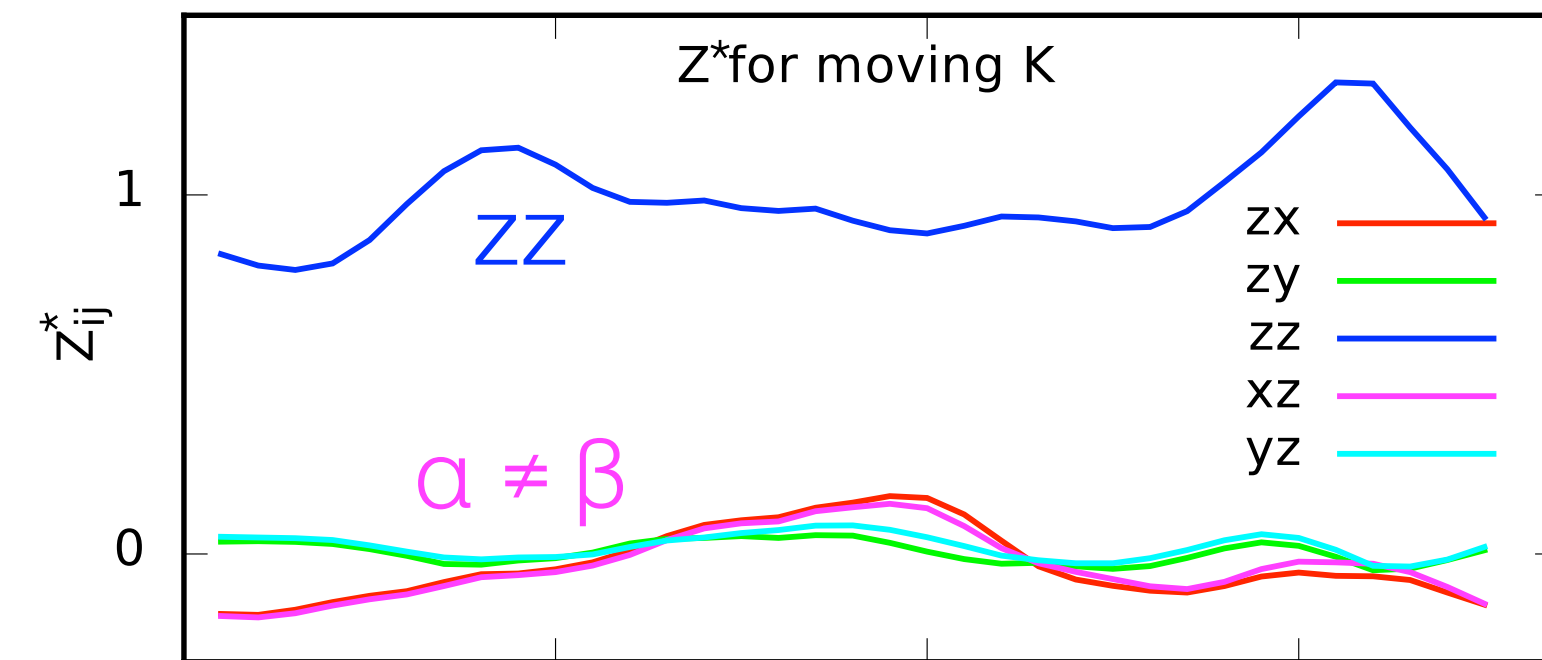
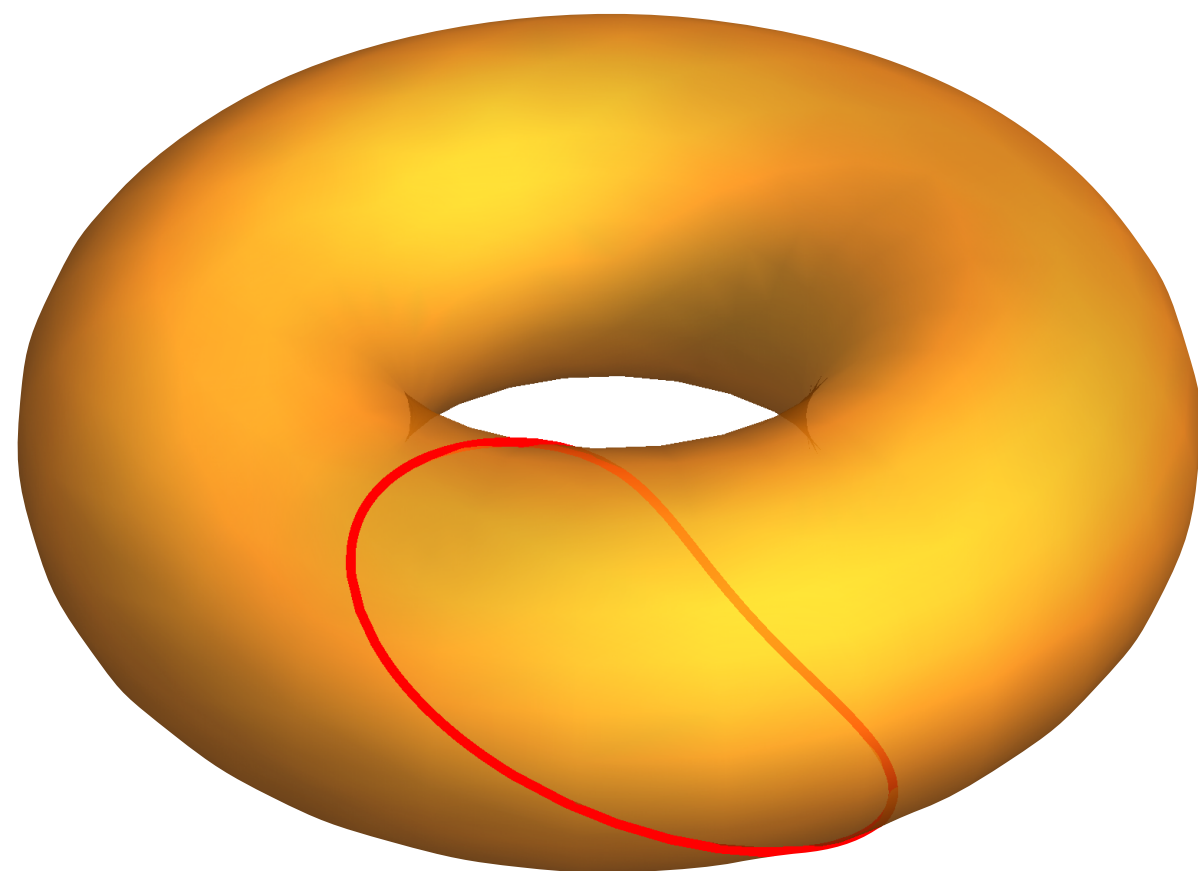
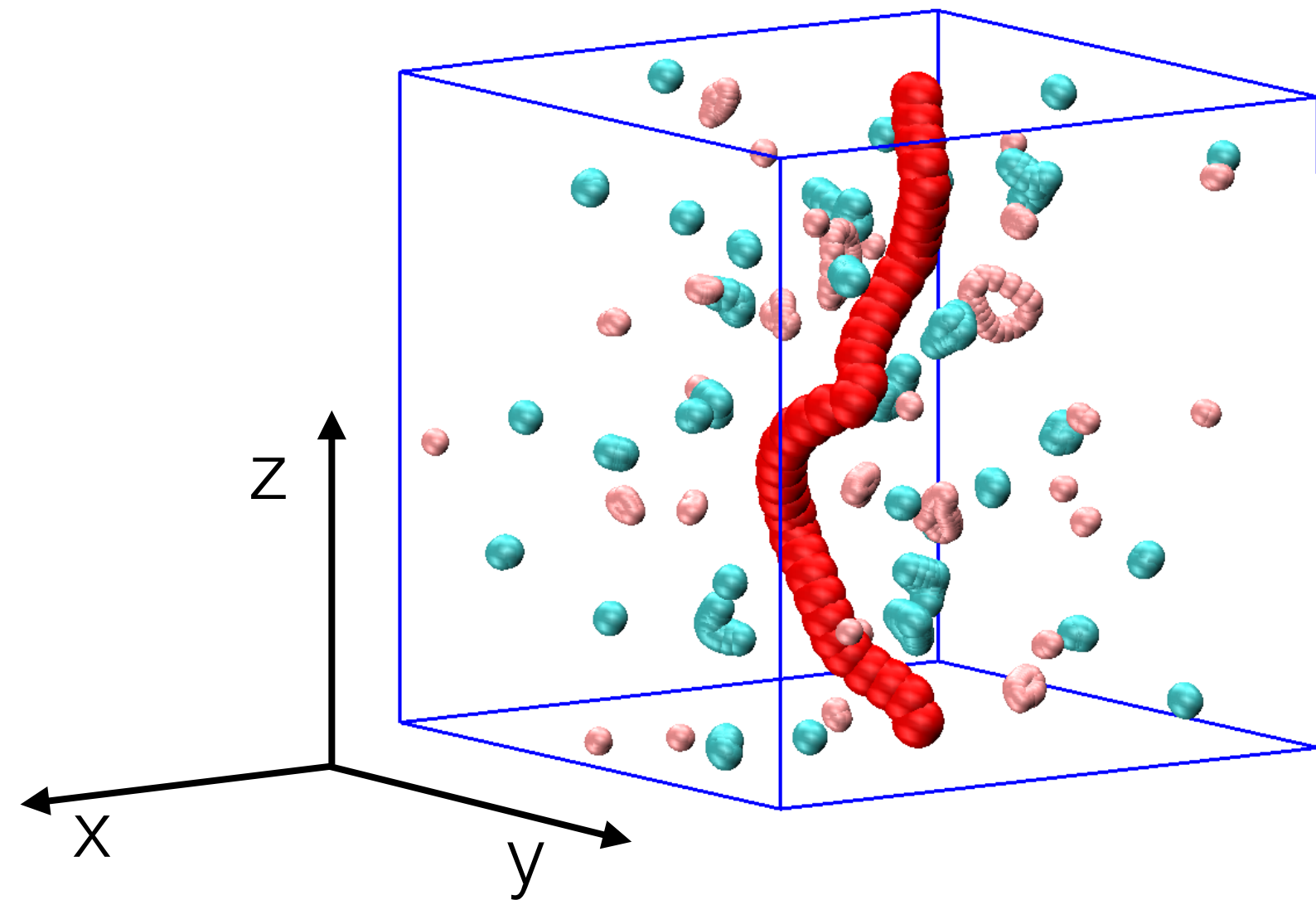
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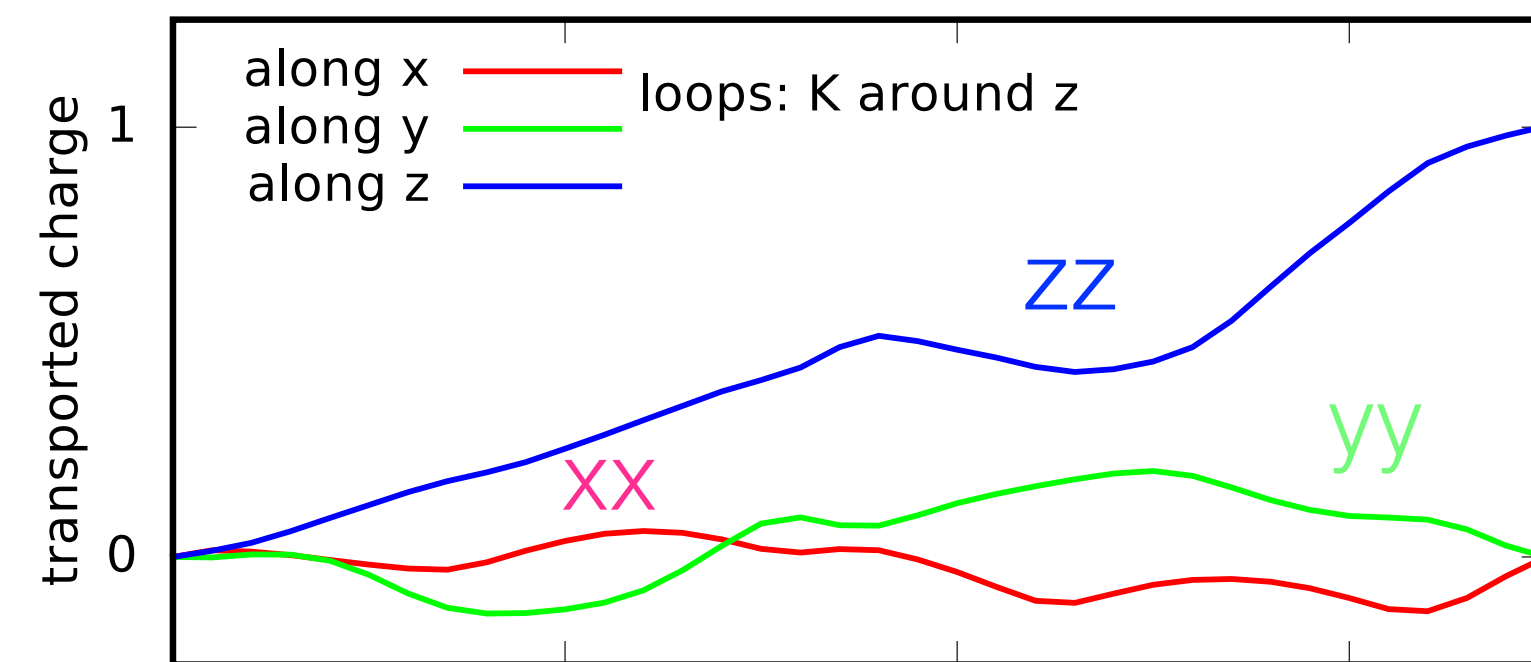
effective
charge



