

**DM@LHC WORKSHOP**  
**CHICAGO - 19 SEPT 2013**

# **VALIDITY OF EFT** **FOR DM @ LHC**

**ANDREA DE SIMONE**



**BASED ON:**

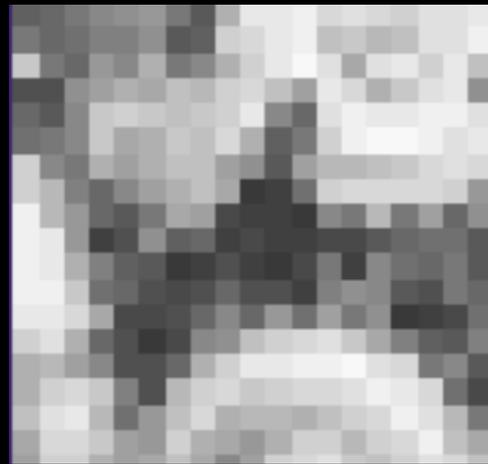
**BUSONI, DS, MORGANTE, RIOTTO - ARXIV:1307.2253**

## Searches for Dark Matter @ LHC

- **is the EFT approach reliable?**
- **dead end? improvements in sight?**

# EFFECTIVE FIELD THEORY DESCRIPTION

effective  
low-energy  
description



( $\Lambda \sim 1$  TeV)

New States

(say, 10 TeV)

$M_Z$

E

EFT OK

Integrate out the UV physics connecting DM-SM and describe interactions with eff. ops.:

$$\frac{1}{\Lambda^2} (\bar{\chi} \Gamma^A \chi) (\bar{q} \Gamma_A q)$$

LHC can access regions **beyond** the validity of the eff. description



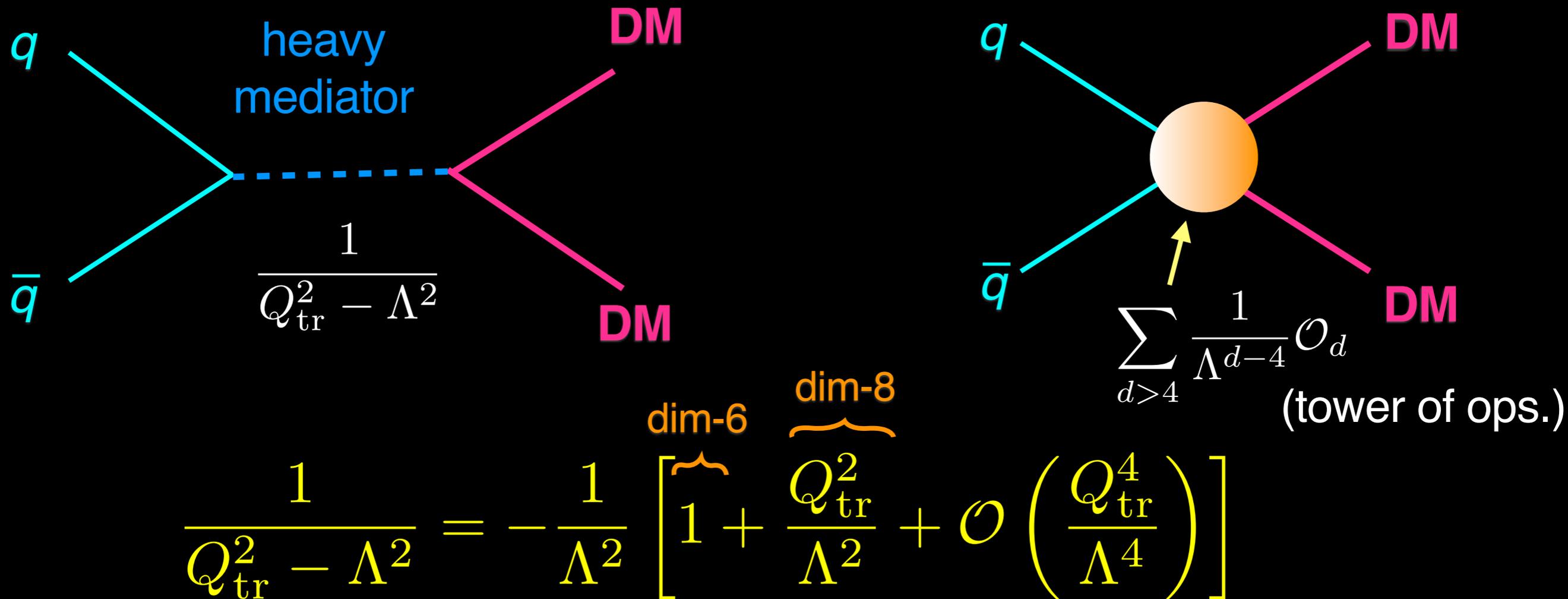
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→ need to use EFT carefully and consistently

# EFFECTIVE FIELD THEORY DESCRIPTION

- the **momentum transfer** in the relevant process must be  $Q_{\text{tr}} \lesssim \Lambda$



- $Q_{\text{tr}}/\Lambda$  measures the badness of the truncation of the tower of effective ops to the lowest dimensional ones
- Usually, lowest order is OK. Not a problem for **direct/indirect** searches. Situation can be different **@ LHC**.

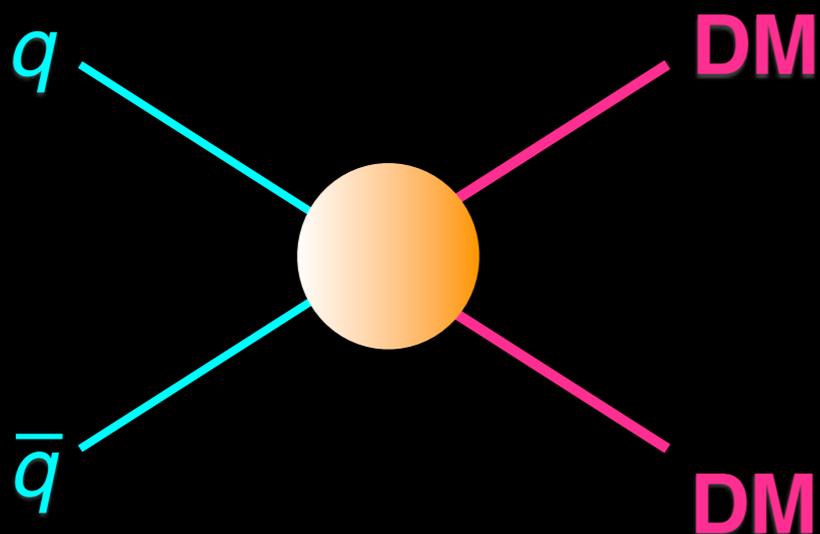
## Standard lore

mediator mass  $M > m_\chi$

$$\Lambda \simeq \frac{M}{\sqrt{g_{\text{SM}} g_\chi}} \gtrsim \frac{M}{4\pi} \quad \Lambda \gtrsim \frac{m_\chi}{4\pi}$$

