

Introduction

The aim of this work is to show the applicability of the **Reduced Basis (RB)** model reduction in nonlinear systems undergoing bifurcations. Bifurcation analysis, i.e. following the different bifurcating branches and determining the **bifurcation points**, is a complex computational task. **Reduced Order Models (ROM)** can reduce the computational burden, enabling fast online evaluation of the solution for arbitrary parameter values. We represent the nonlinear PDE with the parametrized mapping $G: V \times \mathcal{D} \rightarrow V'$, so that the weak form reads: given $\lambda \in \mathcal{D}$, find $X(\lambda) \in V$ s.t.

$$g(X(\lambda), Y; \lambda) \doteq \langle G(X(\lambda); \lambda), Y \rangle = 0, \quad \forall Y \in V.$$

We say that $\lambda^* \in \mathbb{R}$ is a bifurcation point for G from the trivial solution, if there is a sequence $(X_n, \lambda_n) \in V \times \mathbb{R}$ with $X_n \neq 0$ and $G(X_n, \lambda_n) = 0$ such that $(X_n, \lambda_n) \rightarrow (0, \lambda^*)$. A necessary condition for λ^* to be a bifurcation point for G is that $D_X G(0; \lambda^*)$ is not invertible.

From the algebraic viewpoint the RB method reads: find $\delta \vec{X}_N \in \mathbb{R}^N$ such that

$$\mathbb{J}_N(\vec{X}_N^k(\lambda); \lambda) \delta \vec{X}_N = G_N(\vec{X}_N^k(\lambda); \lambda), \quad \text{and} \quad X_N^{k+1} = X_N^k - \delta X_N.$$

Von Kármán equations for plates

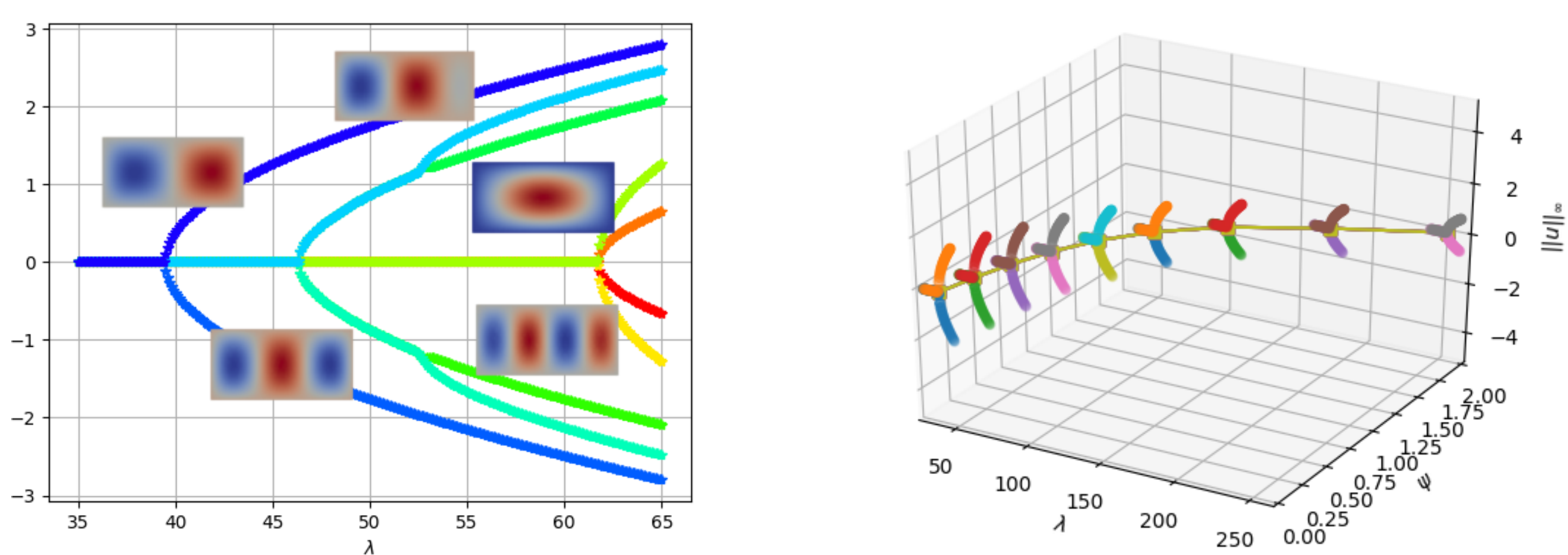
Consider a rectangular plate $\Omega = [0, L] \times [0, 1]$ in its undeformed state, subject to a λ -parametrized external load acting on its edges. The displacement u and the **Airy stress potential** ϕ satisfy the equations