a numerical experiment on molten KCl



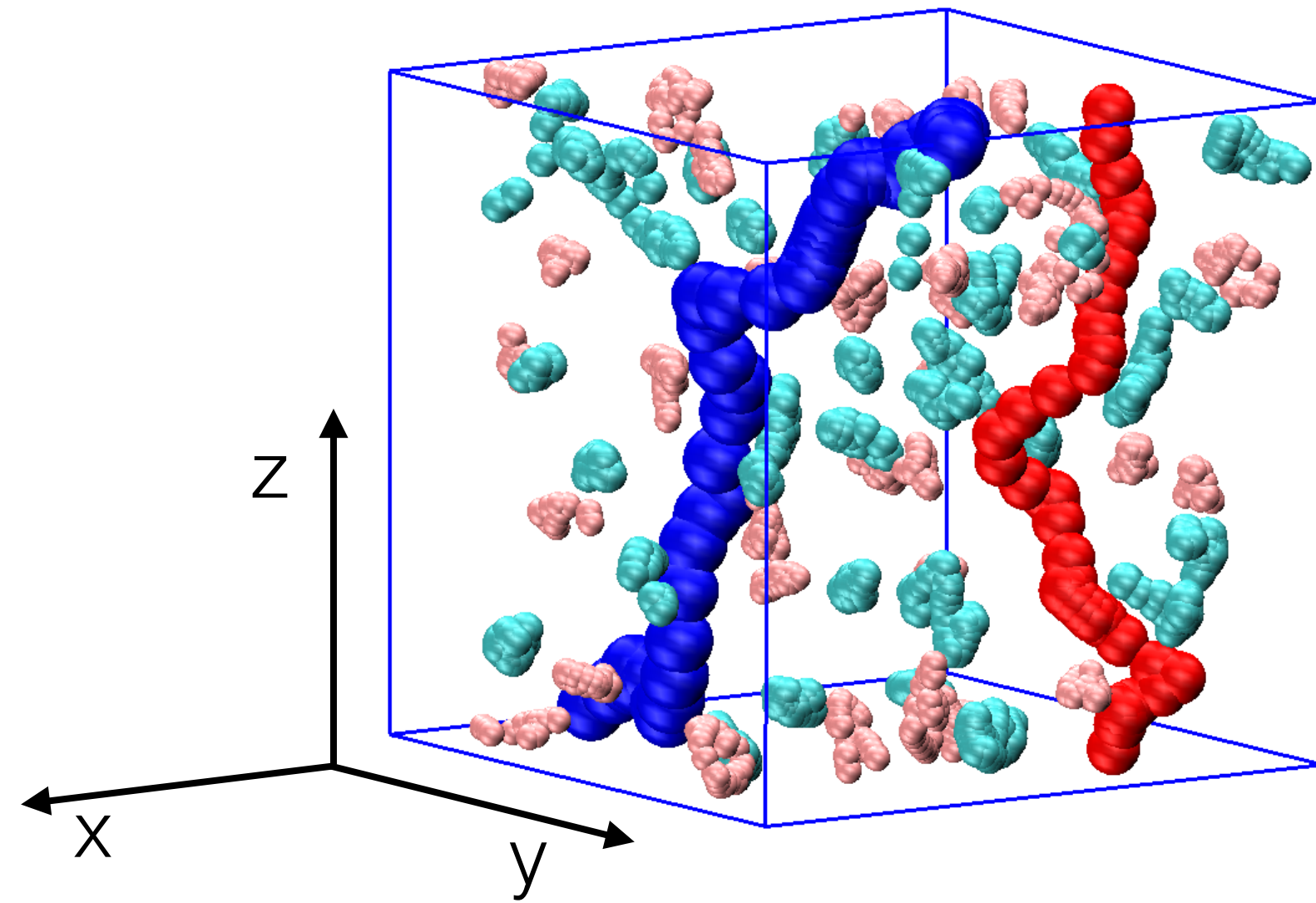
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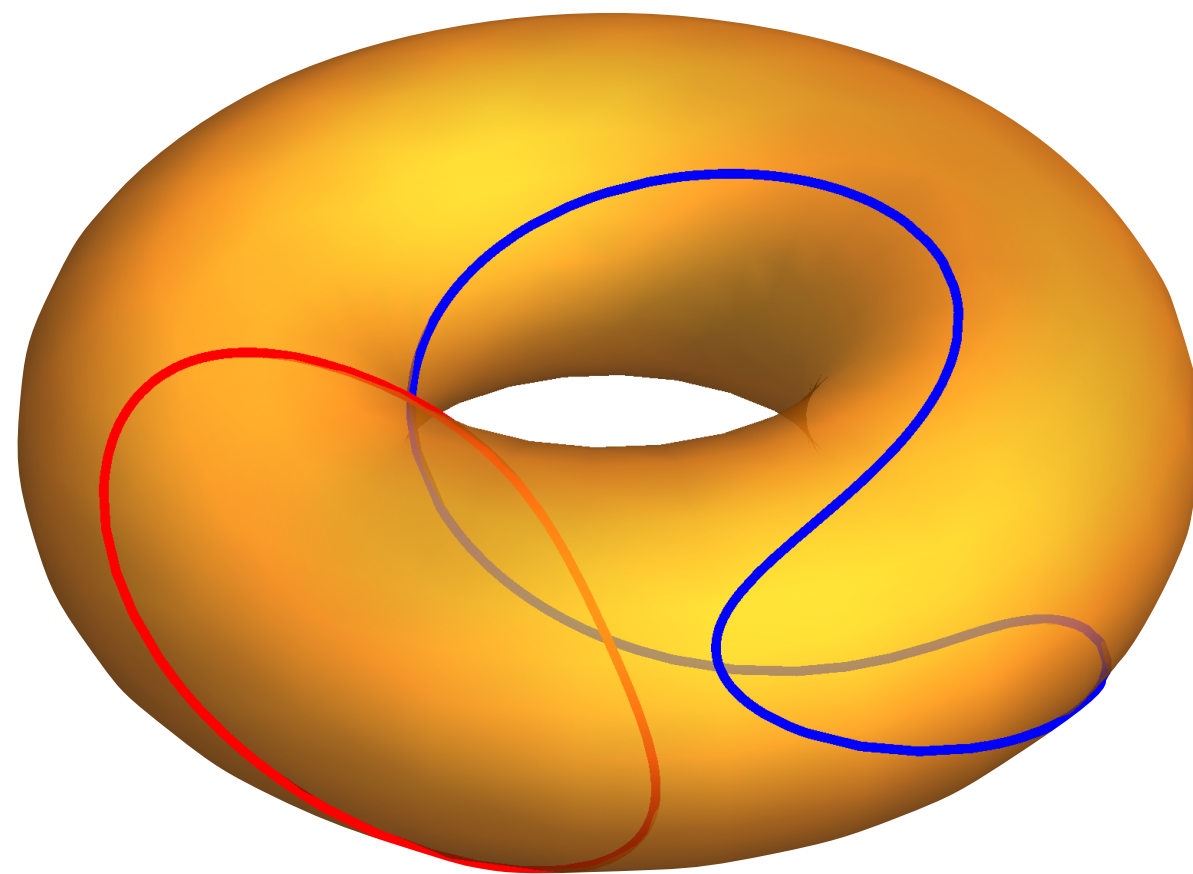
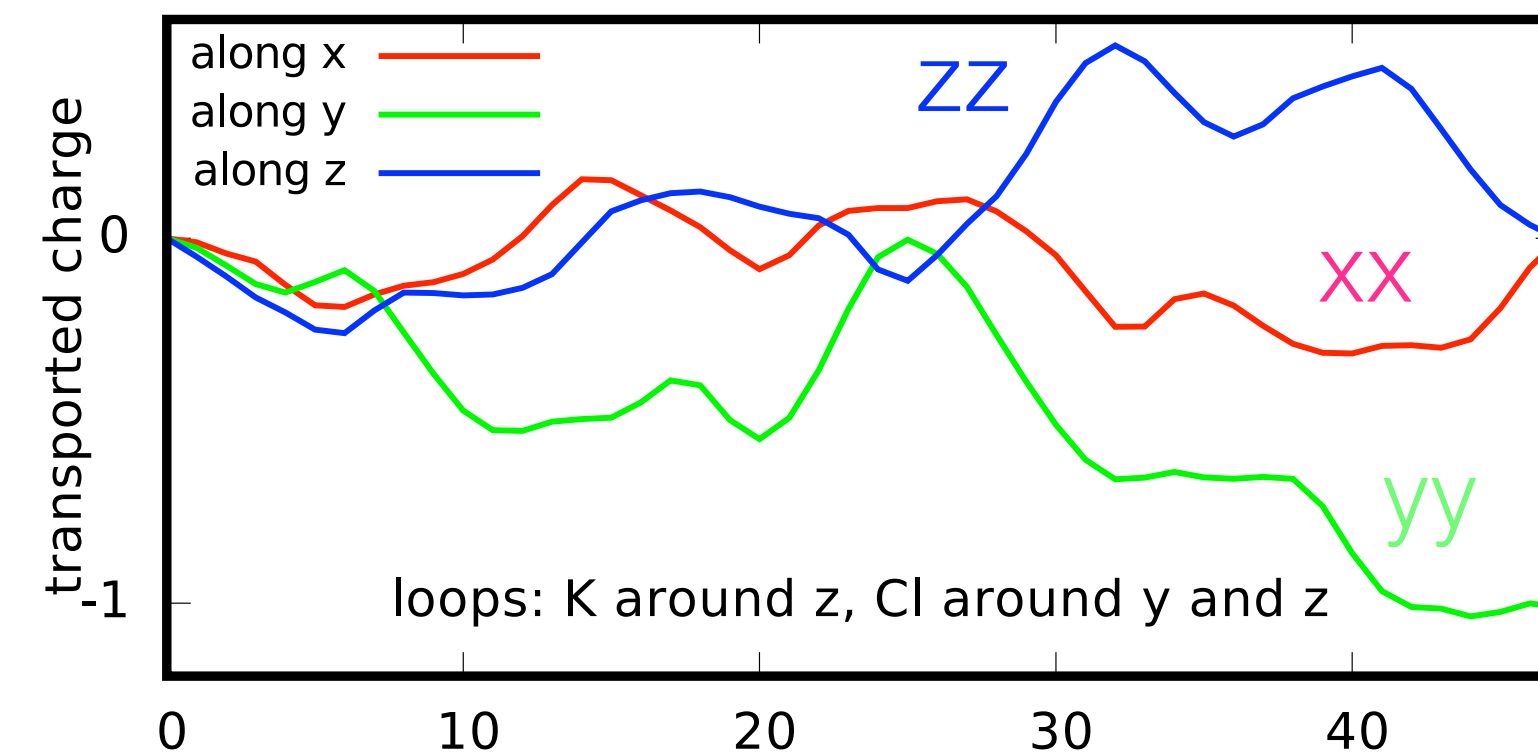
topological
charge

$$Q_x = -0.000(6); \quad Q_y = 0.000(2); \quad Q_z = 1.00(18)$$

a numerical experiment on molten KCl

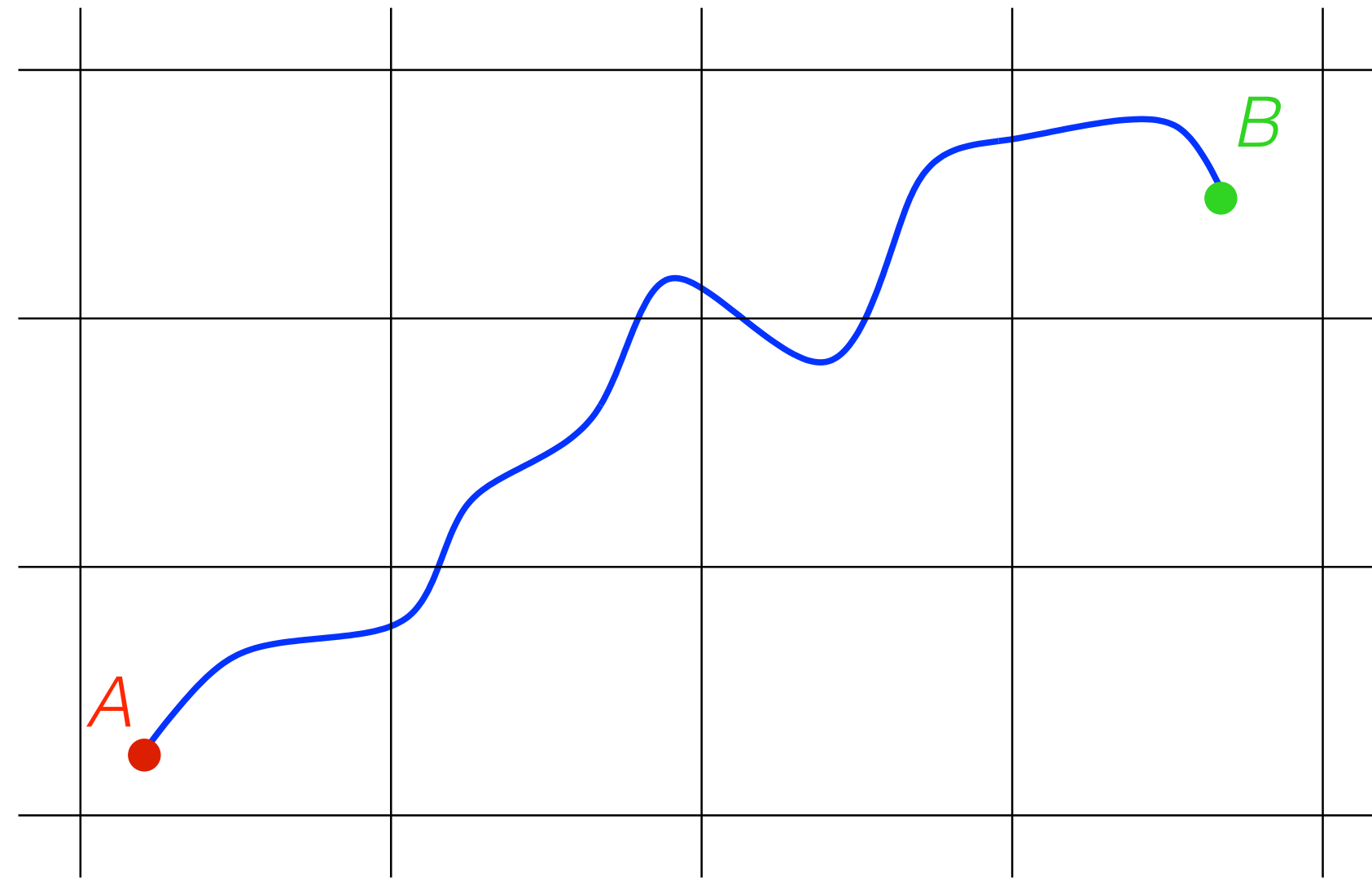


$$\begin{array}{ll} Q_z[\text{Cl}] = -1 & Q_y[\text{Cl}] = -1 \\ Q_z[\text{K}] = 1 & Q_z[\text{K}] = 0 \end{array}$$



the charges transported by K and Cl
around z cancel exactly

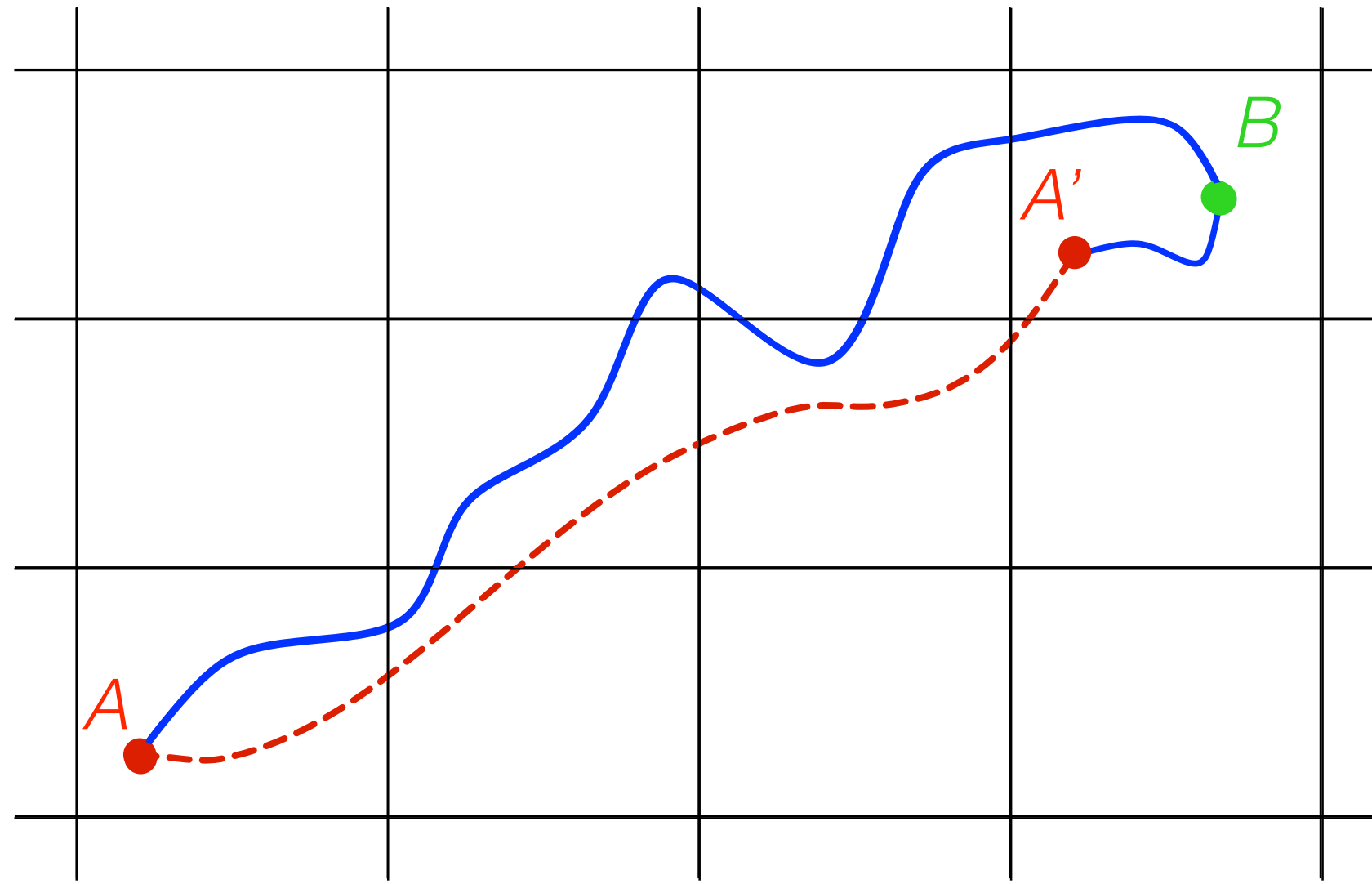
gauge invariance of charge transport



$$\sigma \propto \lim_{t \rightarrow \infty} \frac{1}{2t} \text{var} [\mu_{AB}(t)]$$
$$\mu_{AB}(t) = \int_0^t J(t') dt'$$

gauge invariance of charge transport

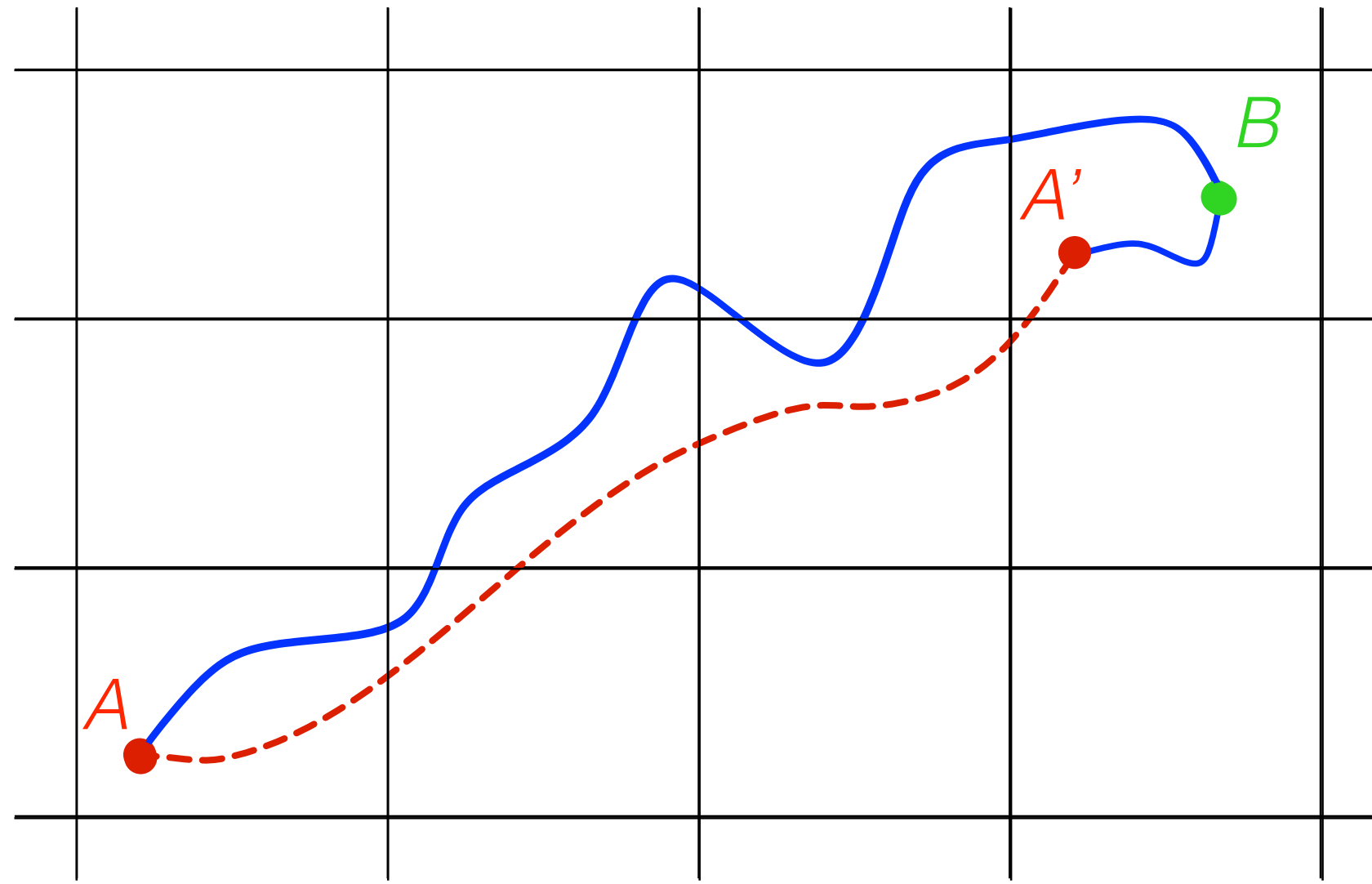
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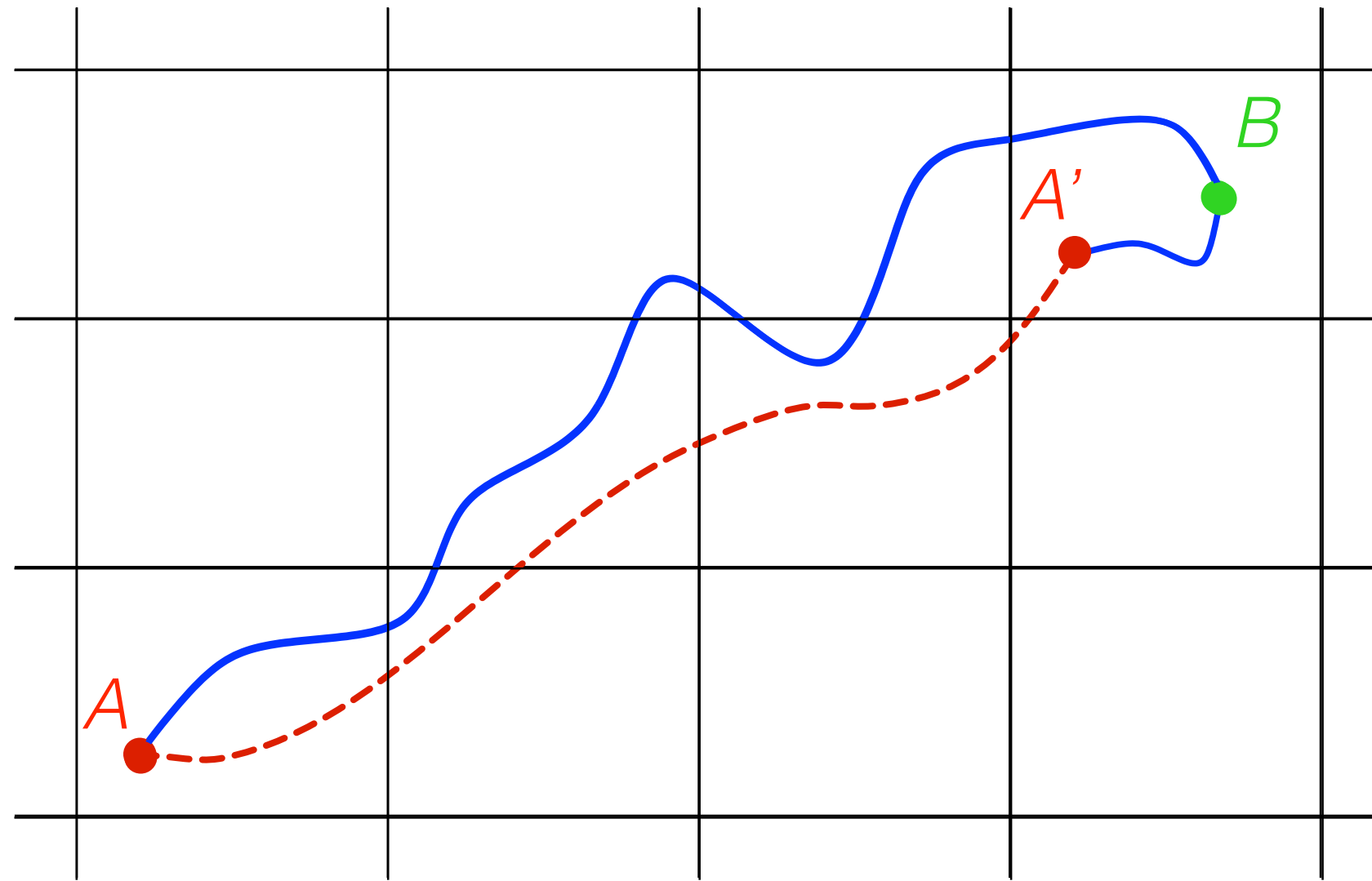