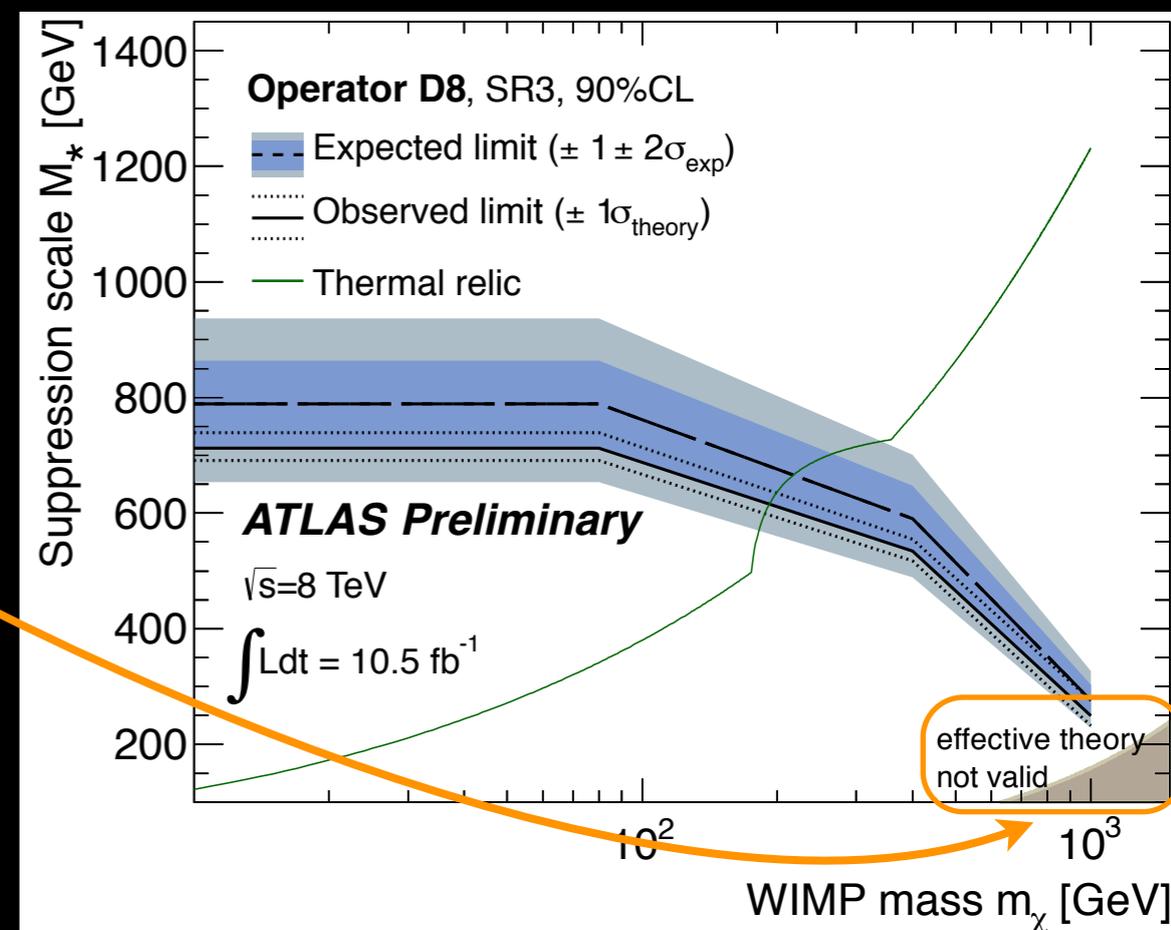
Actual limits can be **stronger** (depending on the process)

## 2 → 2 process



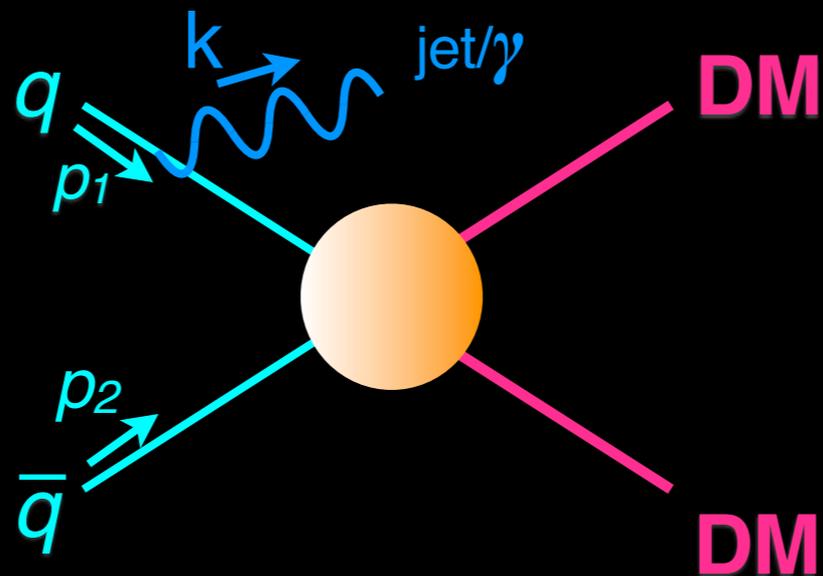
$$Q_{\text{tr}}^2 \geq 4m_\chi^2 \longrightarrow \Lambda > 2m_\chi$$