$$\begin{cases} \Delta^2 u = [\lambda h + \phi, u] + f, & \text{in } \Omega \\ \Delta^2 \phi = -[u, u], & \text{in } \Omega \end{cases} \quad \begin{cases} u = \Delta u = 0, & \text{in } \partial\Omega \\ \phi = \Delta \phi = 0, & \text{in } \partial\Omega \end{cases}$$

where h and f are some given external forces acting on our plate, while

$$\Delta^2 := \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)^2, \quad [u, \phi] := \frac{\partial^2 u}{\partial x^2} \frac{\partial^2 \phi}{\partial y^2} - 2 \frac{\partial^2 u}{\partial x \partial y} \frac{\partial^2 \phi}{\partial x \partial y} + \frac{\partial^2 u}{\partial y^2} \frac{\partial^2 \phi}{\partial x^2},$$

are the biharmonic operator and the **Monge-Ampère bracket**, respectively.

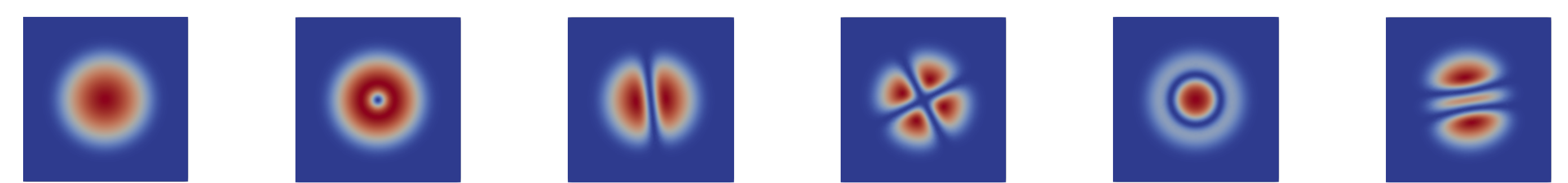


Schrodinger equations in quantum theory with A. Quaini (Houston)

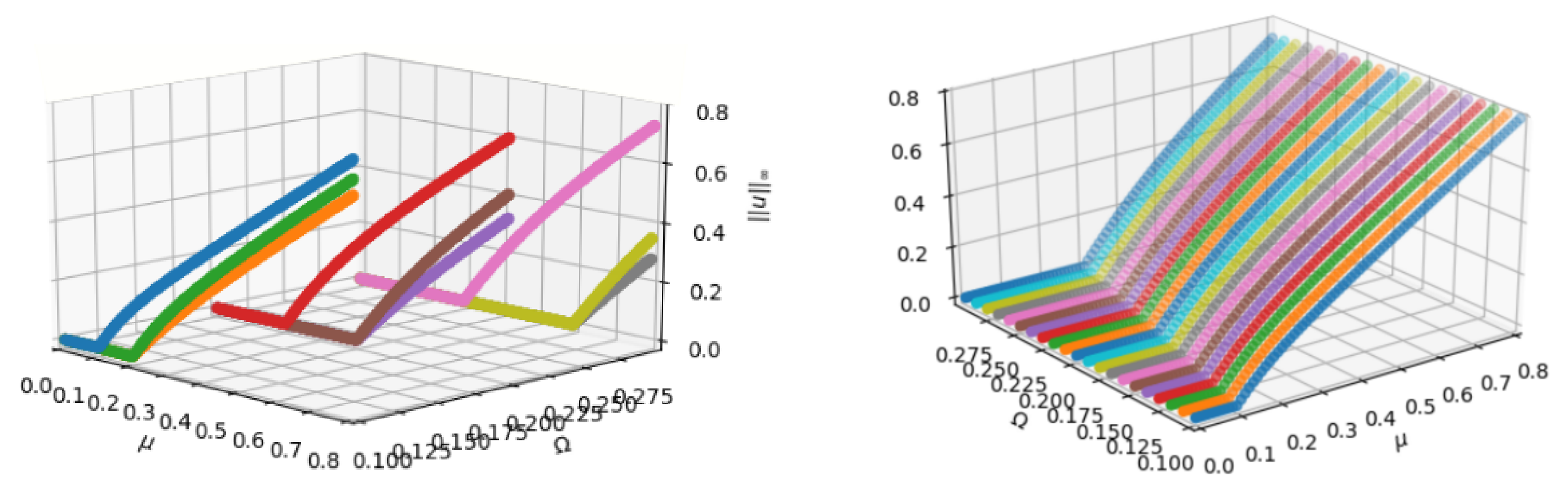
The **Gross-Pitaevskii** equation, which models certain classes of Bose-Einstein condensates, describes the ground state of a quantum system through the single-particle wave function $\phi: D \rightarrow \mathbb{C}$ which satisfies