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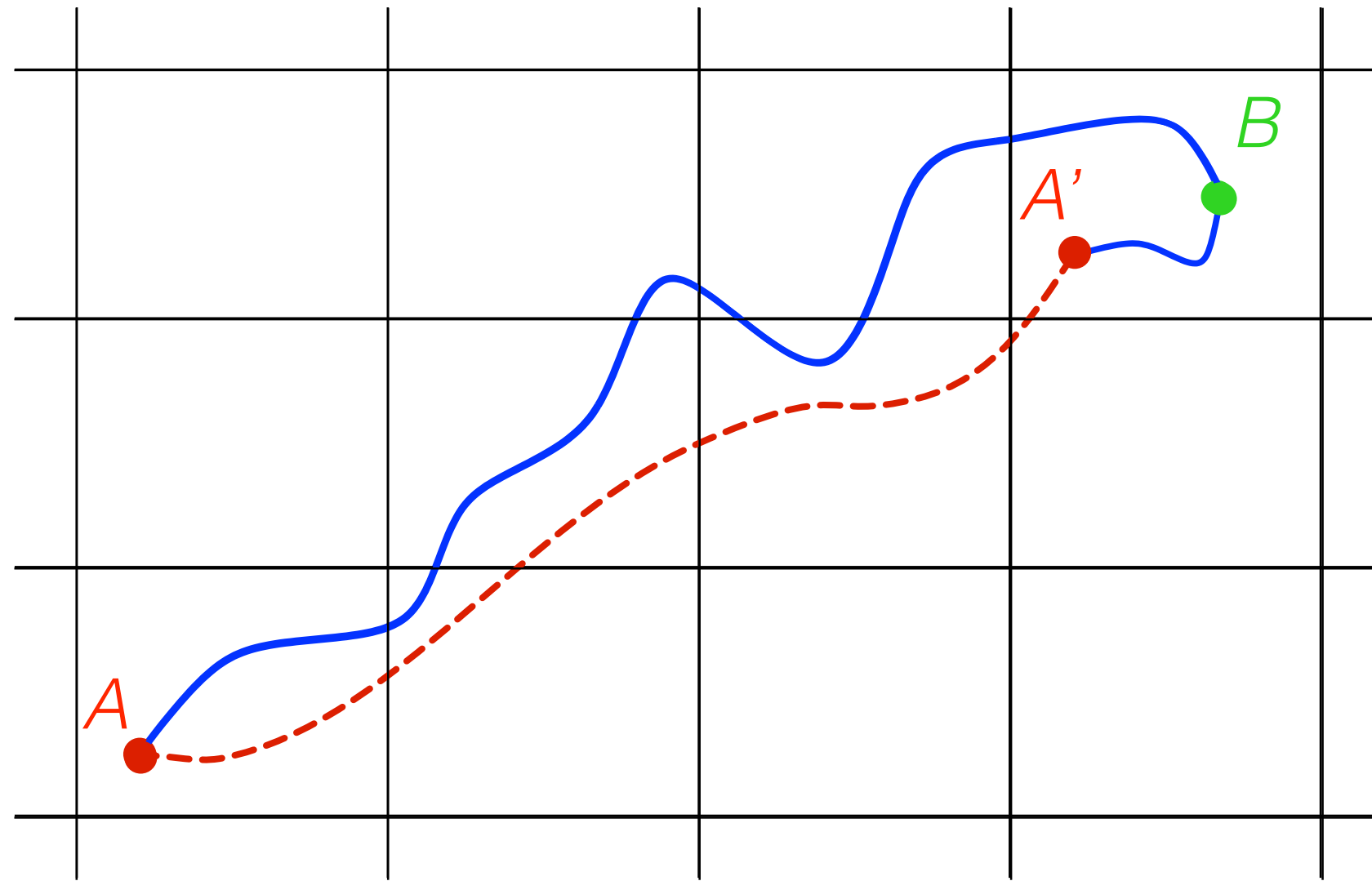
$$= \mu_{AA'} + \cancel{\mu_{A'B}}$$

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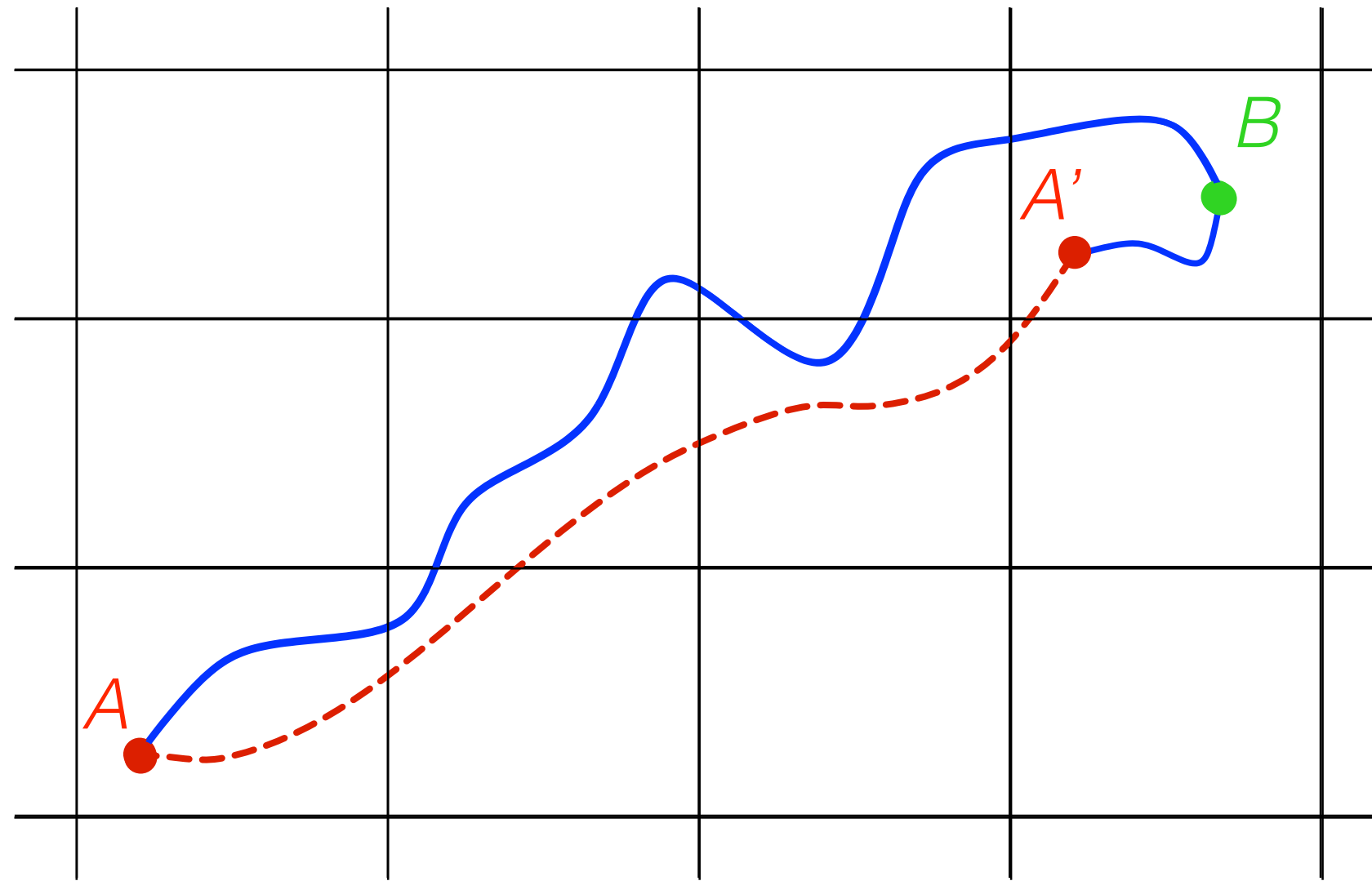
$$J'_\alpha = \sum_i q_{S(i)} v_{i\alpha}$$

$$\mu_{AB}(t) = \mu_{AA'} + \mathcal{O}(1)$$

$$= \overset{||}{\mu'}_{AA'} + \mathcal{O}(1) \quad (\text{Thouless})$$

gauge invariance of charge transport

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$$= \mu_{AA'} + \cancel{\mu_{A'B}}$$

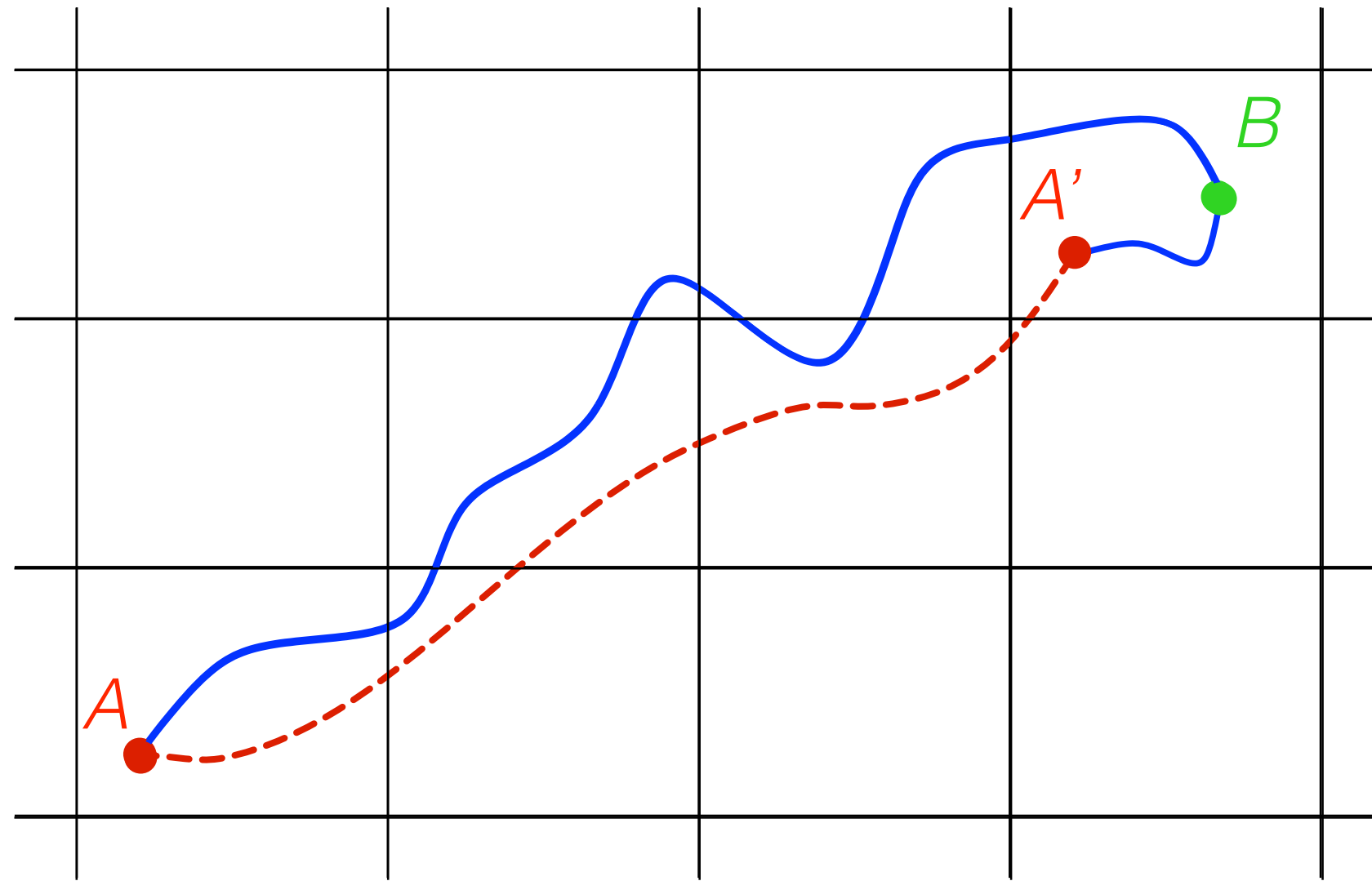
$$\begin{aligned}\mu_{AB}(t) &= \mu_{AA'} + \mathcal{O}(1) \\ &\stackrel{||}{=} \mu'_{AA'} + \mathcal{O}(1) \quad (\text{Thouless}) \\ &= \mu'_{AB} + \mathcal{O}(1)\end{aligned}$$



gauge invariance of charge transport

$$\hat{H}(B) \neq \hat{H}(A)$$

$$\hat{H}(A') = \hat{H}(A)$$



$$\sigma \propto \lim_{t \rightarrow \infty} \left(\frac{1}{2t} \right) \text{var} [\mu_{AB}(t)]$$