below this bound, the contribution of higher-dim ops becomes important



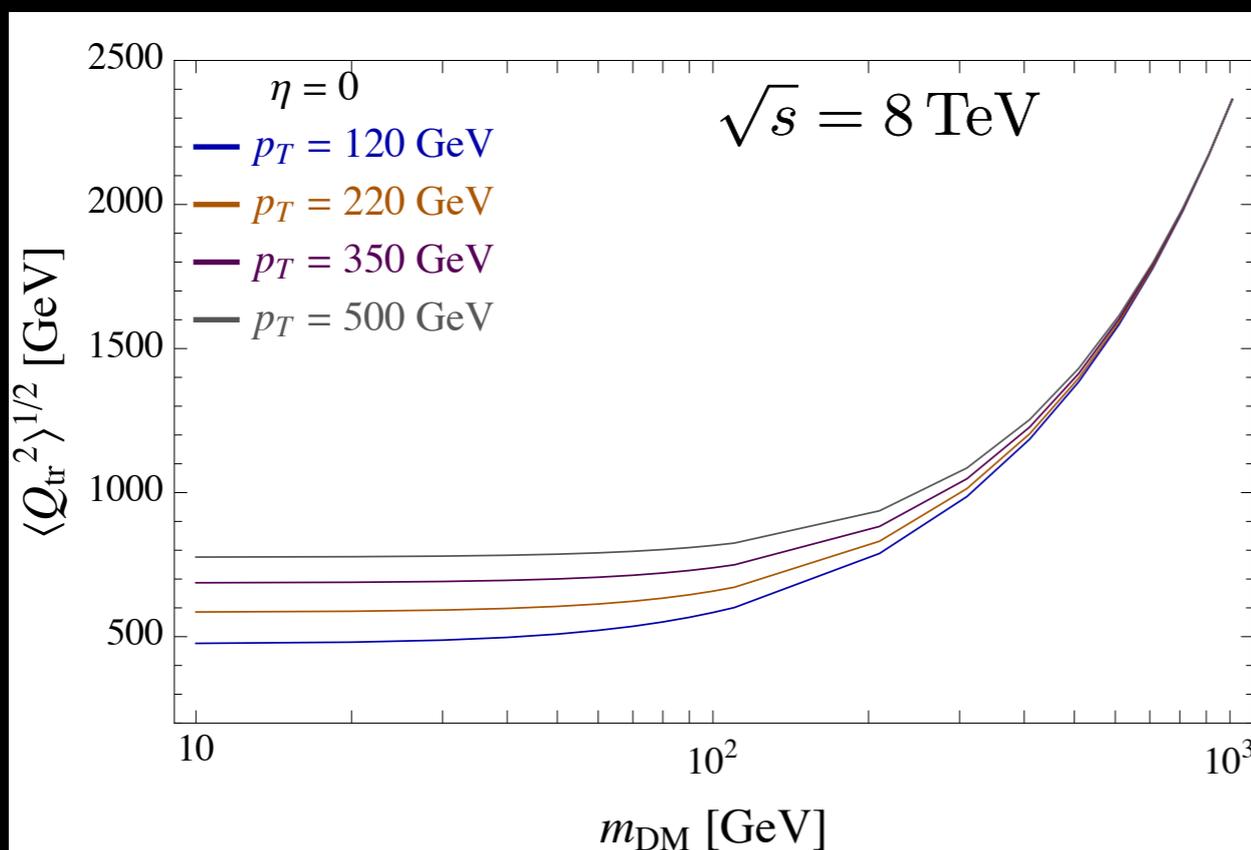
[ATLAS-CONF-2012-147]

## 2 → 3 process



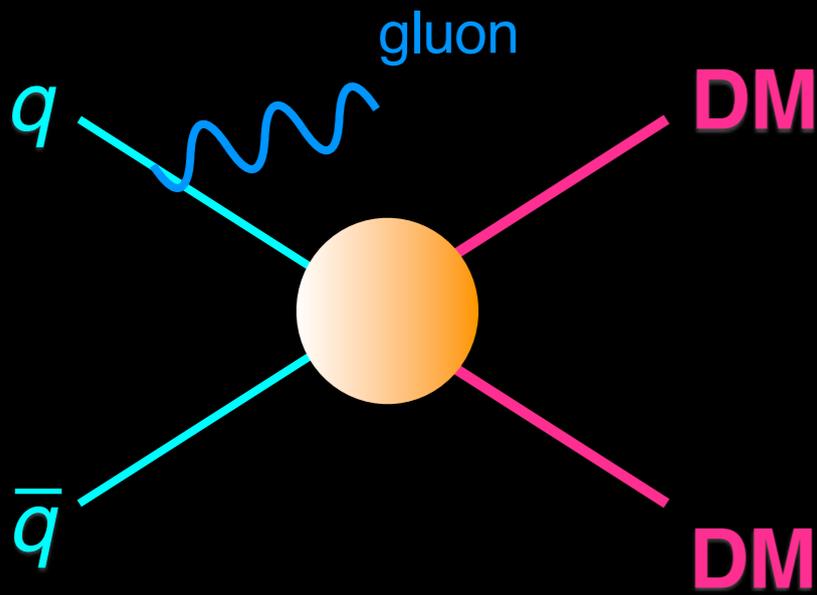
$$Q_{\text{tr}}^2 \equiv (p_1 + p_2 - k)^2 = x_1 x_2 s - \sqrt{s} p_T (x_1 e^{-\eta} + x_2 e^{\eta})$$

for s-channel momentum transfer



the momentum transfer is larger for larger DM mass

# AN EXAMPLE WITH SCALAR MEDIATOR

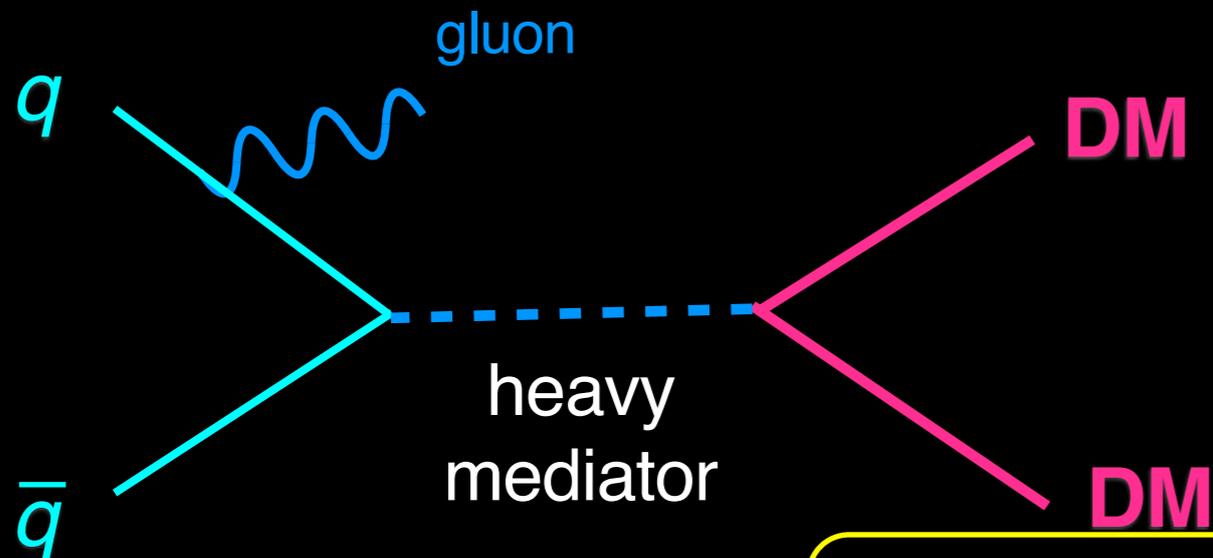


eff. operator  $\mathcal{O}_S = \frac{1}{\Lambda^2} (\bar{\chi}\chi)(\bar{q}q)$

parton-level differential cross section:

$$\frac{d^2\hat{\sigma}_{\text{eff}}}{dp_T d\eta} = \frac{\alpha_s}{36\pi^2} \frac{1}{p_T} \frac{1}{\Lambda^4} \frac{[Q_{\text{tr}}^2 - 4m_{\text{DM}}^2]^{3/2}}{Q_{\text{tr}}} \left[ 1 + \frac{Q_{\text{tr}}^4}{(x_1 x_2 s)^2} \right]$$

matching:  $\frac{1}{\Lambda^2} = \frac{g_\chi g_q}{M^2}$



$$\mathcal{L}_{\text{UV}} \supset \frac{1}{2} M^2 S^2 - g_q \bar{q}q S - g_\chi \bar{\chi}\chi S$$

parton-level differential cross section:

