

## Introduction

When dealing with **Parametrized Partial Differential Equations** the computational cost required by a large number of solutions for each new value of the involved parameters may be unaffordably large. To mitigate that, different methods have been studied in order to find solutions in a more efficient way ([1]).

In our work we exploit a **POD approach** to obtain the basis functions used to project the original problem and to reconstruct an approximate solution manifold. We focus on a new **reduced segregated approach** for laminar flows and on a **mixed projection-data driven technique** for turbulent flows, in a **Finite Volume** framework ([3]).

## Segregated methods: the reduced SIMPLE algorithm

Incompressible Navier-Stokes equations:

$$\begin{cases} (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{\nabla P}{\rho} = \nu \nabla^2 \mathbf{u} \\ \nabla \cdot (\rho \mathbf{u}) = 0 \end{cases}$$

Since the equations are **coupled** and block solvers require **big matrices** to be stored at each step, **segregated methods** employed, e.g., in SIMPLE solver, are to be preferred: momentum and continuity equations

are solved iteratively and one at a time. In order to be as **coherent** as possible, we developed a segregated algorithm also at the reduced level.

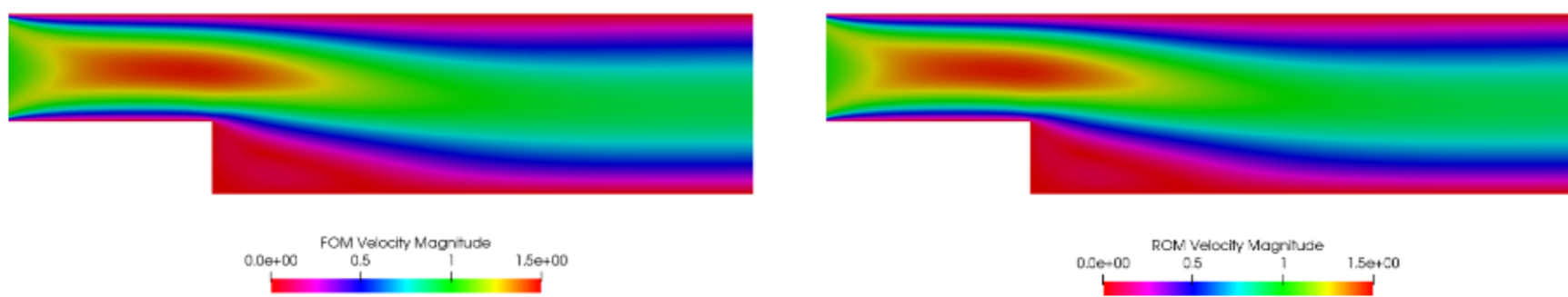
### The Reduced SIMPLE Algorithm

**Input:** Tentative value of the velocity and pressure coefficient vectors  $\mathbf{a}_0^u$  and  $\mathbf{a}_0^p$ ,  $k = 1$ , tolerance  $tol$ ,  $res = tol + 1$ .

- 1: Reconstruct the first attempt full velocity and pressure fields:  $\mathbf{u}_0 = L_u \mathbf{a}_0^u$ ,  $\mathbf{p}_0 = L_p \mathbf{a}_0^p$ ;
- 2: **while**  $res > tol$  **do**
- 3: assemble reduced momentum equation  $\mathbf{A}_u^r = L_u^T \mathbf{A}_u(\mathbf{u}_{k-1}, \mu) L_u$ ;  $\mathbf{b}_u^r = L_u^T \mathbf{b}_u(\mu, \mathbf{p}_{k-1})$ ;
- 4: compute reduced velocity residual  $\rightarrow r_u = |\mathbf{A}_u^r \mathbf{a}_{k-1}^u - \mathbf{b}_u^r|$ ;
- 5: solve  $\mathbf{A}_u^r \mathbf{a}_k^u = \mathbf{b}_u^r$  and reconstruct  $\mathbf{u}_k = L_u \mathbf{a}_k^u$ ;
- 6: assemble reduced pressure equation  $\mathbf{A}_p^r = L_p^T \mathbf{A}_p(\mu) L_p$ ;  $\mathbf{b}_p^r = L_p^T \mathbf{b}_p(\mu, \mathbf{u}_k)$ ;
- 7: compute reduced pressure residual  $\rightarrow r_p = |\mathbf{A}_p^r \mathbf{a}_{k-1}^p - \mathbf{b}_p^r|$ ;
- 8: solve  $\mathbf{A}_p^r \mathbf{a}_k^p = \mathbf{b}_p^r$  and reconstruct  $\mathbf{p}_k = L_p \mathbf{a}_k^p$ ;
- 9:  $res = \max(r_u, r_p)$ ,  $k = k + 1$ ;
- 10: **end while**