$$\mu_{AB}(t) = \int_0^t J(t') dt'$$

$$= \mu_{AA'} + \cancel{\mu_{A'B}}$$

$$J_\alpha = \sum_{i\beta} Z_{i\alpha\beta}^* v_{i\beta}$$

$$J'_\alpha = \sum_i q_{S(i)} v_{i\alpha}$$

$$\mu_{AB}(t) = \mu_{AA'} + \mathcal{O}(1)$$

$$= \overset{||}{\mu'}_{AA'} + \mathcal{O}(1) \quad (\text{Thouless})$$

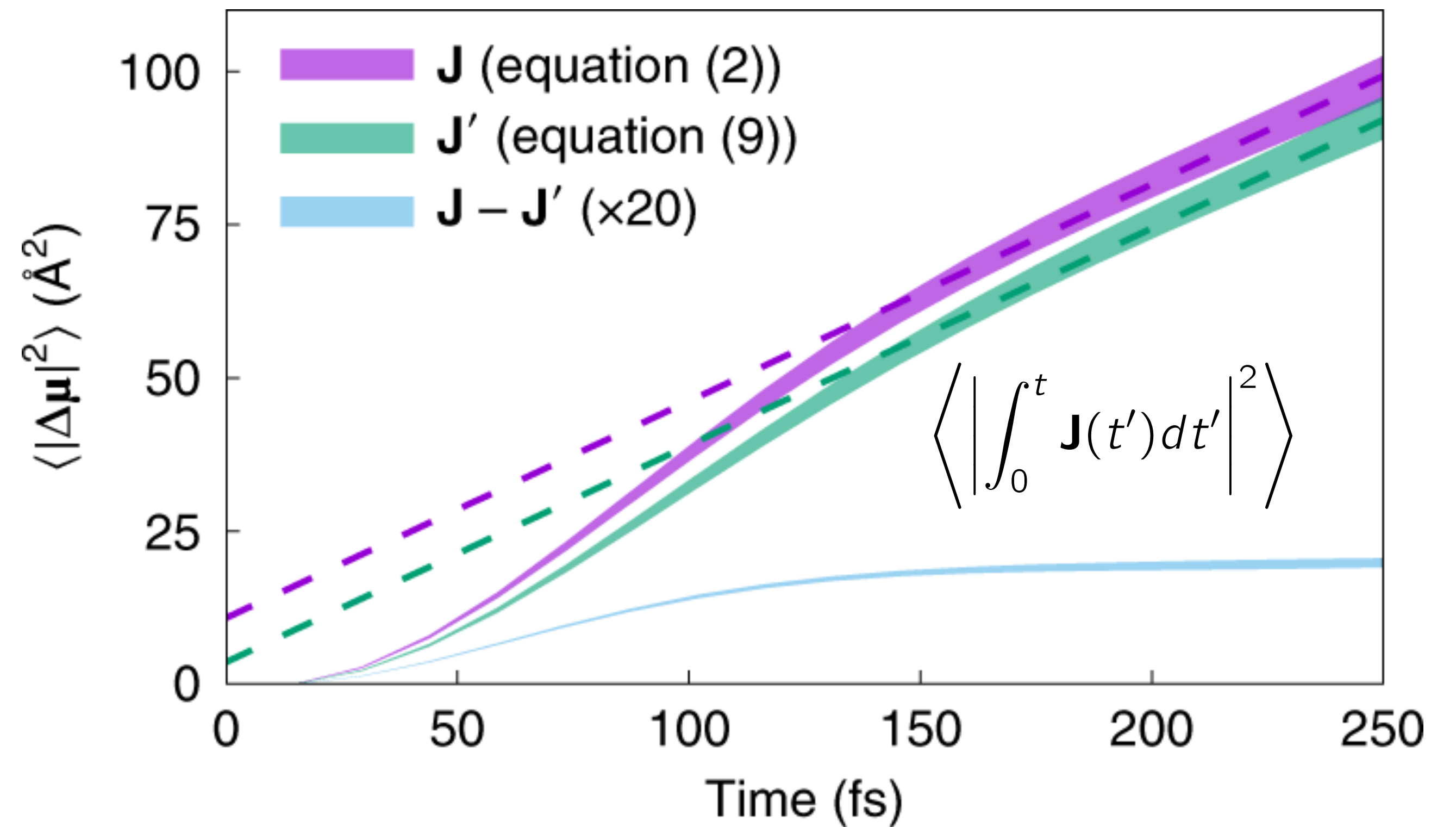
$$= \mu'_{AB} + \mathcal{O}(1)$$

$$\sigma = \sigma'$$

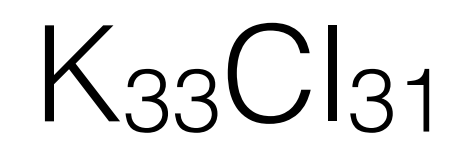
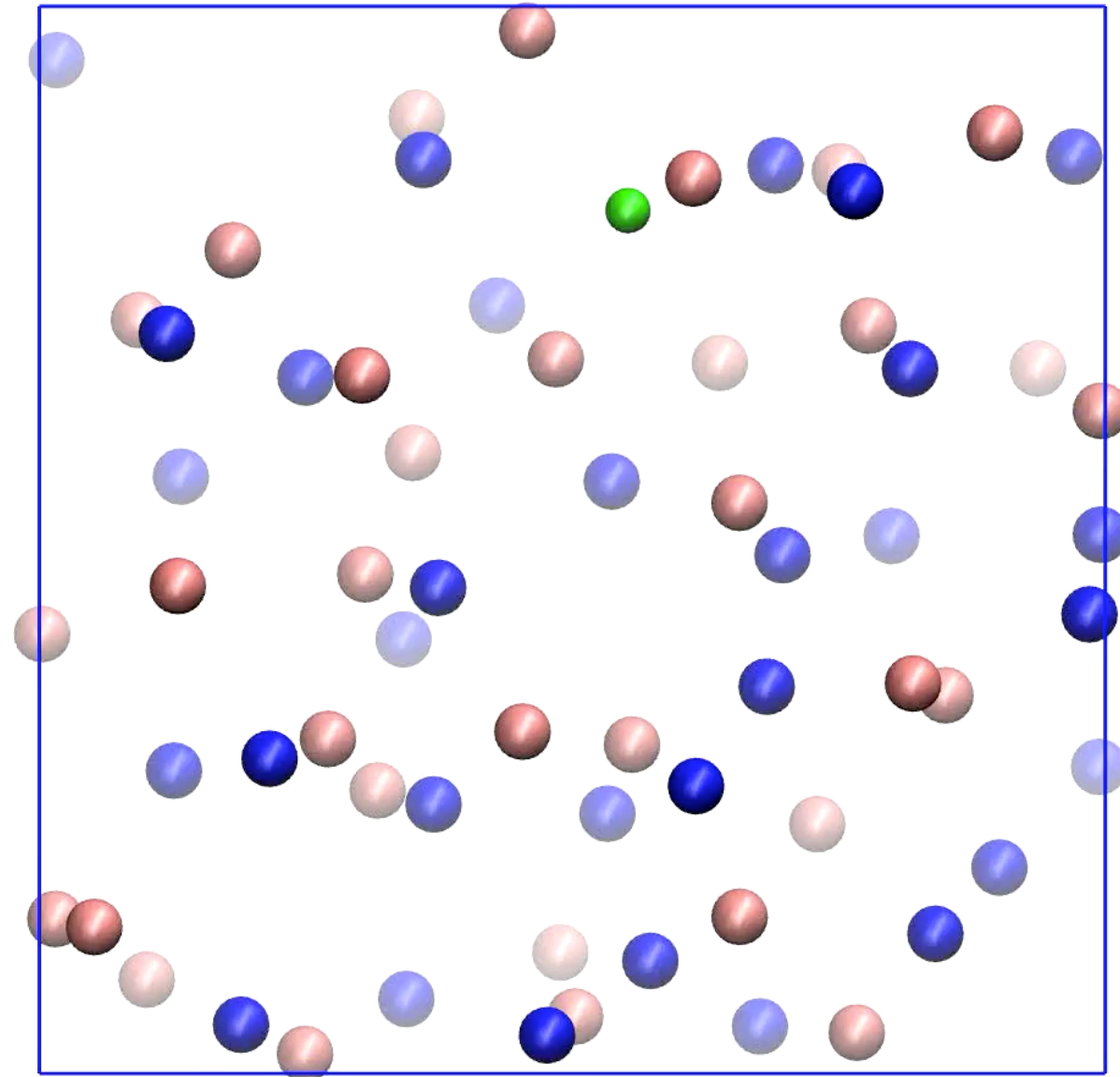
currents from atomic oxidation numbers

$$(2) \quad J_{\alpha} = \sum_{i\beta} Z_{i\alpha\beta}^* v_{i\beta}$$

$$(9) \quad J'_{\alpha} = \sum_i q_{S(i)} v_{i\alpha}$$

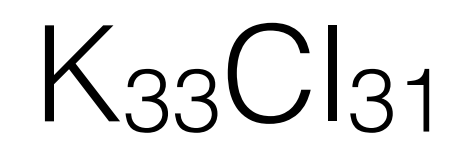
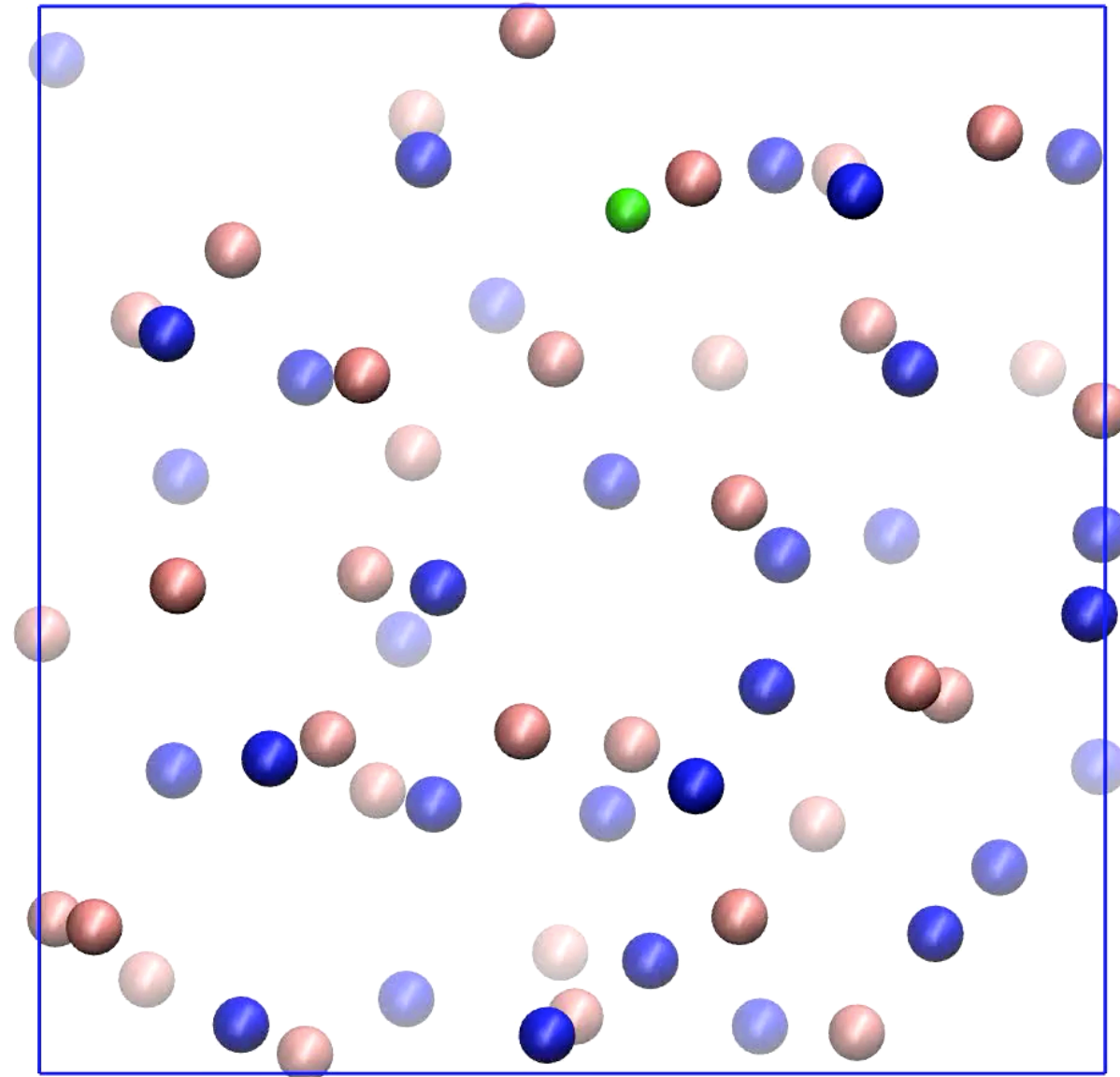


non-stoichiometric melts



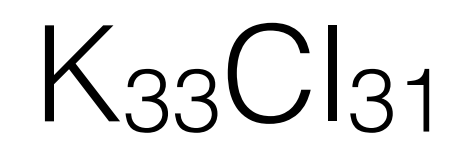
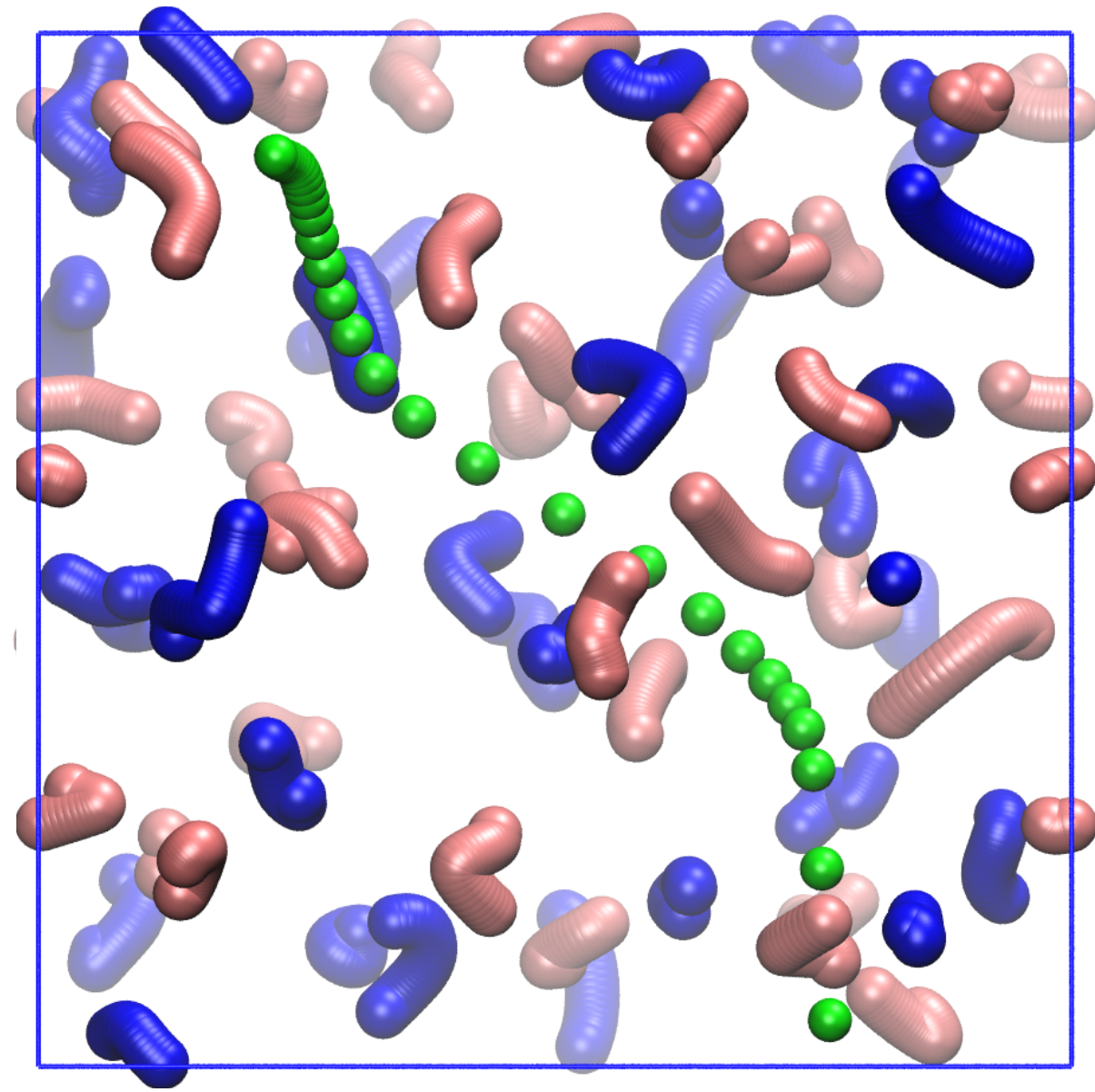
$$x \approx 0.06$$

non-stoichiometric melts



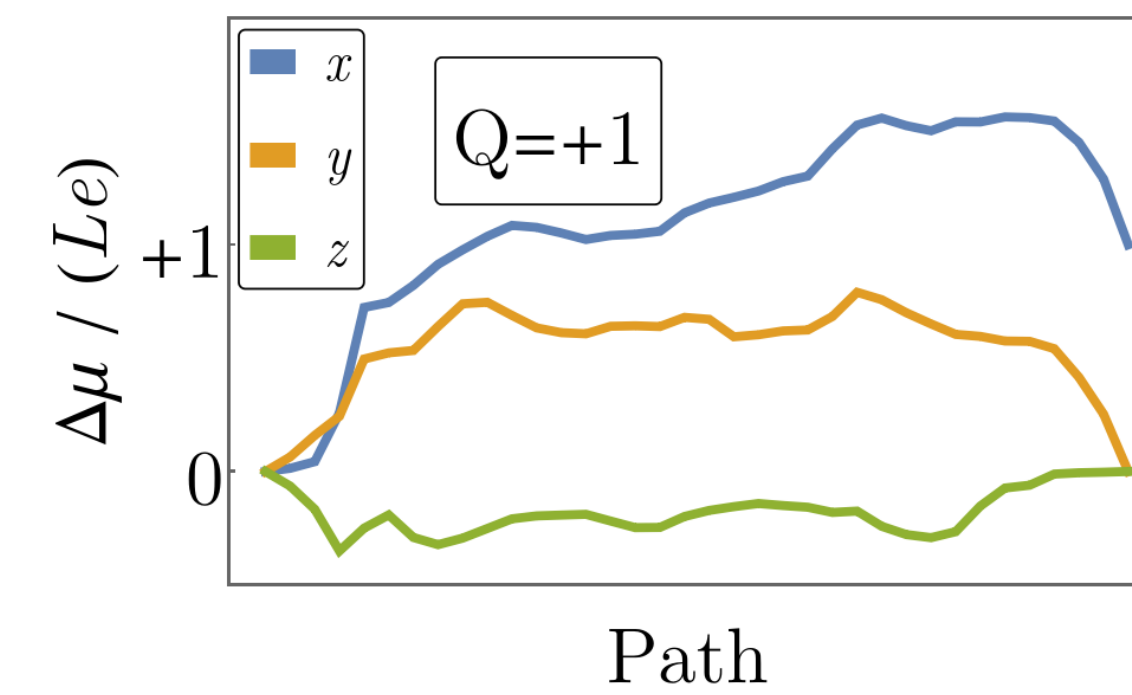
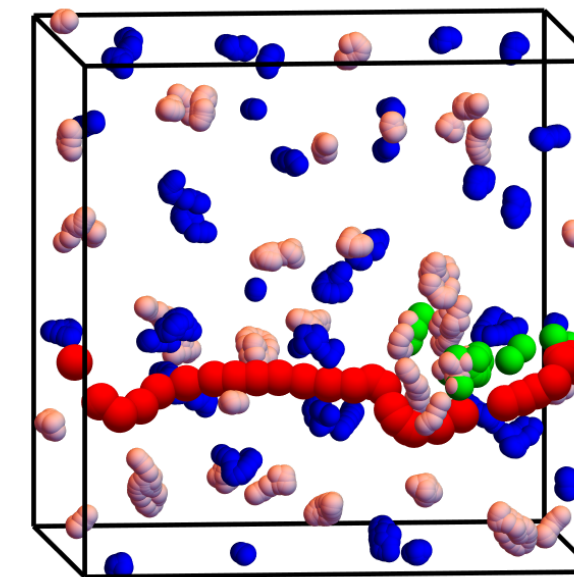
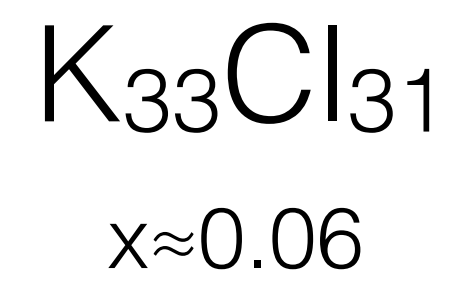
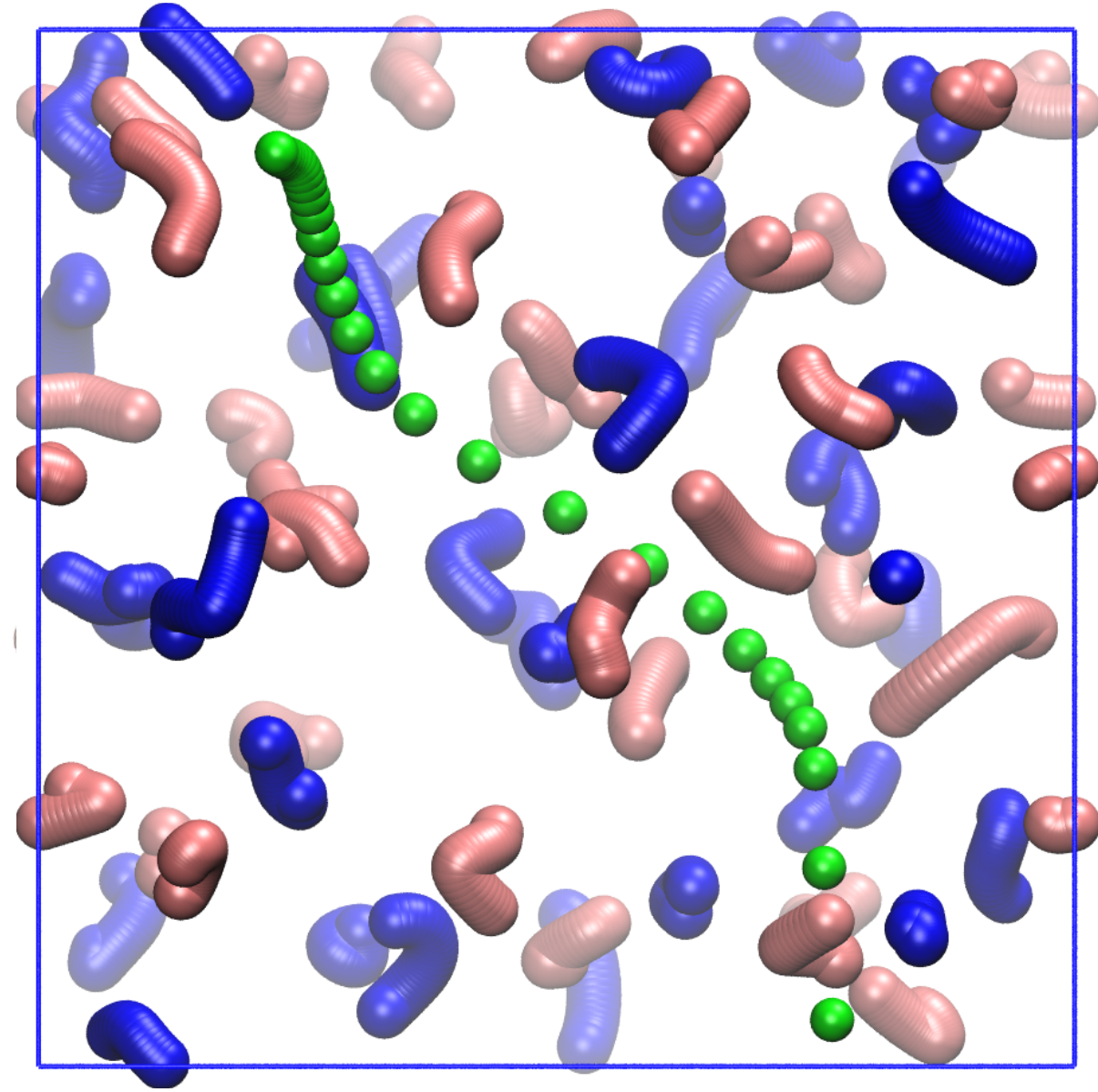
$$x \approx 0.06$$