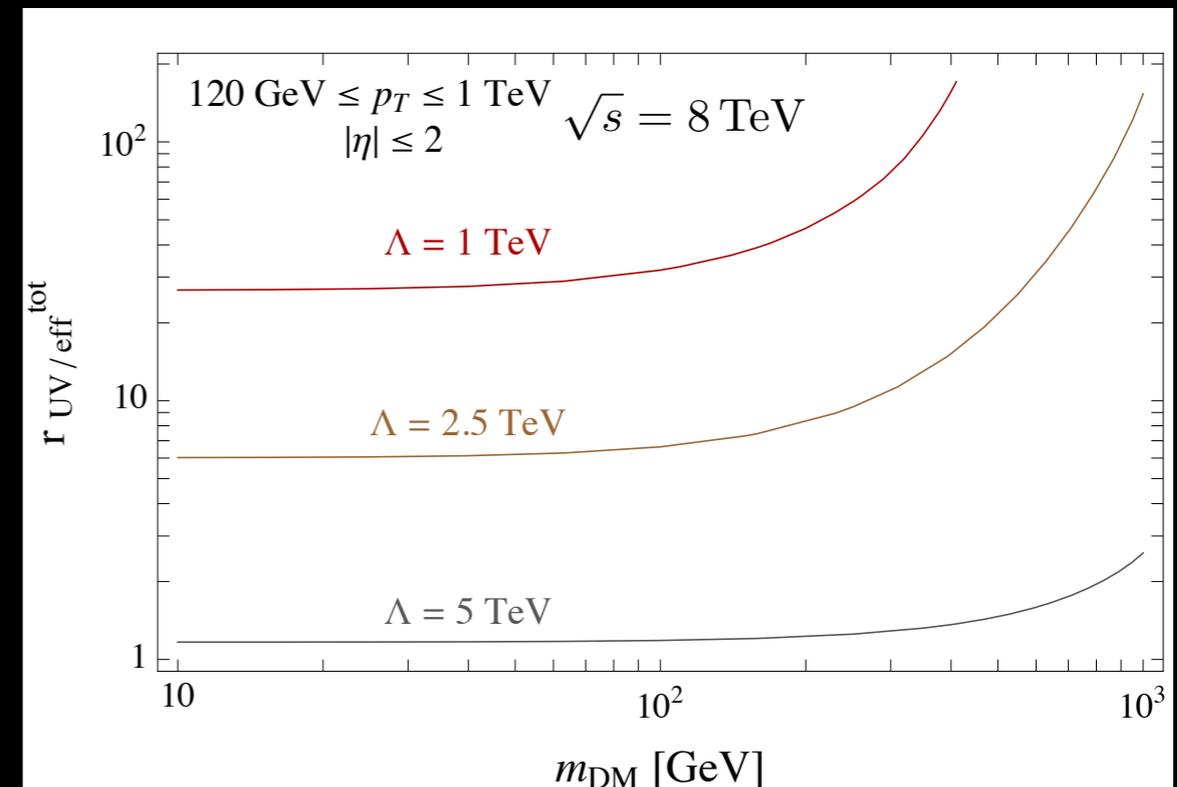
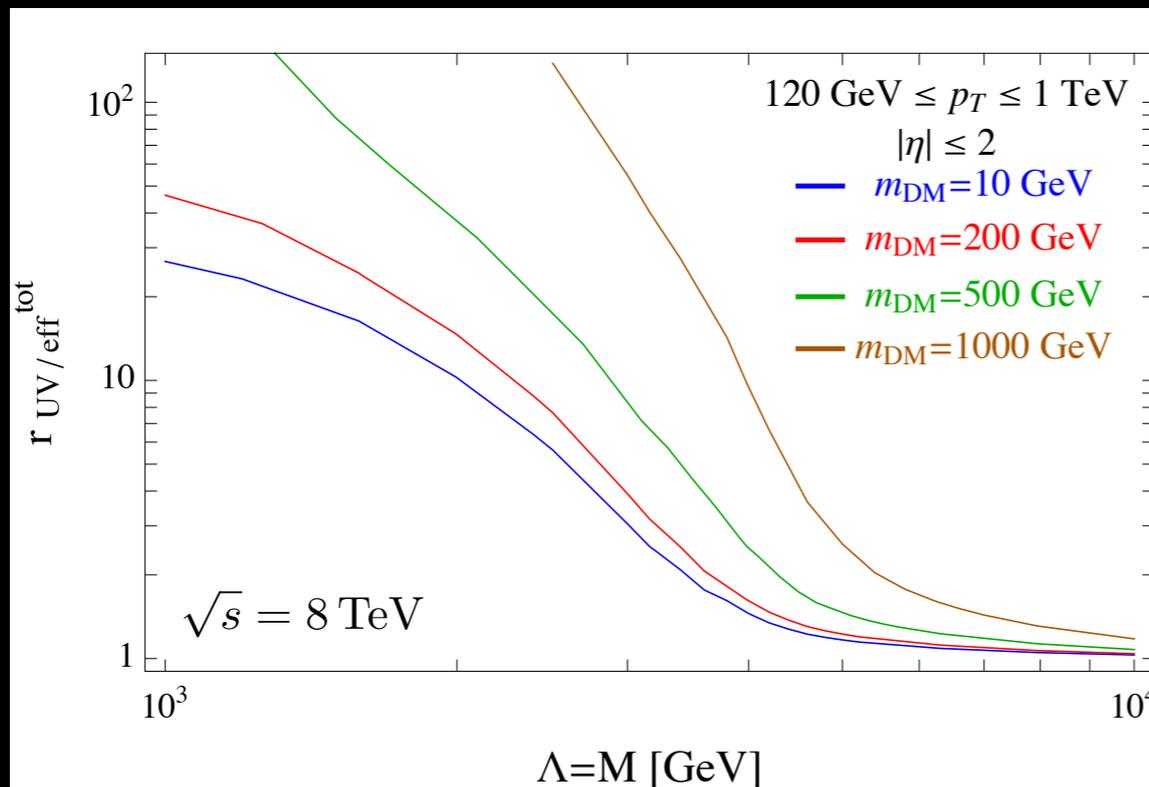
$$\frac{d^2\hat{\sigma}_{\text{UV}}}{dp_T d\eta} = \frac{\alpha_s}{36\pi^2} \frac{1}{p_T} \frac{g_q^2 g_\chi^2}{[Q_{\text{tr}}^2 - M^2]^2} \frac{[Q_{\text{tr}}^2 - 4m_{\text{DM}}^2]^{3/2}}{Q_{\text{tr}}} \left[ 1 + \frac{Q_{\text{tr}}^4}{(x_1 x_2 s)^2} \right]$$

1. What is the difference between interpreting data with an effective operator and with its UV completion?
2. In what regions of the parameter space  $\{\Lambda, m_{\text{DM}}\}$  is the effective description accurate/reliable?