## Segregated methods: a parametrized viscosity problem

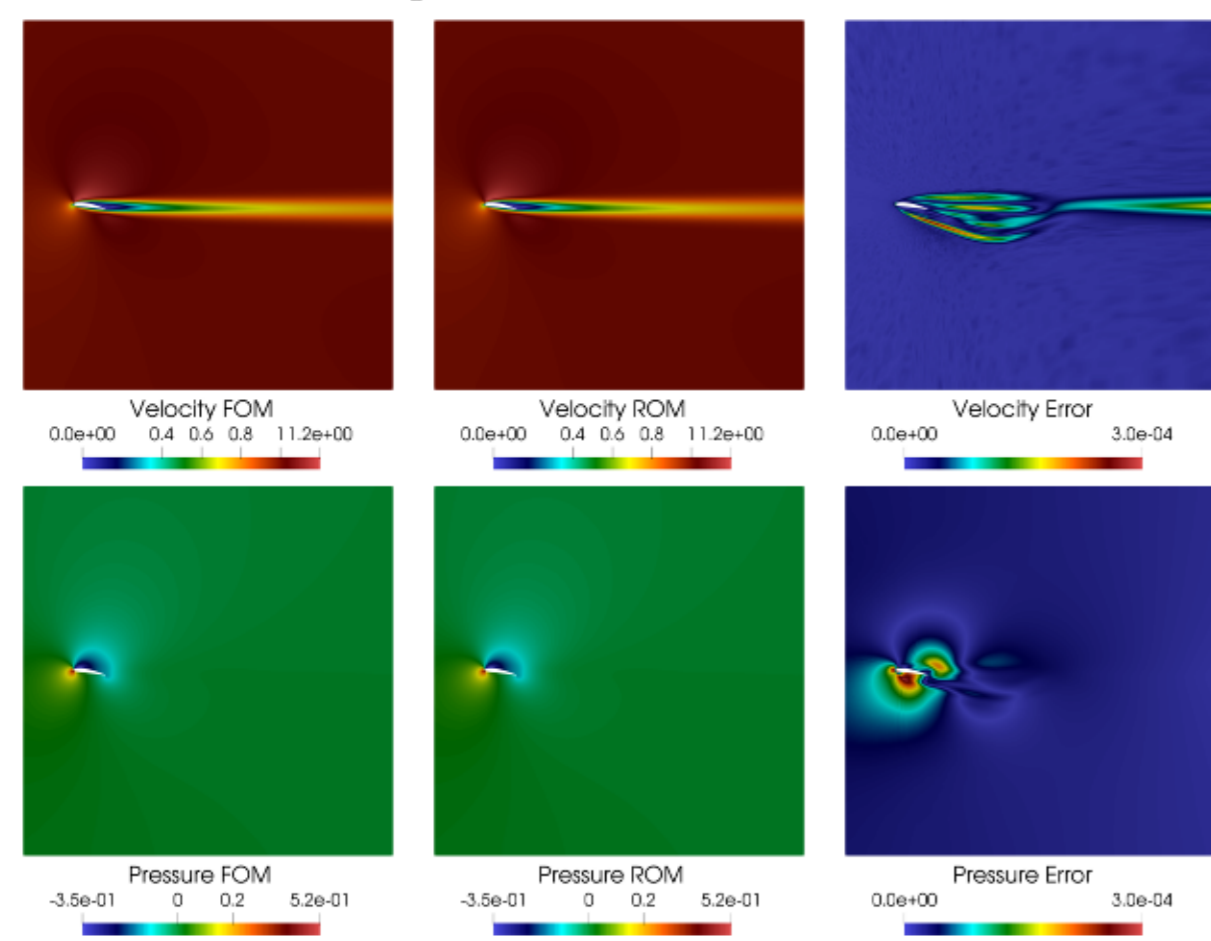
The geometry is represented by a back step and the parameter we have chosen is the **viscosity**  $\mu \in [0.01, 1]$ . 50 snapshots have been solved to apply the POD.