non-stoichiometric melts

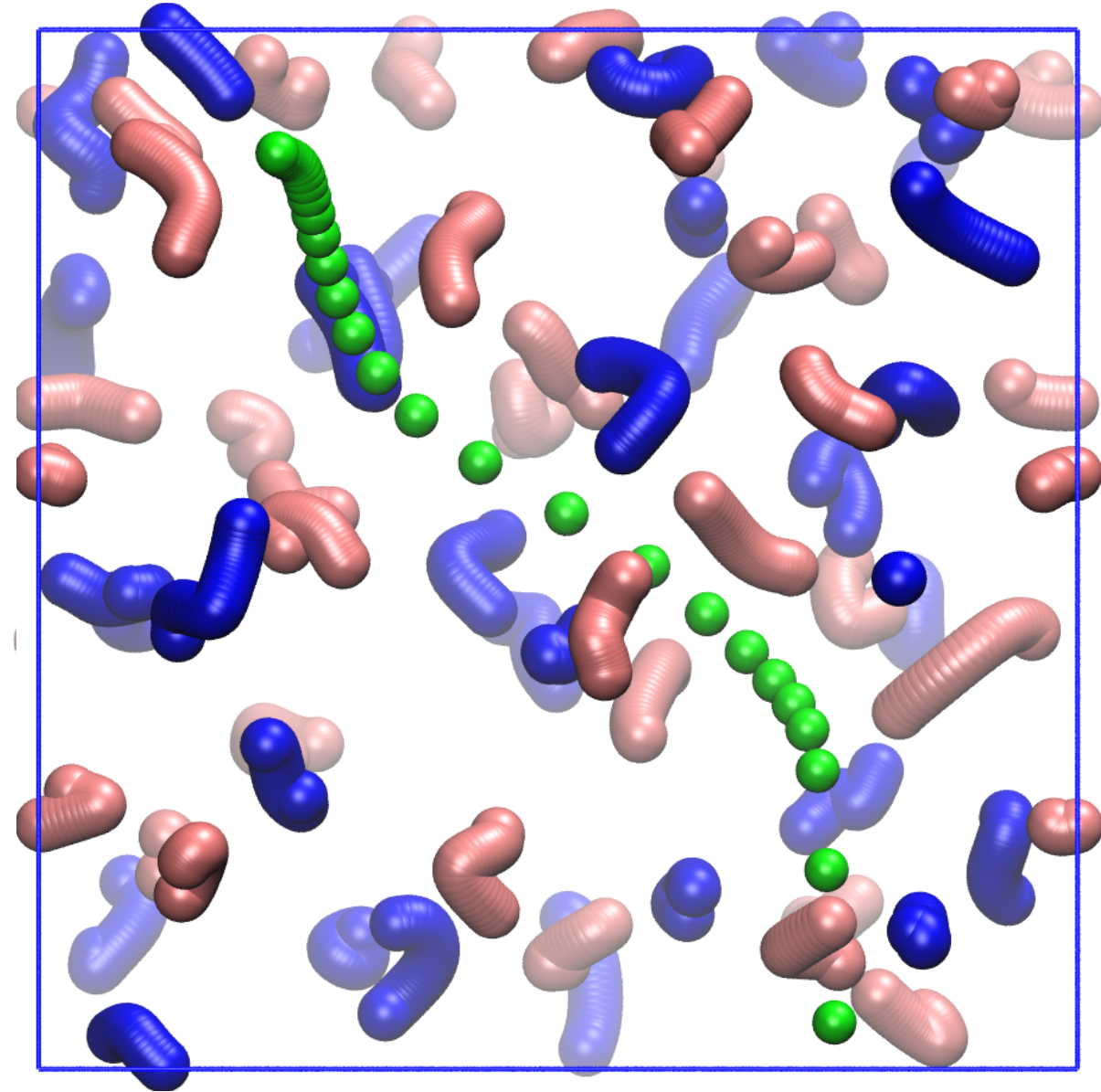


$$x \approx 0.06$$

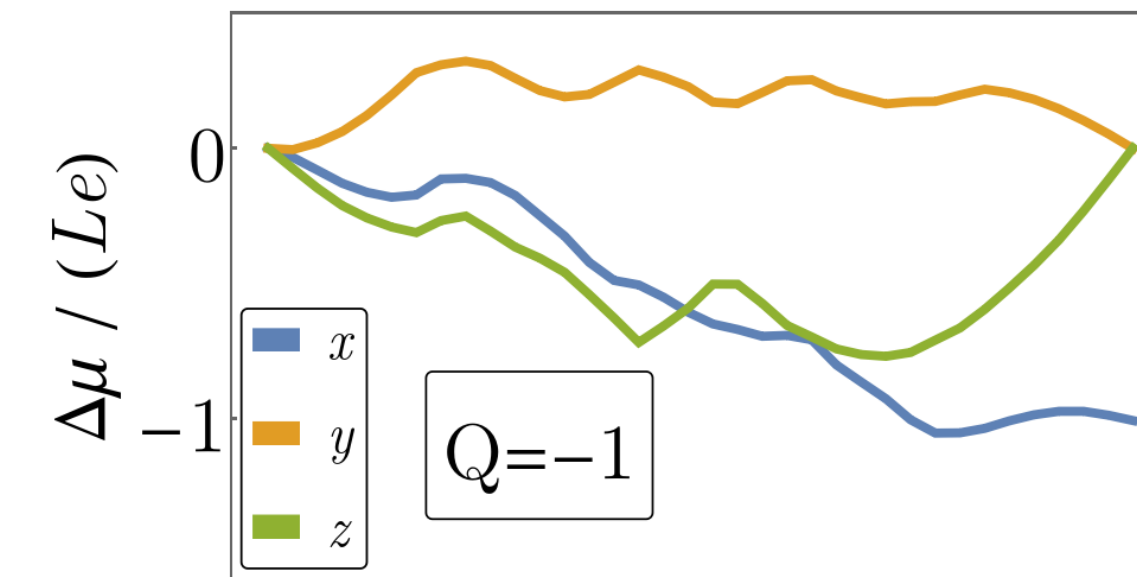
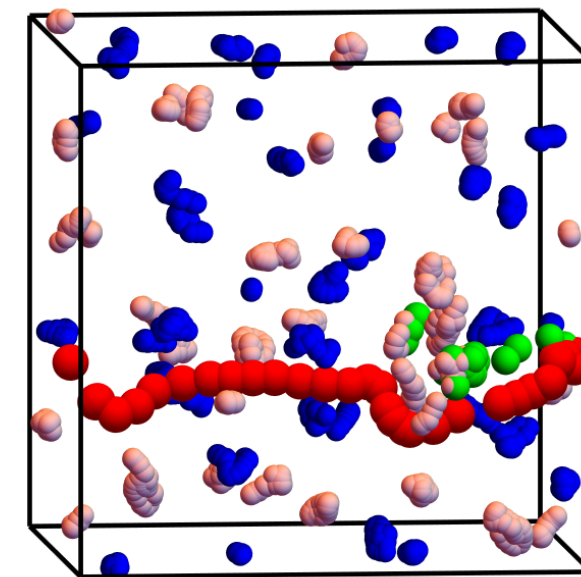
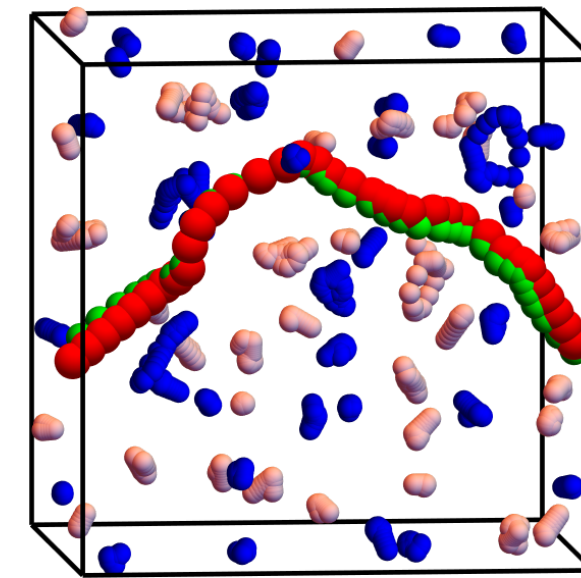
non-stoichiometric melts



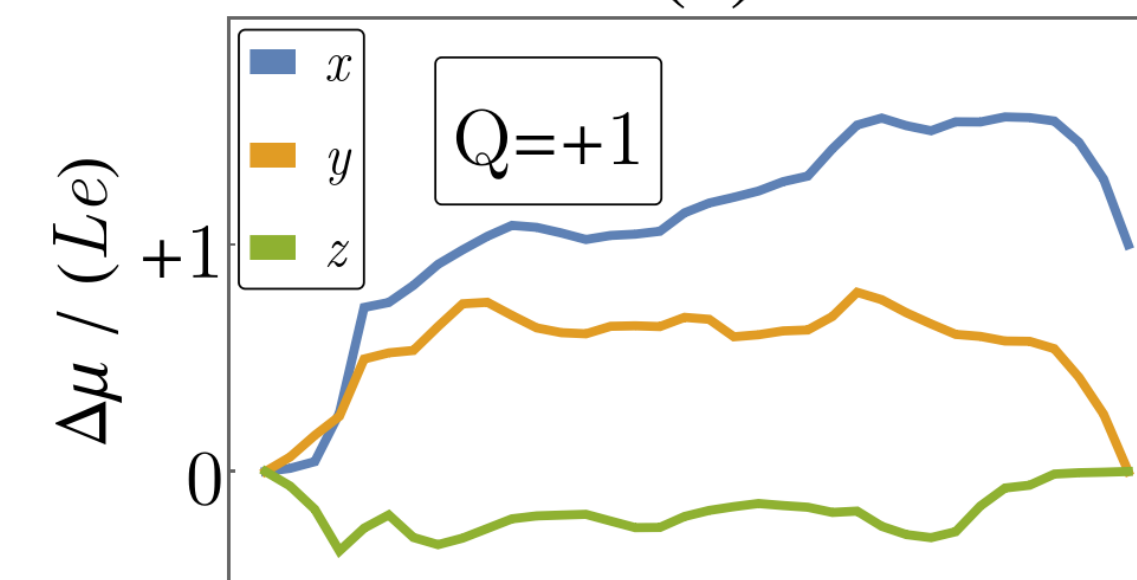
non-stoichiometric melts



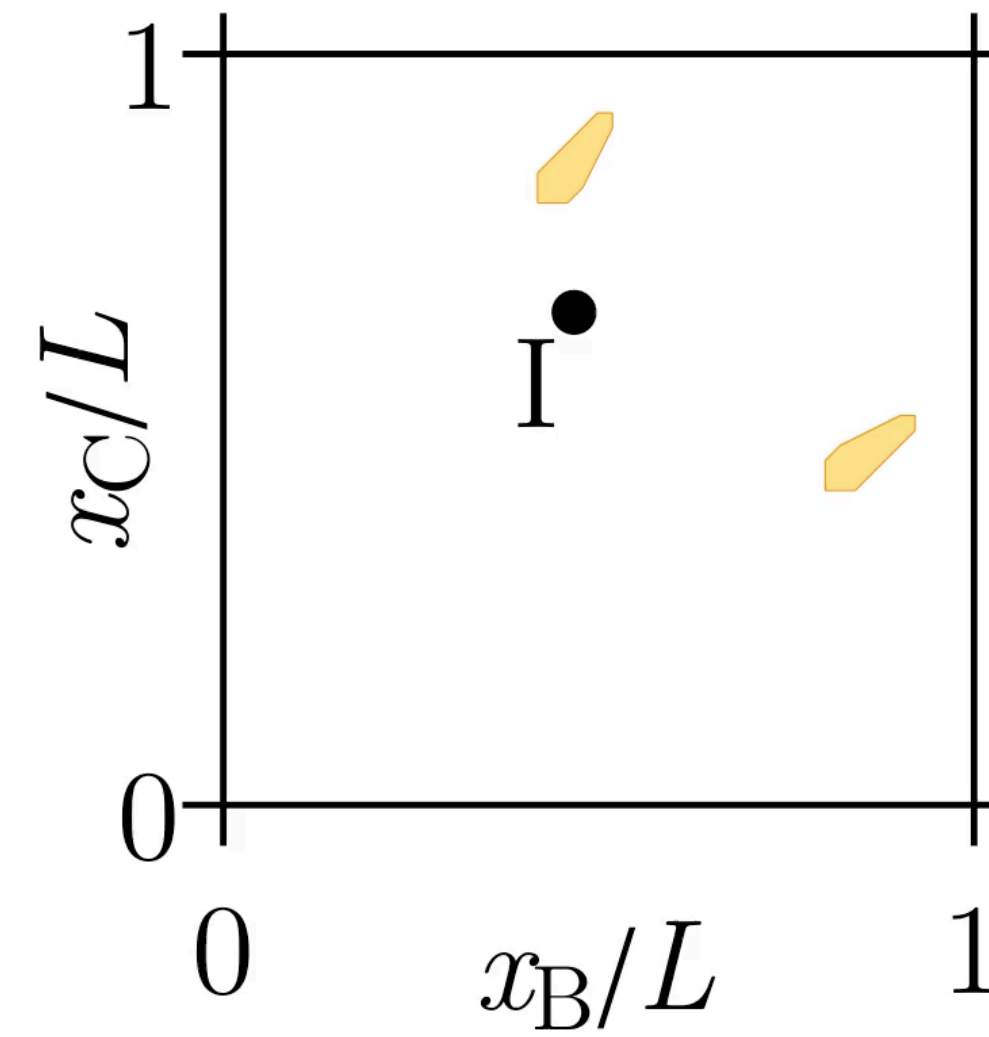
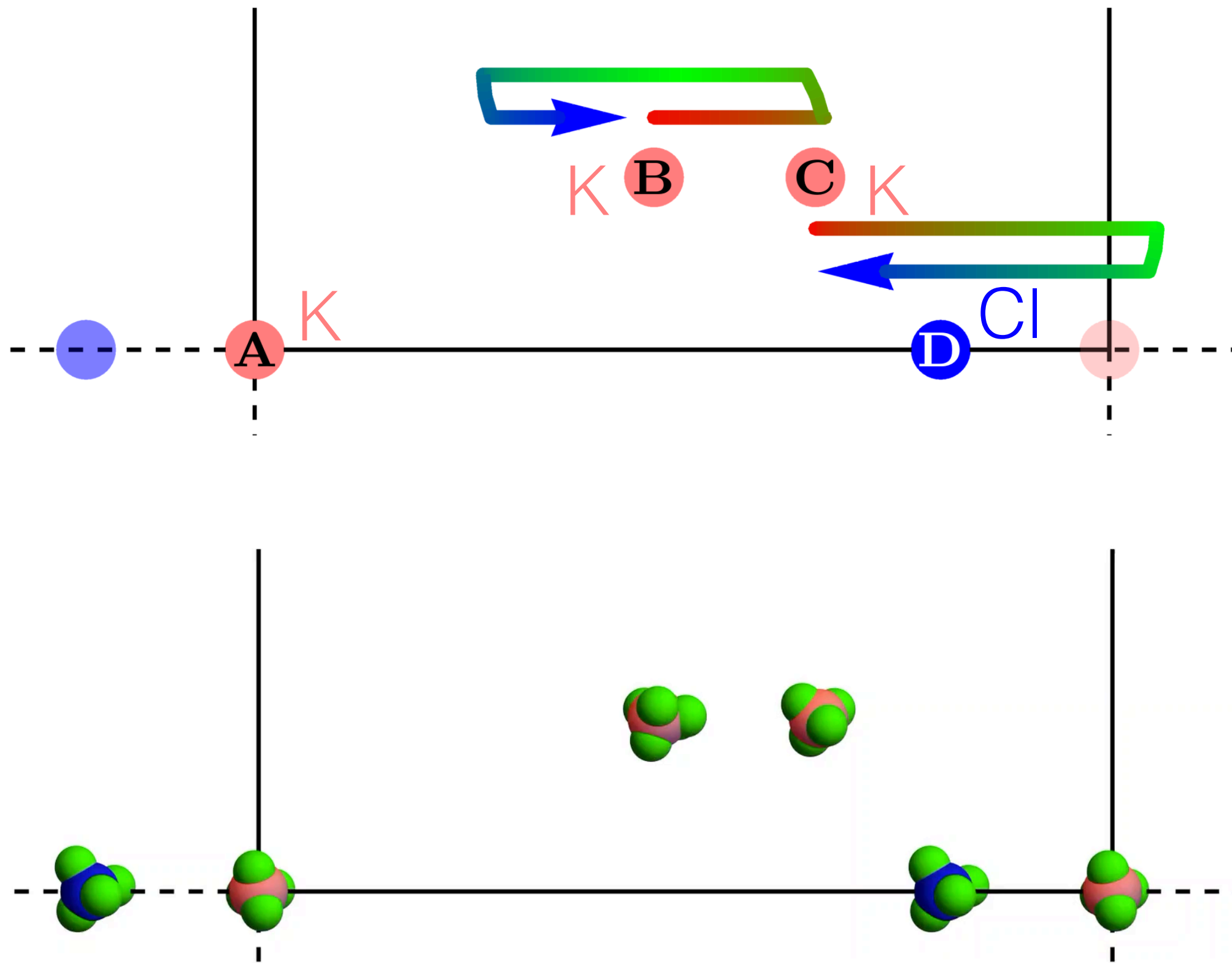
$$x \approx 0.06$$