# EFT vs UV COMPLETION

$$r_{\text{UV/eff}}^{\text{tot}} \equiv \frac{\sigma_{\text{UV}}|_{Q_{\text{tr}} < M}}{\sigma_{\text{eff}}|_{Q_{\text{tr}} < \Lambda}}$$

error of using EFT (truncated at dim-6)  
instead of full theory

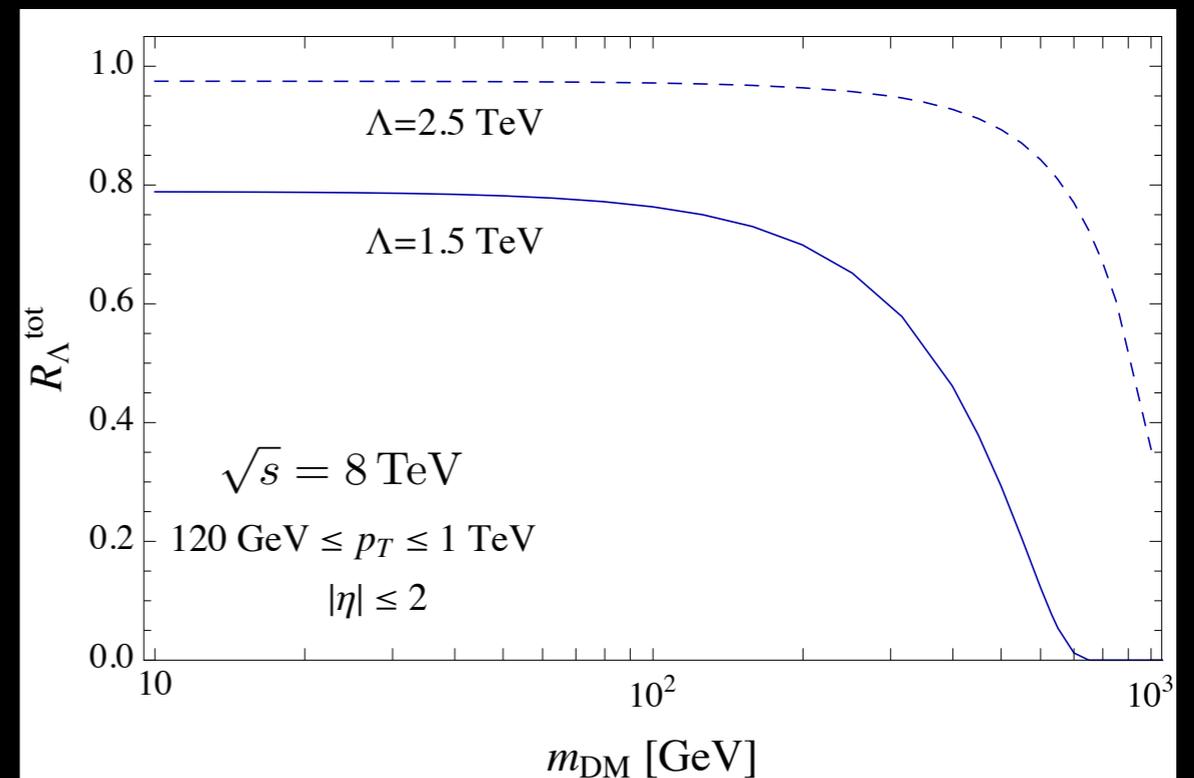
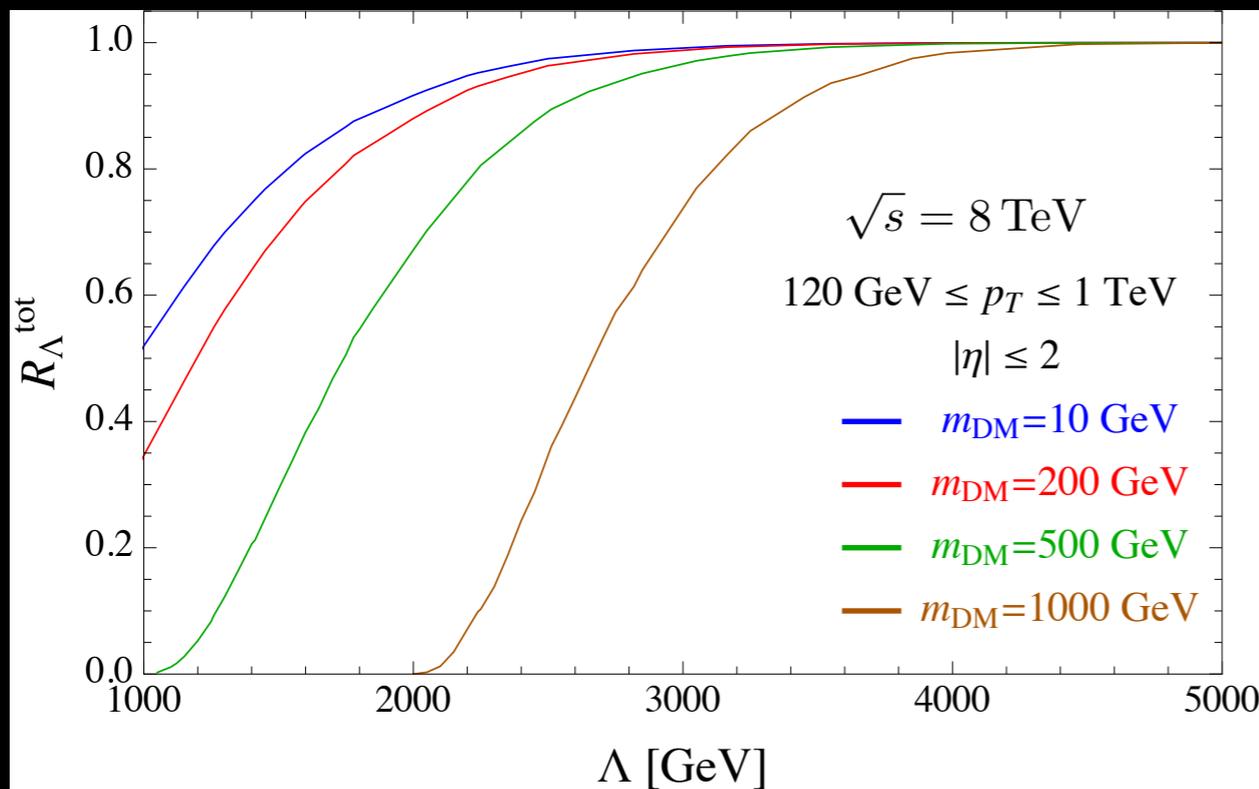