$$-\frac{1}{2} \Delta \phi + |\phi|^2 \phi + \left(\frac{1}{2} \Omega^2 r^2 \right) \phi - \mu \phi = 0, \quad \text{in } D$$

where Ω is the trap strength and μ is the chemical parameter.



By applying the **Empirical Interpolation Method (EIM)**, it takes just $t_{RB} = 7$ s for the complete construction of the reduced basis bifurcation diagrams, with respect to the $t_{HF} = 246$ s for the high fidelity one.



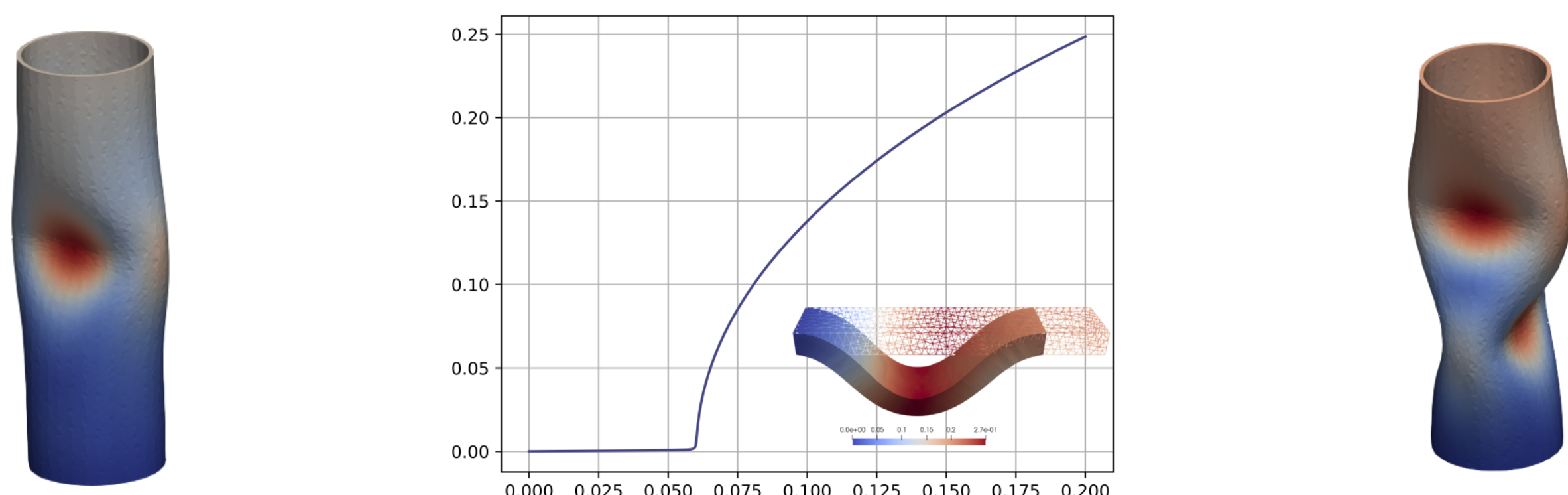
Hyperelastic beam in 2D and 3D cases

with A. Patera (MIT)