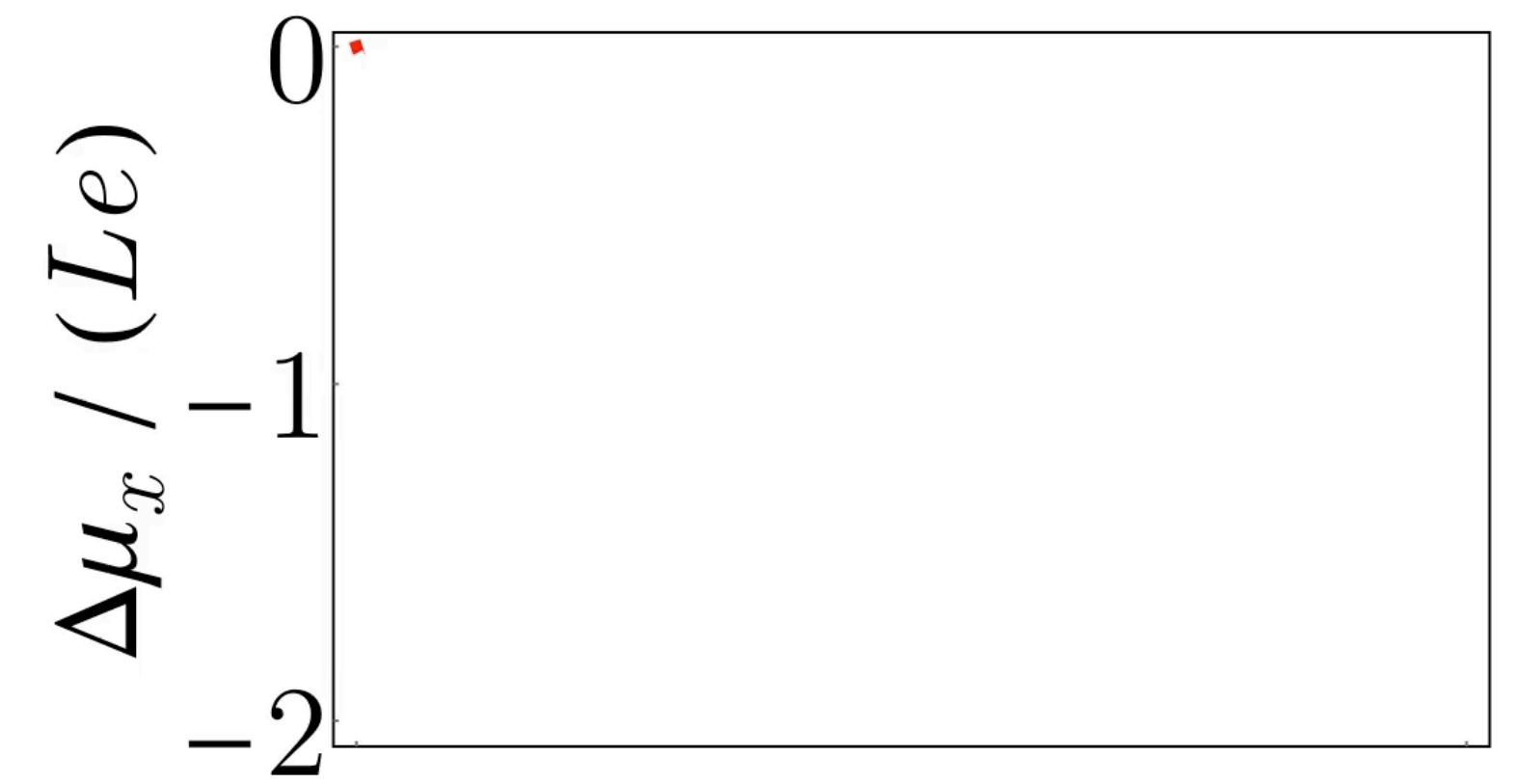
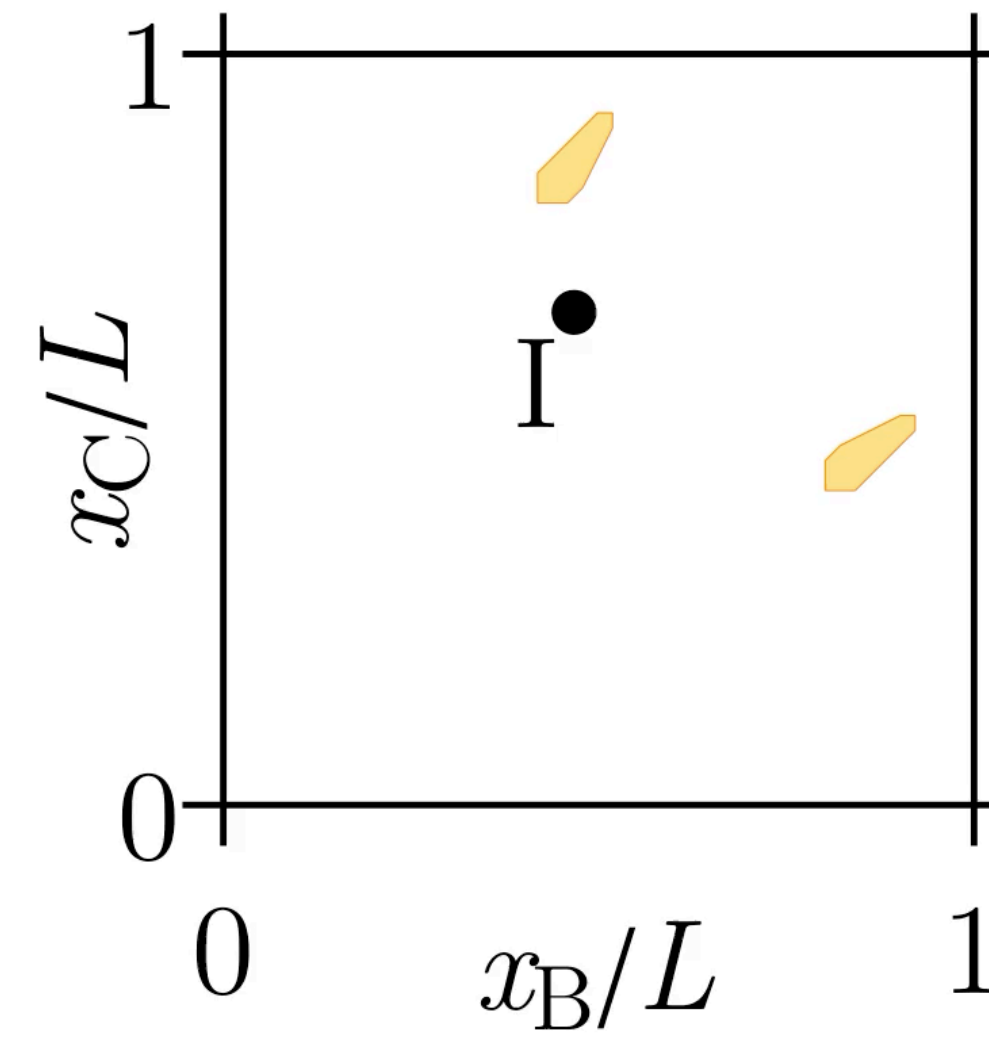
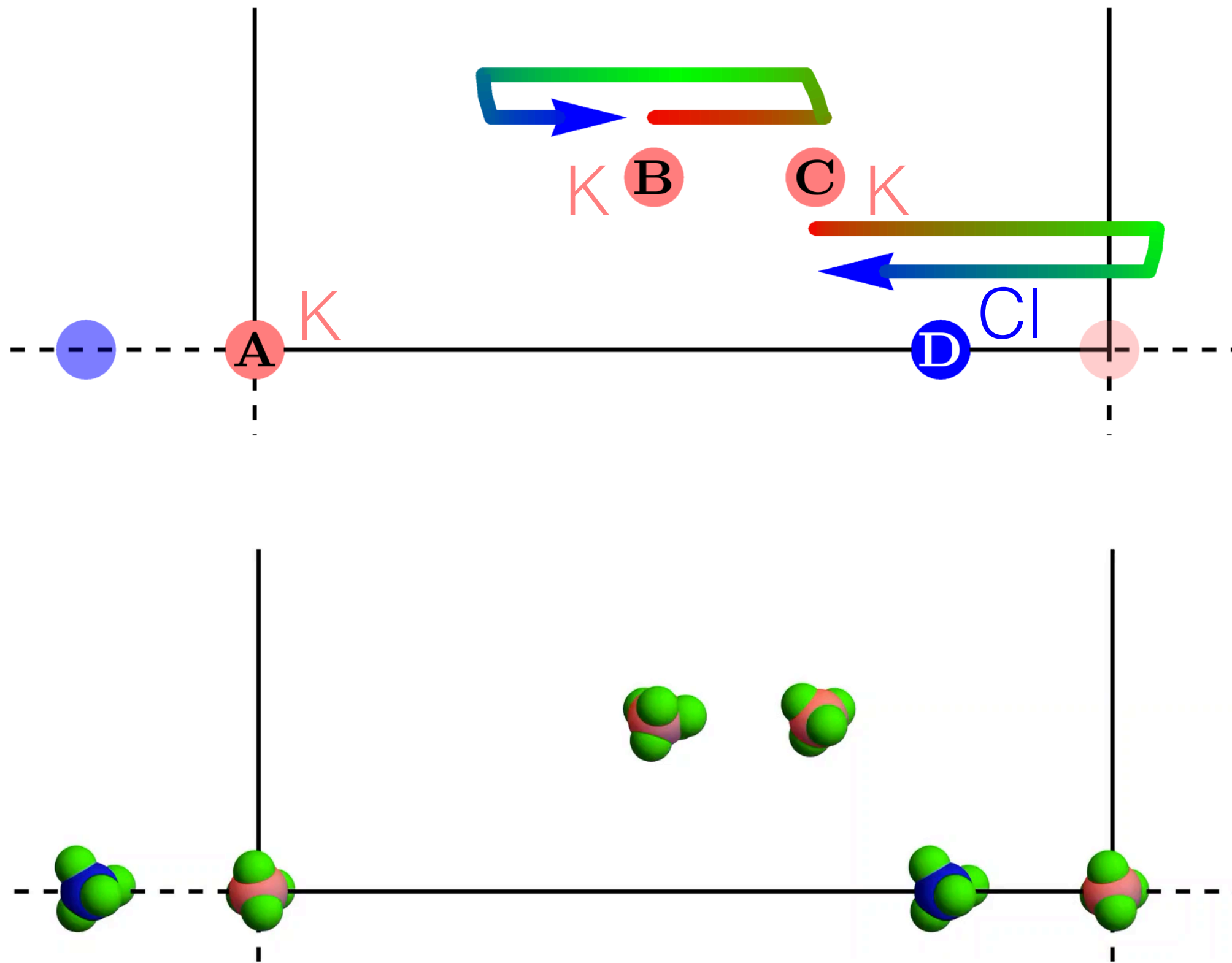
(b)



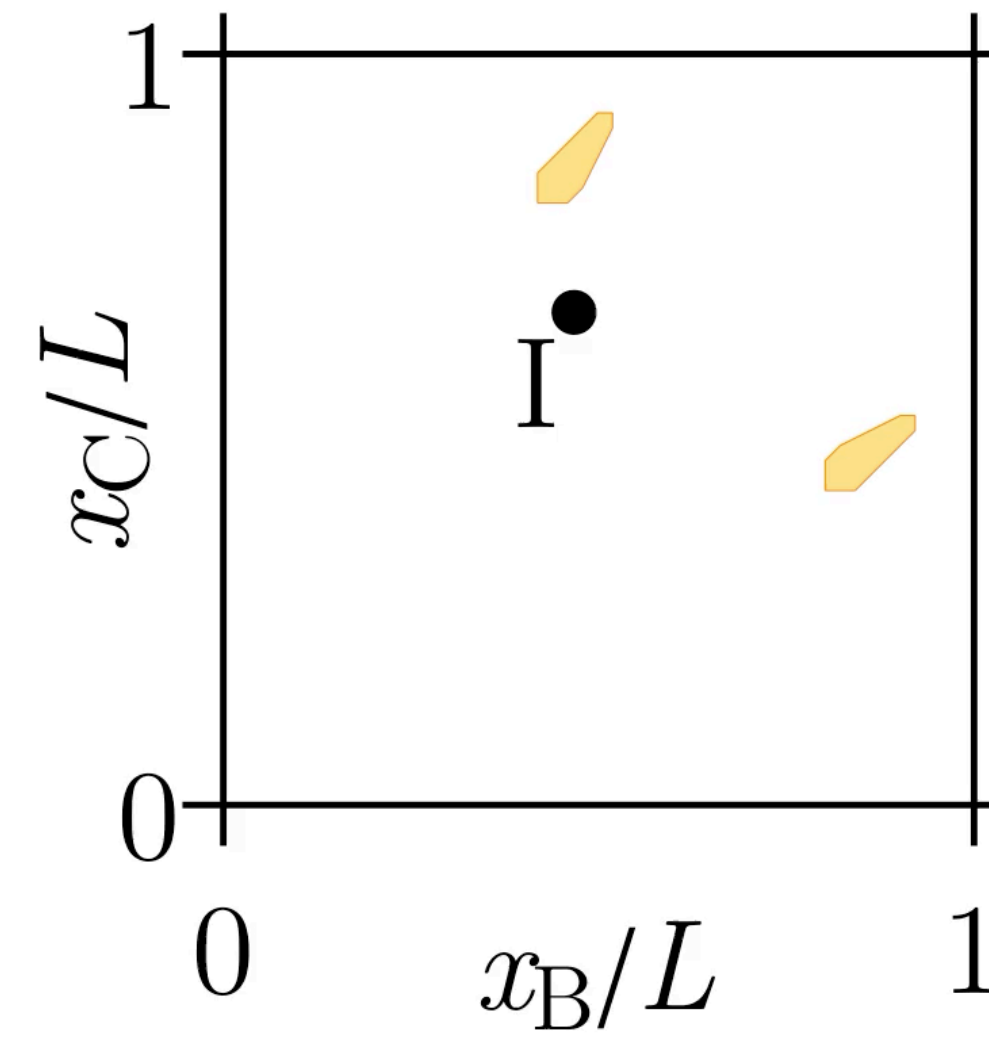
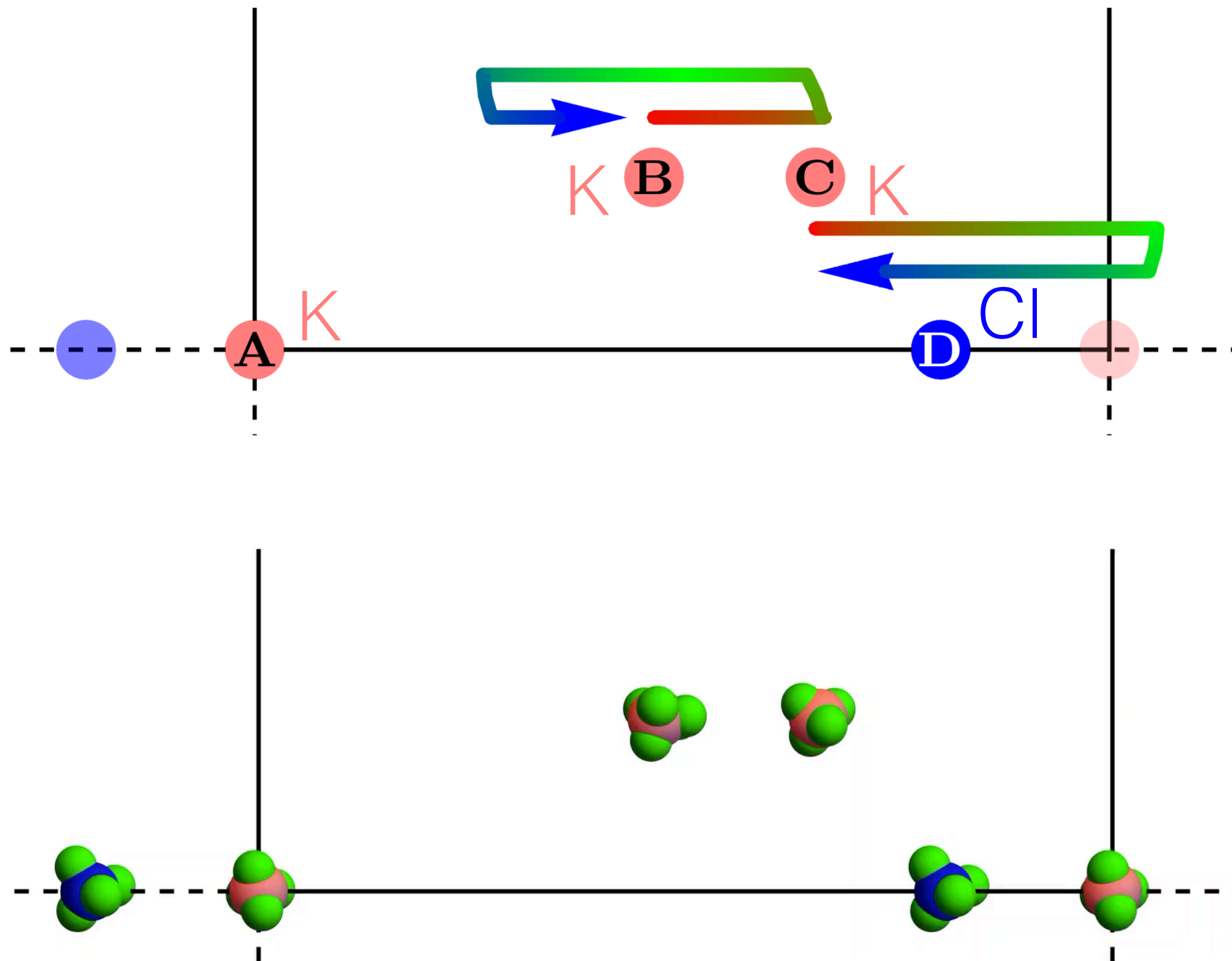
non-trivial particle transport