- $\sigma_{\text{UV}}$  easily larger than  $\sigma_{\text{eff}}$   $\longrightarrow$  mono-jet data can place more stringent bounds on mediator masses
- **direct** exclusion bounds from negative searches of heavy mediators? (e.g. di-jet searches)

# THE EFFECT OF THE EFT CUTOFF

$$R_{\Lambda}^{\text{tot}} \equiv \frac{\sigma_{\text{eff}} |_{Q_{\text{tr}} < \Lambda}}{\sigma_{\text{eff}}}$$

fraction of eff. cross section  
at low momentum transfer

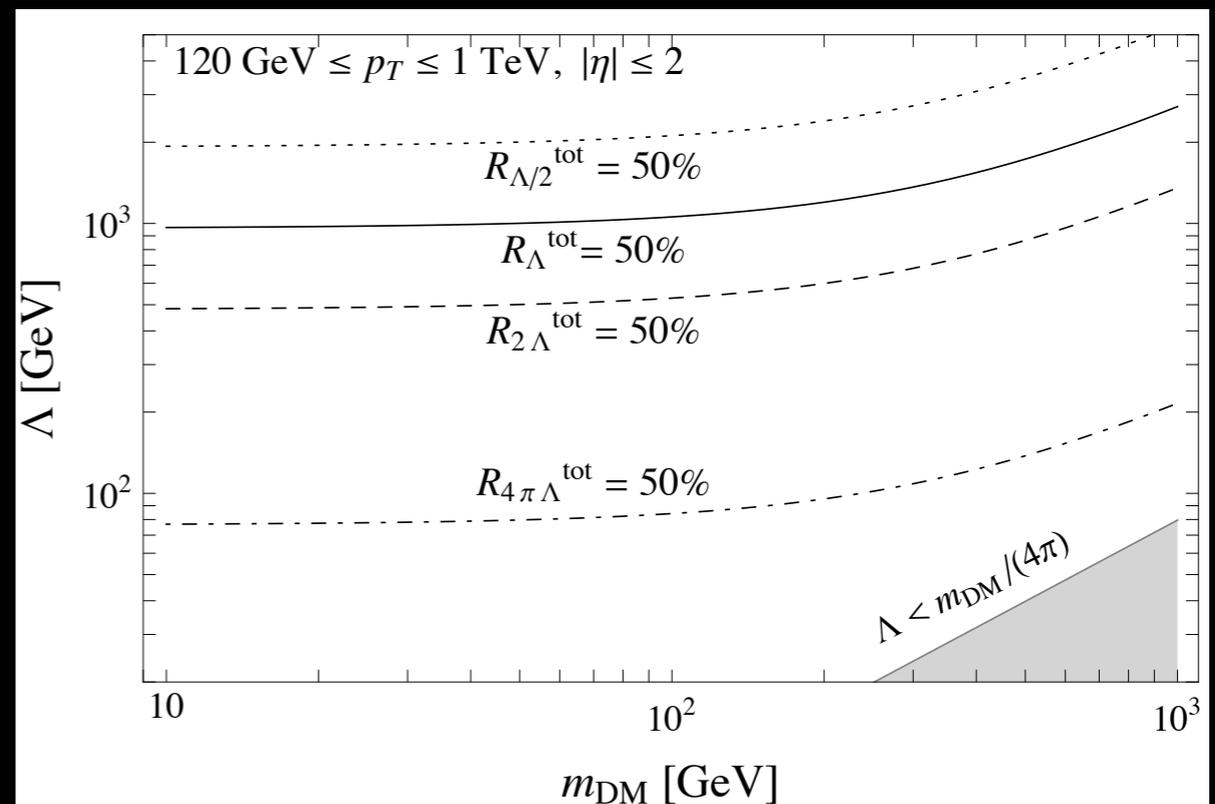
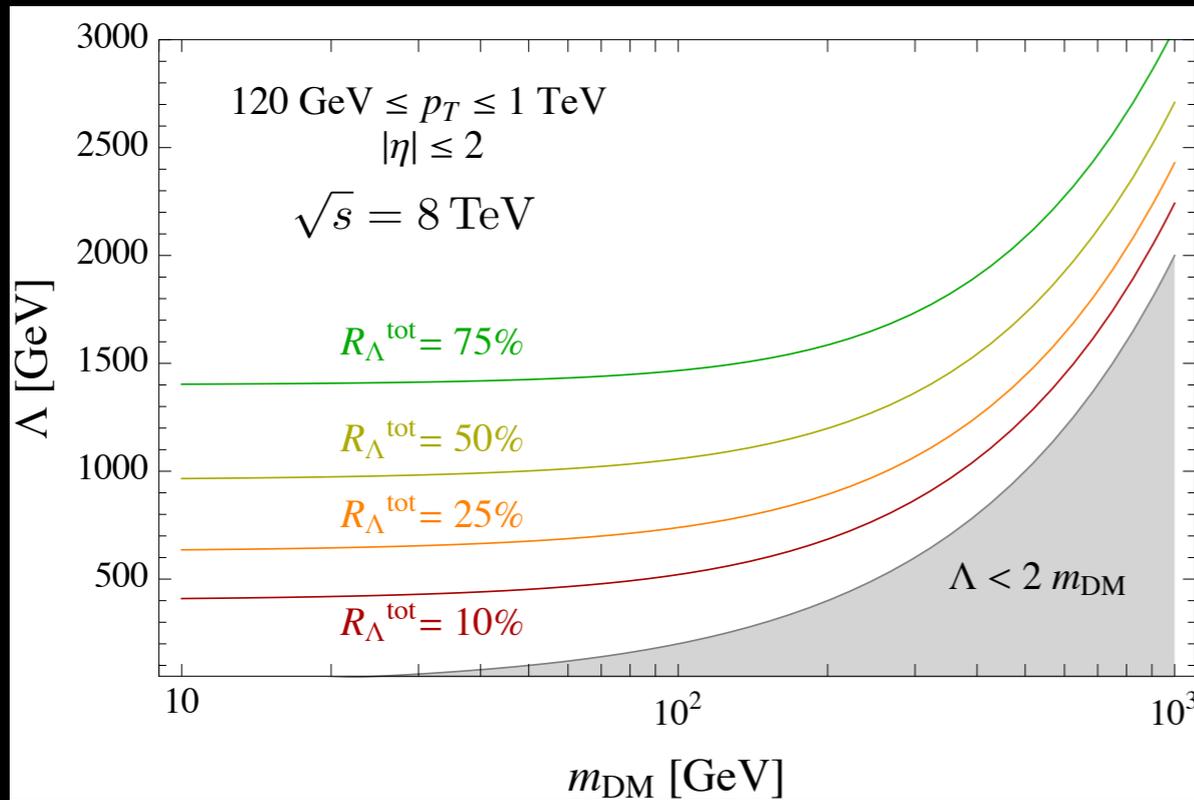