To study the deformation of an hyperelastic beam, we can transform it in a **minimization problem**, defining the potential energy as

$$\Pi(\mathbf{u}) = \int_{\Omega} \psi(\mathbf{u}) \, dx - \int_{\Omega} \mathbf{B} \cdot \mathbf{u} \, dx - \int_{\partial\Omega} \mathbf{T} \cdot \mathbf{u} \, ds,$$

where \mathbf{u} is the in plane displacement and $\psi(\mathbf{u})$ is the strain energy function, which becomes $\psi(\mathbf{u}) = \frac{\mu}{2}(I_c - 3) - \mu \ln J + \frac{\lambda}{2}(\ln J)^2$ with **Neo-hookean** law.



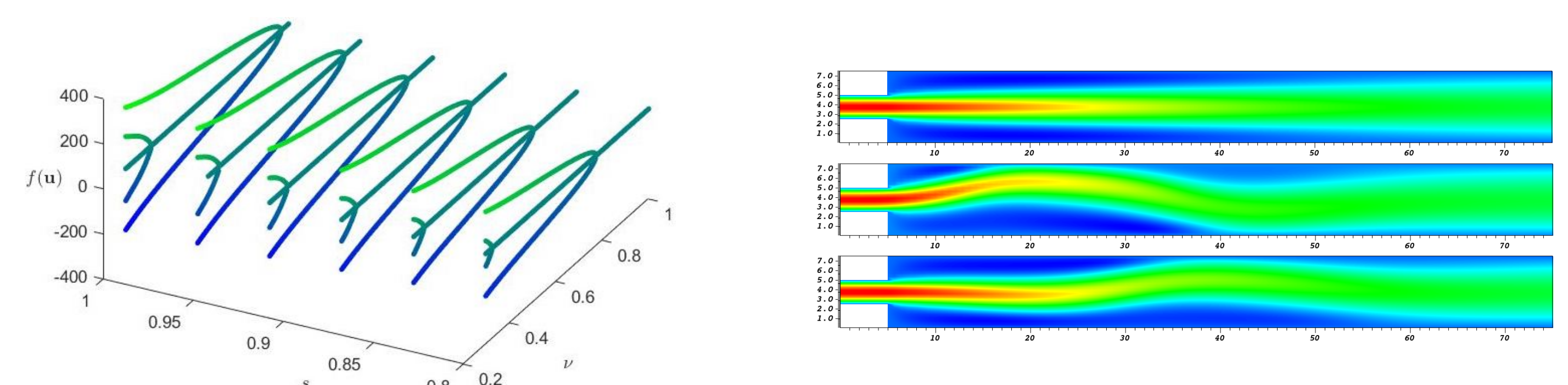
Navier-Stokes equations in CFD

with M. Hess (SISSA)

The test case is the study of a **viscous, steady and incompressible flow** in a planar straight channel with a narrow inlet, described by N-S equations

$$\begin{cases} -\nu \Delta \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p = \mathbf{0}, & \text{in } \Omega \\ \nabla \cdot \mathbf{u} = 0, & \text{in } \Omega \end{cases}$$

where \mathbf{u} is the velocity, p the pressure and ν the kinematic viscosity. The **multi-parameter case** is obtained by varying the inlet BC magnitude.



References

- [1] F. Pichi, A. Quaini, and G. Rozza. Reduced technique in bifurcating phenomena: application to the Gross-Pitaevskii equation. *In preparation*, 2019.
- [2] F. Pichi and G. Rozza. Reduced basis approaches for parametrized bifurcation problems held by nonlinear Von Kármán equations. *Journal of Scientific Computing*, 2019, in press. doi: 10.1007/s10915-019-01003-3.

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