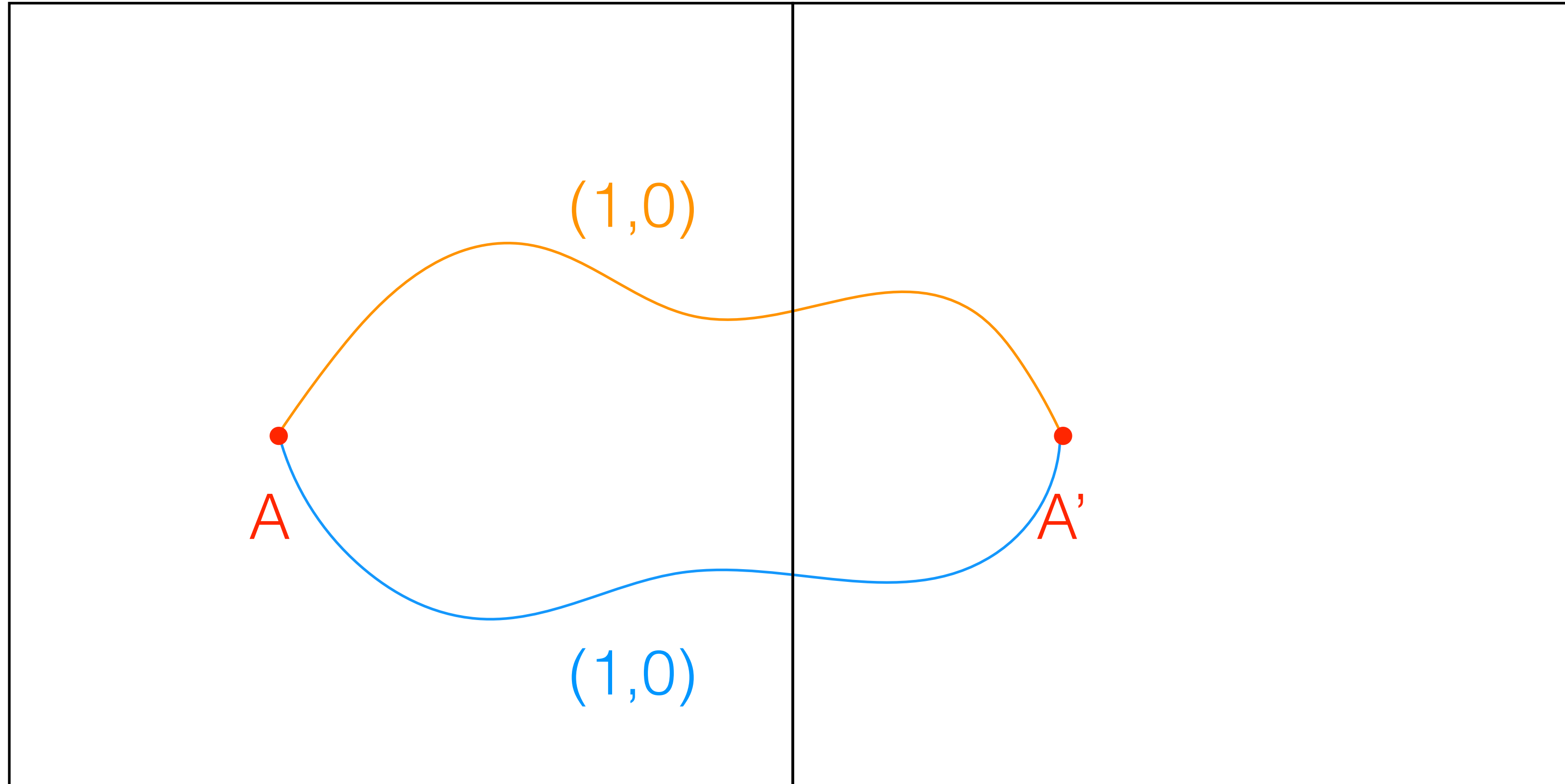
non-trivial particle transport



non-trivial particle transport

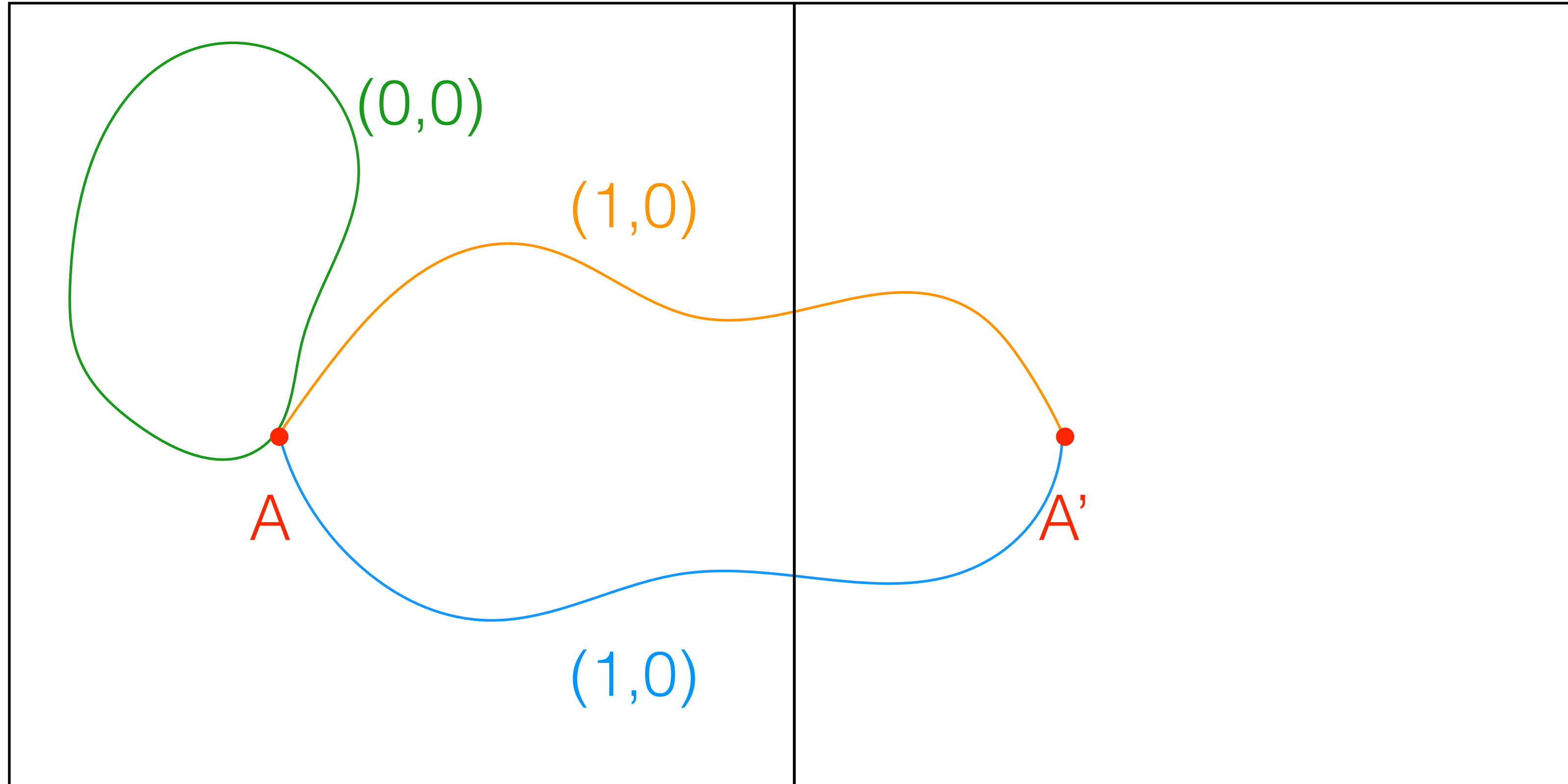


breach of strong adiabaticity



$$\mu = \mu^*$$

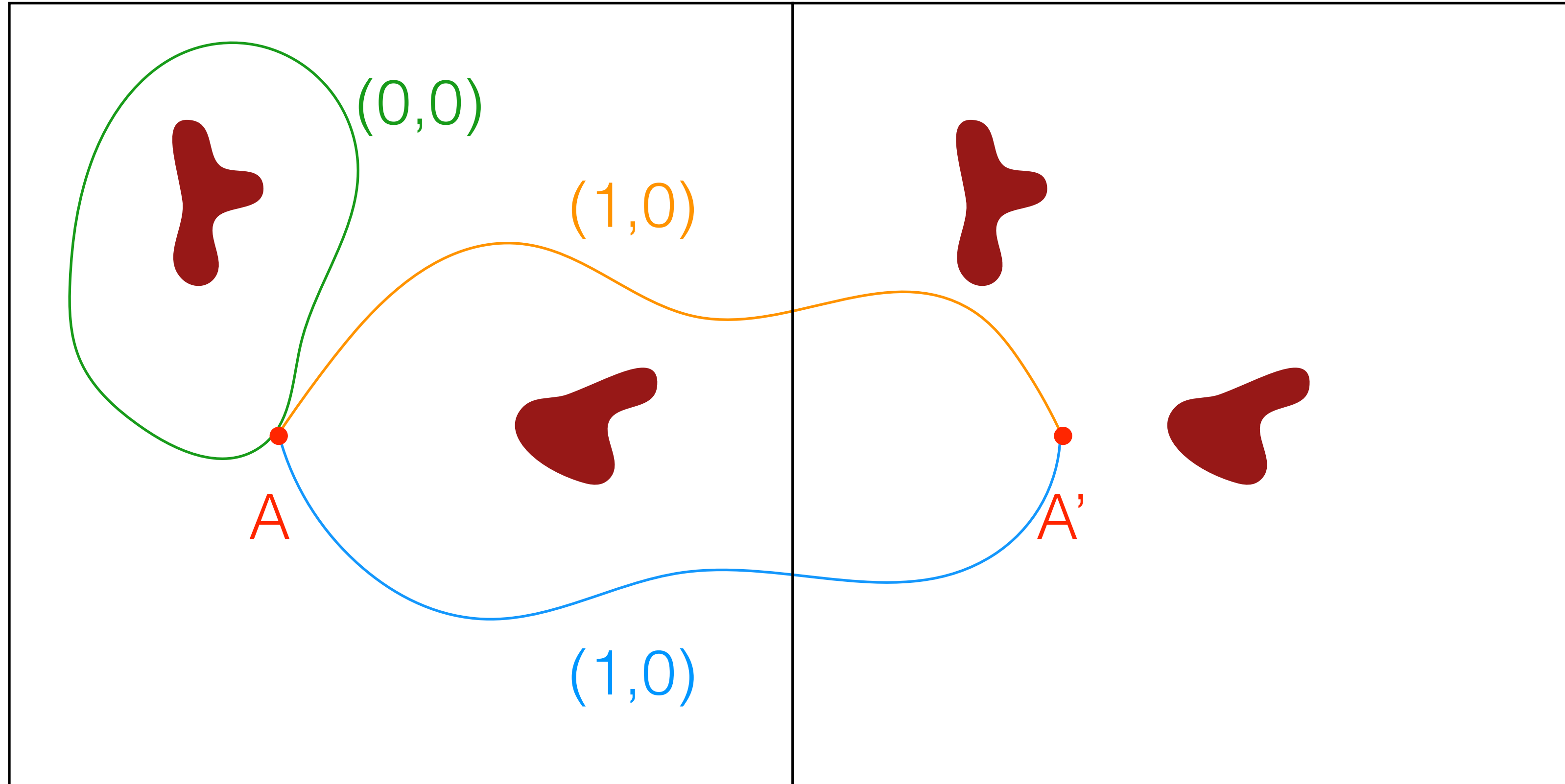
breach of strong adiabaticity



$$\mu = \mu^*$$

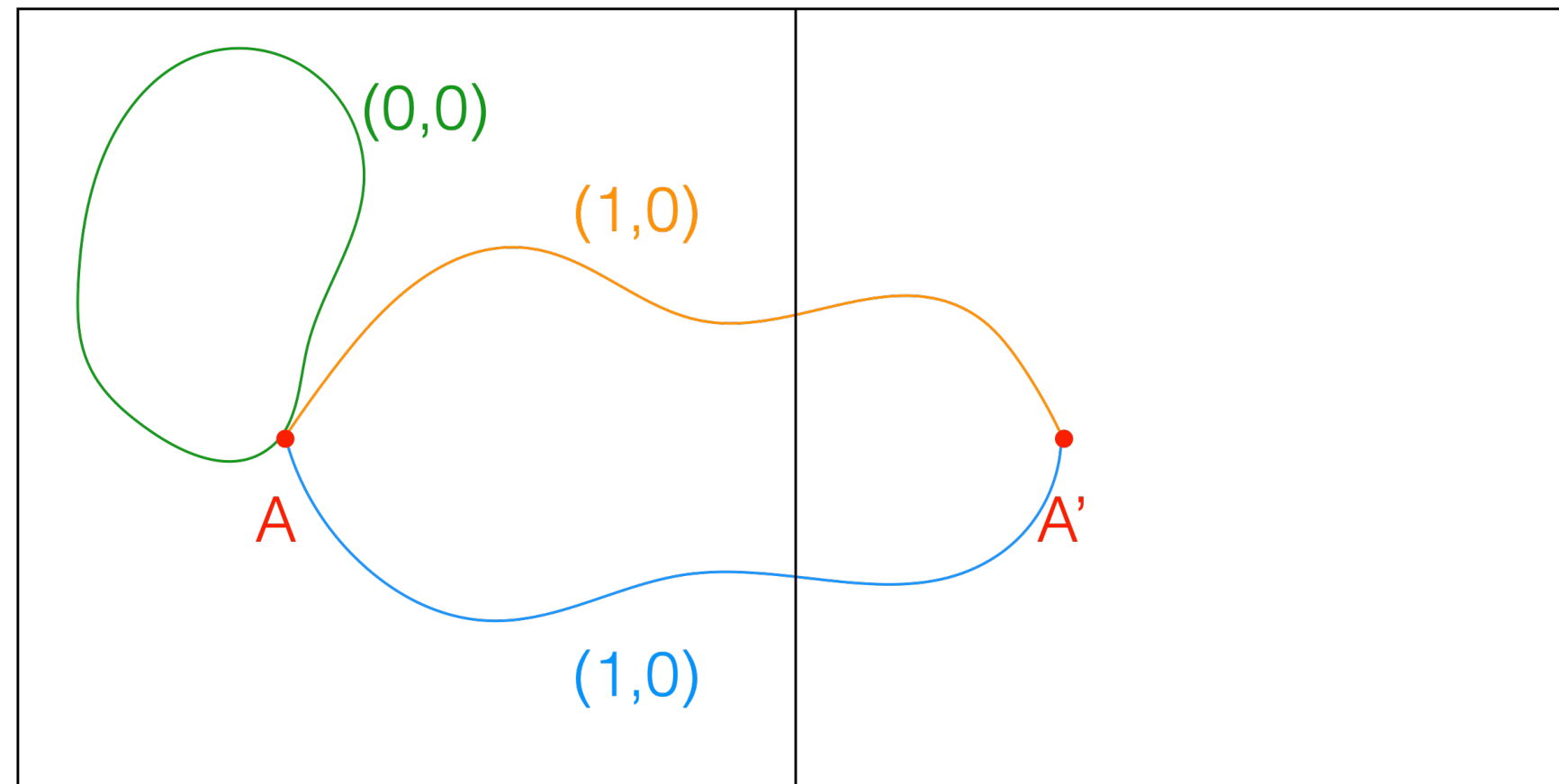
$$\mu = 0$$

breach of strong adiabaticity



$$\begin{aligned} \mu &\neq \mu^* \\ \mu &\neq 0 \end{aligned}$$

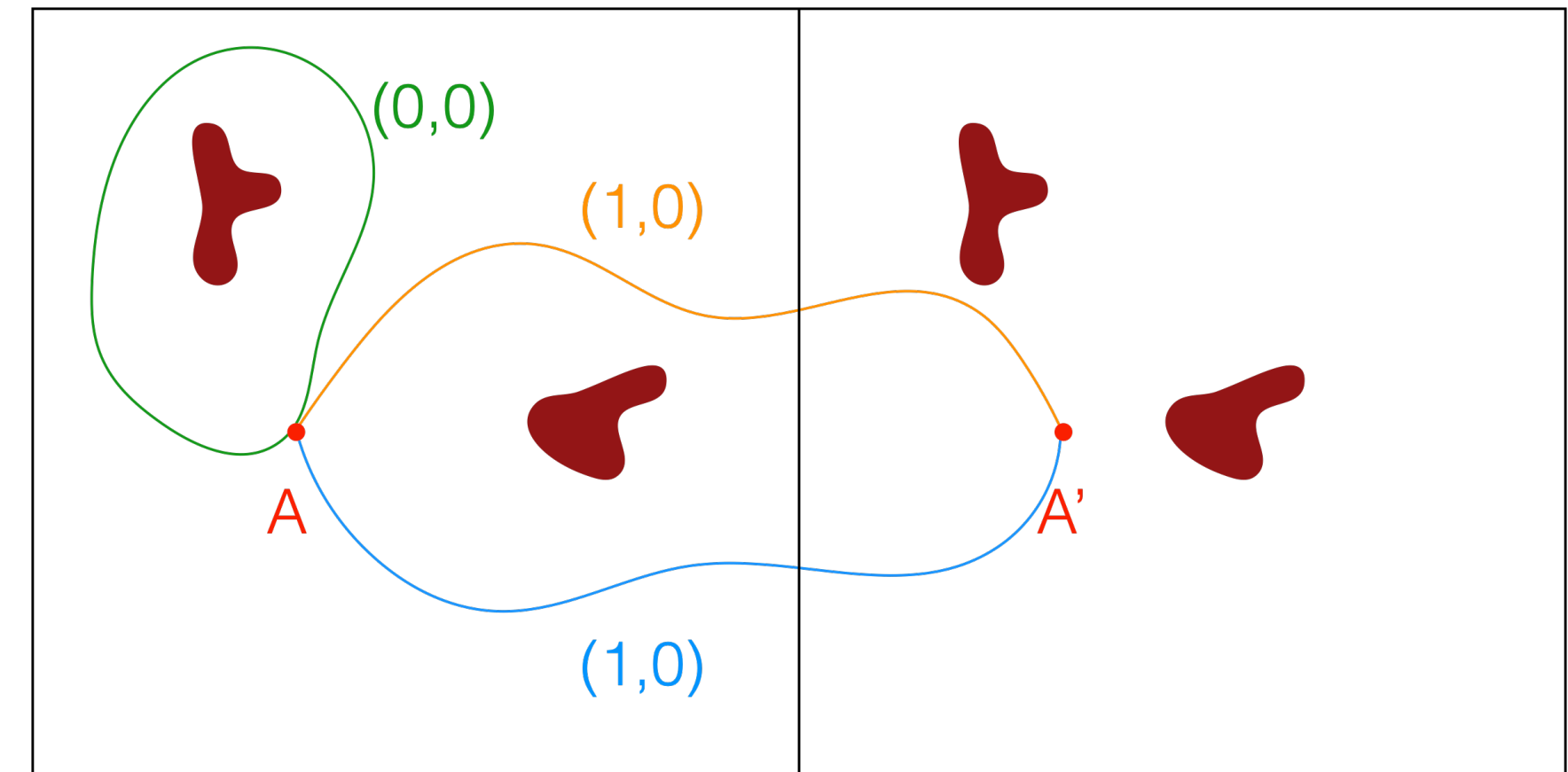
strongly adiabatic transport



$$\begin{aligned} \mu &= \mu^* \\ \mu &= 0 \end{aligned}$$



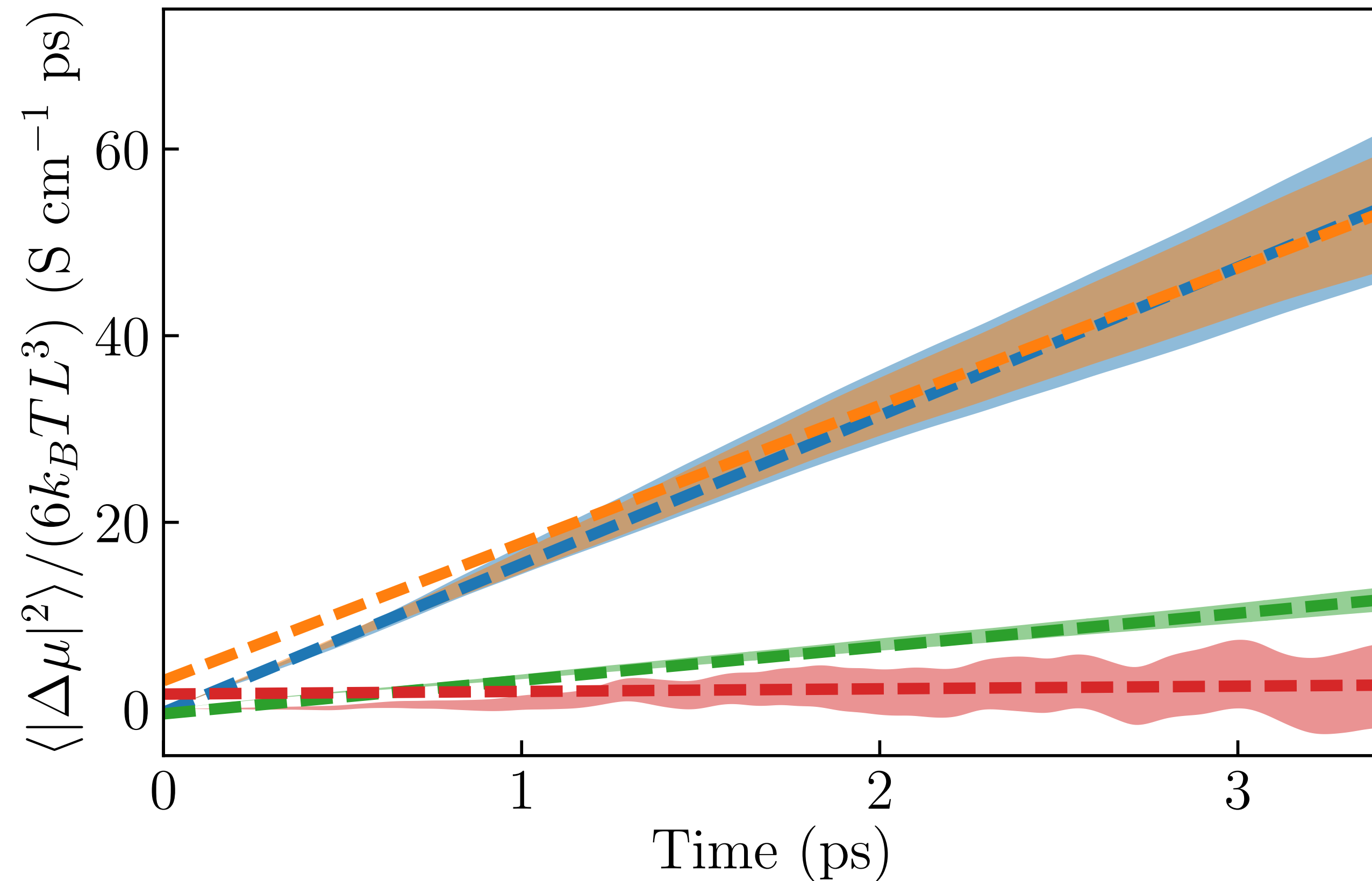
weakly adiabatic transport



$$\begin{aligned} \mu &\neq \mu^* \\ \mu &\neq 0 \end{aligned}$$



not trivial weakly adiabatic conductivity



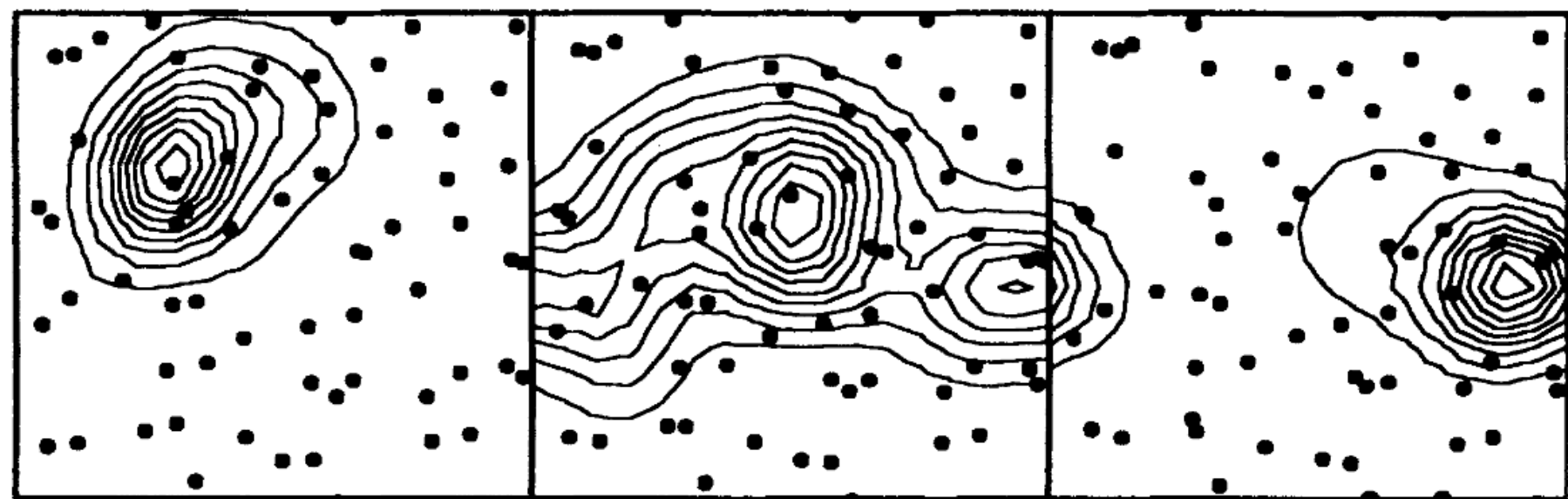
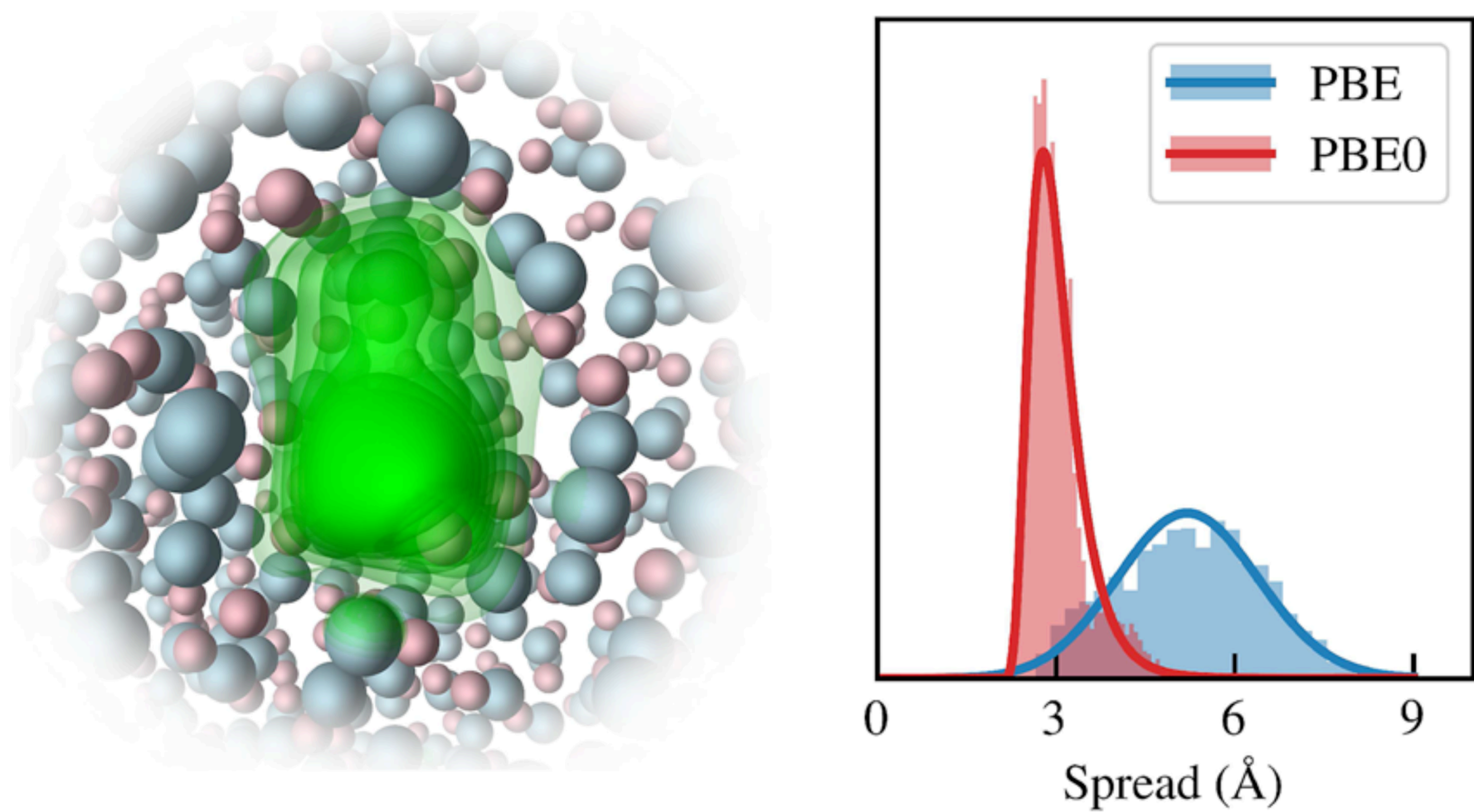
$$\Delta\boldsymbol{\mu} = e \int_0^t \mathbf{J}(t') dt'$$

$$J_\alpha(t) = \sum_{i\beta} Z_{i\alpha\beta}^*(t) v_{i\beta}(t)$$

$$J_\alpha(t) = \sum_i q_{S(i)} v_{i\alpha}(t) - 2v_\alpha^{lp}(t)$$

cross term

effects of self-interactions



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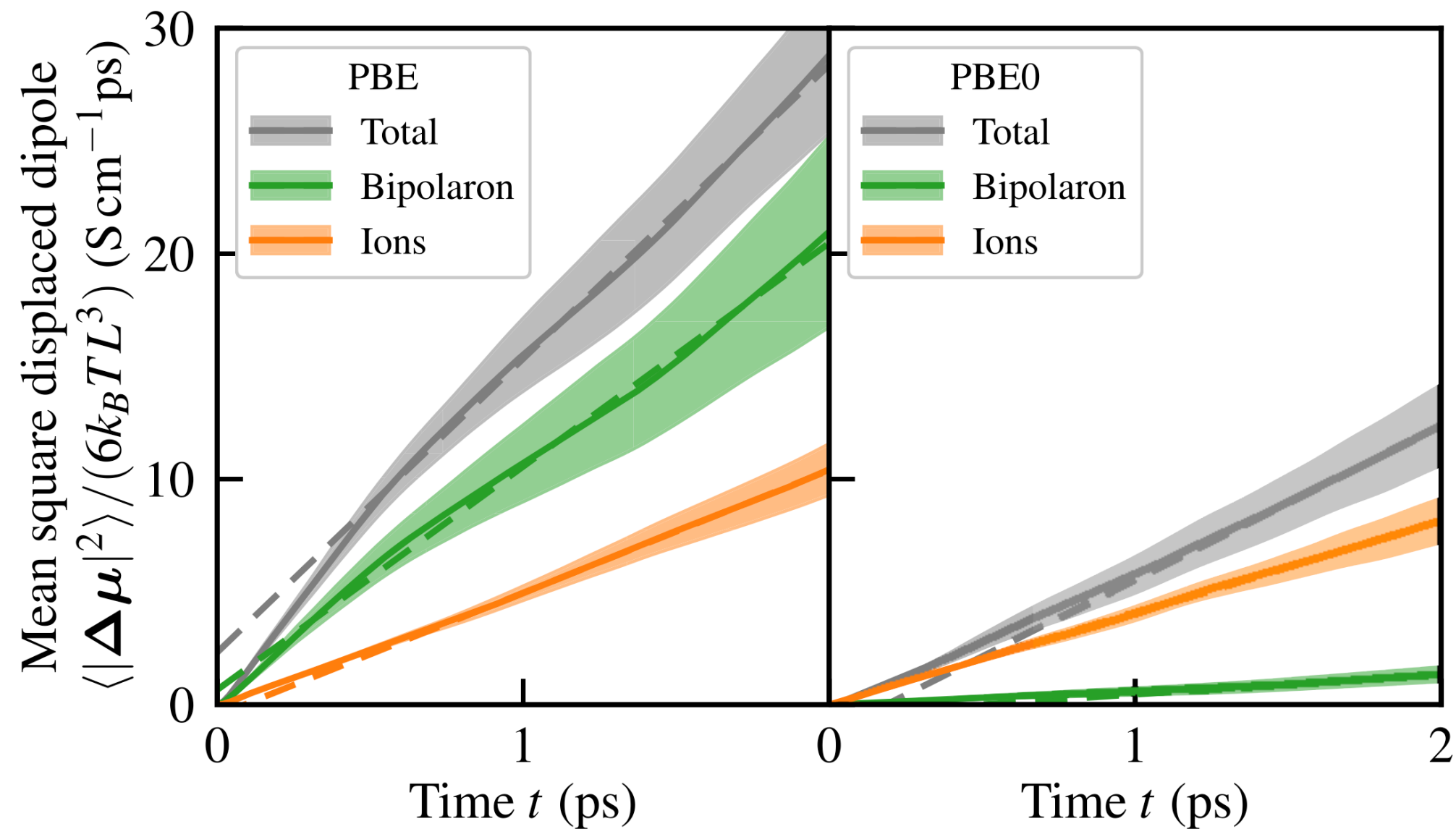
Localization, Hopping, and Diffusion of Electrons in Molten Salts

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IBM European Center for Scientific and Engineering Computing, Roma, Italy

and

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International School for Advanced Studies, Trieste, Italy
(Received 24 April 1986; revised manuscript received 3 April 1987)



conclusions



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- Topological quantization of adiabatic charge transport provides a rigorous definition of atomic oxidation states.;

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- In non-stoichiometric electrolytes, a breakdown of strong adiabaticity gives rise to an anomalous transport regime—intermediate between metallic and ionic—where charge can be transferred without accompanying mass displacement;