- Ratio  $\sim 1$ : negligible contribution from higher-dim ops. Accurate cross sections can be extracted without considering the cutoff on the momentum transfer.
- EFT works better for **larger  $\Lambda$**  and **smaller  $m_{\text{DM}}$**

# THE EFFECT OF THE EFT CUTOFF

$$R_{\Lambda}^{\text{tot}} \equiv \frac{\sigma_{\text{eff}} |_{Q_{\text{tr}} < \Lambda}}{\sigma_{\text{eff}}}$$

fraction of eff. cross section  
at low momentum transfer



- Cross sections are measured experimentally with  $\sim O(10\%)$  accuracy. Worry about EFT validity is justified.
- The precise definition of cutoff scale is somewhat arbitrary ( $\Lambda/2$ ,  $2\Lambda$ ?)  
**Most conservative choice:**  $Q_{\text{tr}} < 4\pi\Lambda$
- Numerical simulations: see next talk by J. Gramling.

1-1 correspondence:  
eff ops  $\longleftrightarrow$  simple toy models

- ✗ 1 or 2 more parameters ( $g$ 's)
- ✗ direct detection limits must be re-expressed

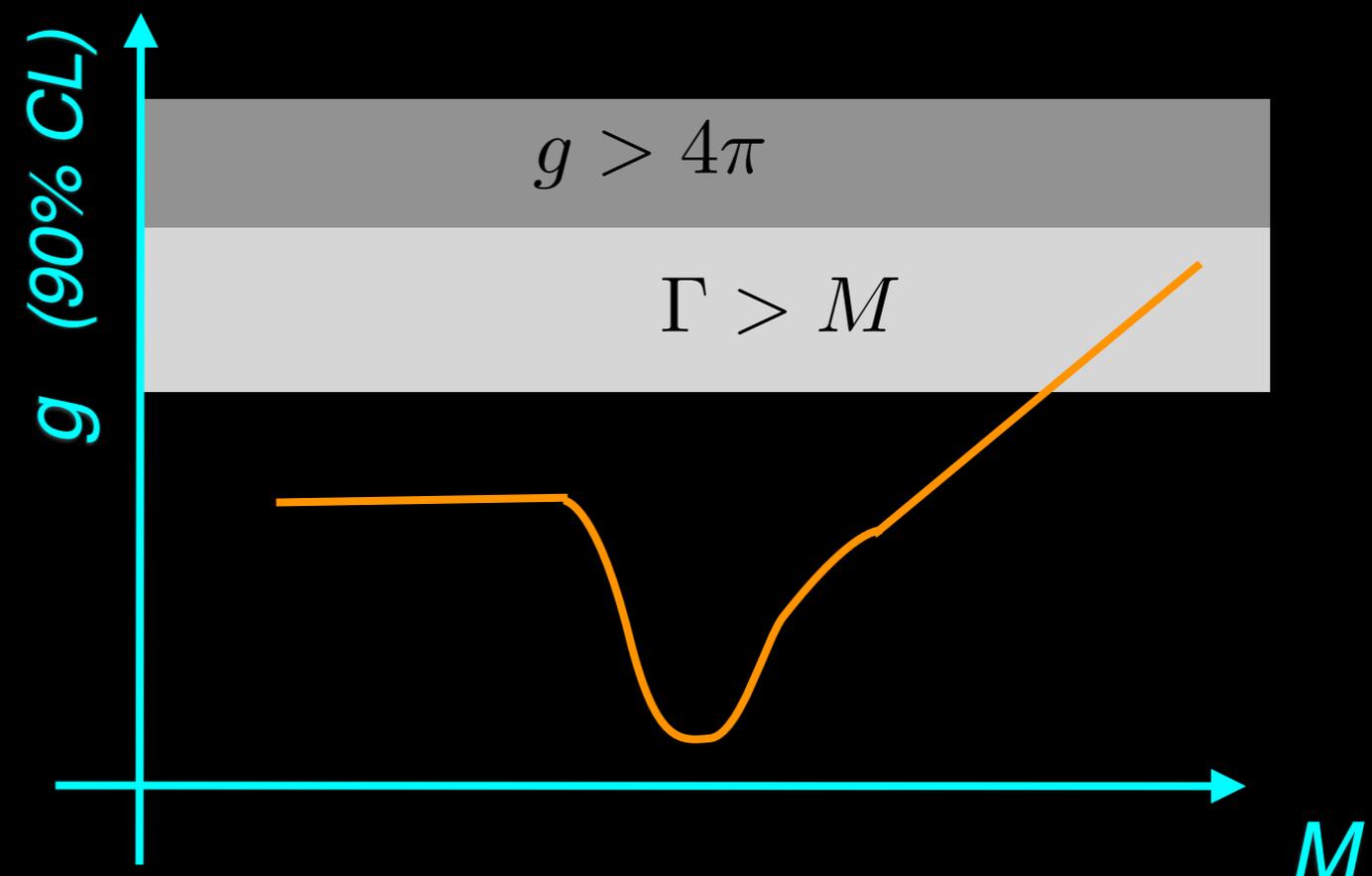
- ✓ exploit other searches for mediators (e.g. di-jet), complementary to mono-jet
- ✓ theoretically consistent, no worries about EFT, widths, etc.