conclusions

- Topological quantization of adiabatic charge transport provides a rigorous definition of atomic oxidation states.;
- Gauge invariance and the quantization of charge transport imply that the electric conductivity of stoichiometric electrolytes depends on the formal oxidation numbers of the ionic species, as expressed through the Green-Kubo relation;
- In non-stoichiometric electrolytes, a breakdown of strong adiabaticity gives rise to an anomalous transport regime—intermediate between metallic and ionic—where charge can be transferred without accompanying mass displacement;
- the degree of electron localization in the ionic broth is strongly influenced by self-interaction effects, which are treated differently by various exchange-correlation functionals; the resulting electrone contribution to the conductivity varies accordingly.



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... we are not as young
as we used to be!

thanks to:



Federico Grasselli



Paolo Pegolo

Topological quantization and gauge invariance of charge transport in liquid insulators

Federico Grasselli¹ and Stefano Baroni^{1,2*}

PHYSICAL REVIEW X

Oxidation States, Thouless' Pumps, and Nontrivial Ionic Transport in Nonstoichiometric Electrolytes

Paolo Pegolo, Federico Grasselli, and Stefano Baroni
Phys. Rev. X **10**, 041031 – Published 12 November 2020



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Paolo Pegolo✉, Stefano Baroni✉, Federico Grasselli✉

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Self-interaction and transport of solvated electrons in molten salts

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Paolo Pegolo,^{1,a)} Stefano Baroni,^{1,2,b)} and Federico Grasselli^{3,c)}



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