## PROPOSAL

*provide upper limits on  $g$  (or  $M/g$ )*

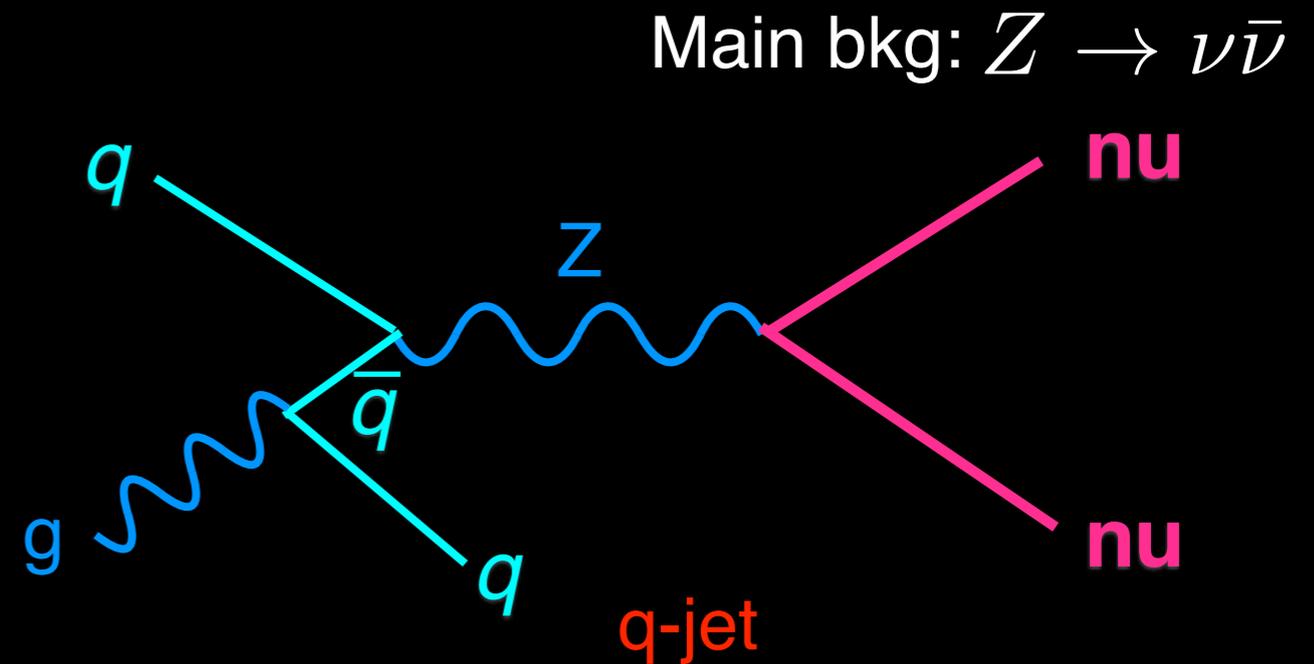
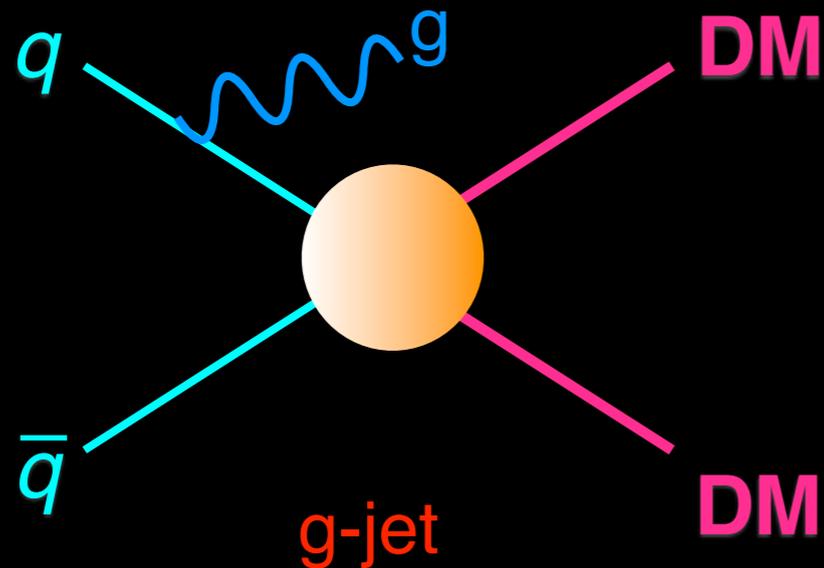
- for each simplified model
- for given  $m_{\text{DM}}$

→ complete and reliable information

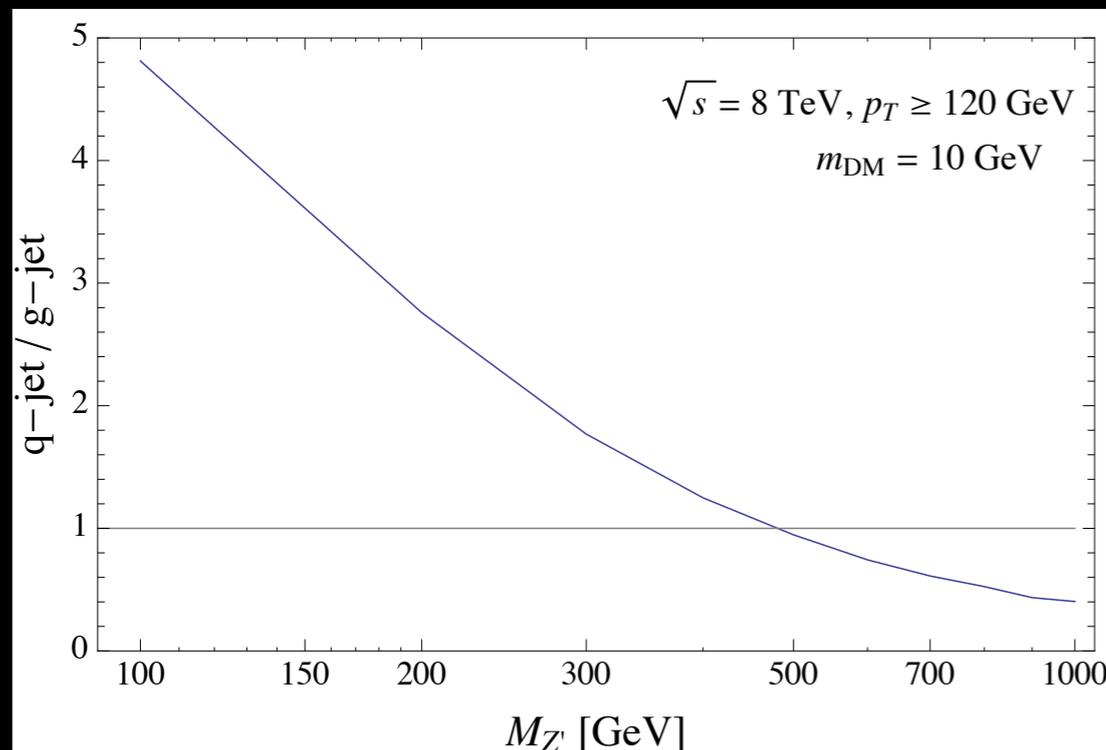


# QUARKS VS GLUON JETS?

Preferentially:



because of parton luminosities



- Ratio q-jet/g-jet can be:  $\sim 1/4$  for **SIGNAL**  
 $\sim 5$  for **BKG**
- Challenge: distinguish q/g jets
- This observation can help improving bounds with 14 TeV data

14 TeV data?

- stat and syst errors are already comparable
- expect stronger bounds on  $\Lambda$  but also larger  $Q_{\text{tr}}$

other ops?

- work in progress (analytically)
- for numerics, see Gramling's talk

$t$ -channel?

- experimental results only for contact  $s$ -channel interactions

Scalar DM?

- anybody working on it?

# CONCLUSIONS



LHC searches for DM using effective operators must be handled with care

without resorting to an explicit model,  
info about the validity of EFT  
can be extracted



take it into account when  
placing bounds

use explicit UV completions  
rather than EFT



stronger/complementary limits  
from direct searches of  
heavy mediators

LHC can discover mediators more  
easily than effective operators



TIME TO SHIFT FROM EFT  
TO SIMPLIFIED MODELS