**Left figure:** full order velocity solution for  $\mu = 0.024$ . **Right figure:** reduced counterpart for the same value of the parameter. Only 10 basis functions have been used during projection for both velocity and pressure fields. **Bottom figure:** comparison between the  $L^2$  norm of the error for each value of the parameter in the training set for the old block ROM solver (red) and for the SIMPLE ROM solver (blue).

## Segregated methods: a parametrized geometry problem

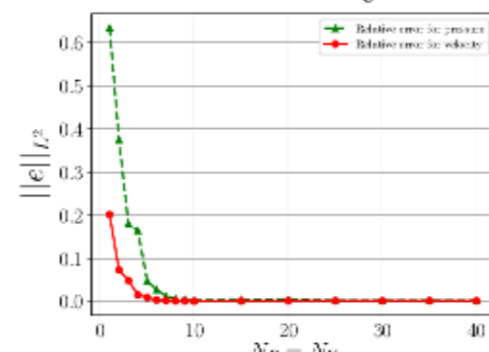
One of the in-progress works in the group is relative to **geometrical parametrization** and **hyper-reduction** ([4]). In this case the **angle of attack** of an airfoil has been parametrized:  $\alpha \in [-10, 10]$ . 100 snapshots have been solved, 40 basis functions have been employed for the projection. The modified mesh is obtained by the use of radial basis functions interpolation.



**Top figures:** from left to right, FOM and ROM velocity solutions and velocity error magnitude for  $\alpha = 9.15$ .

**Bottom figures:** from left to right, FOM and ROM pressure solutions and pressure error for  $\alpha = 9.15$ .

**Plot:** errors decay.



## Turbulent flows: mixing projection and data-driven

The starting point in developing the ROM is the usual decomposition of the fields into a sum of global spatial modes multiplied by temporal coefficients:

$$\begin{aligned} \bar{\mathbf{u}}(\mathbf{x}, t; \boldsymbol{\mu}) &\approx \sum_{i=1}^{N_u} a_i(t; \boldsymbol{\mu}) \phi_i(\mathbf{x}), & \bar{p}(\mathbf{x}, t; \boldsymbol{\mu}) &\approx \sum_{i=1}^{N_p} b_i(t; \boldsymbol{\mu}) \chi_i(\mathbf{x}), \\ \nu_t(\mathbf{x}, t; \boldsymbol{\mu}) &\approx \sum_{i=1}^{N_{\nu_t}} g_i(t; \boldsymbol{\mu}) \eta_i(\mathbf{x}). \end{aligned}$$

The reduced basis  $\phi_i(\mathbf{x})$ ,  $\chi_i(\mathbf{x})$  and  $\eta_i(\mathbf{x})$  are computed by means of **Proper Orthogonal Decomposition (POD)**, which is implemented using the snapshots method. The velocity snapshots matrix  $\mathbf{S}_u$  is given by:

$$\mathbf{S}_u = \{\bar{\mathbf{u}}(\mathbf{x}, t_1; \boldsymbol{\mu}_1), \dots, \bar{\mathbf{u}}(\mathbf{x}, t_{N_T}; \boldsymbol{\mu}_M)\} \in \mathbb{R}^{N_u^h \times N_s}.$$

We proceed to the projection step of the momentum equation of RANS. This step will give the following dynamical system with the unknowns being the vectors of coefficients  $\mathbf{a}$  and  $\mathbf{b}$ :

$$M \dot{\mathbf{a}} = \nu(\mathbf{B} + \mathbf{B}_T) \mathbf{a} - \mathbf{a}^T \mathbf{C} \mathbf{a} + \mathbf{g}^T (\mathbf{C}_{T1} + \mathbf{C}_{T2}) \mathbf{a} - \mathbf{H} \mathbf{b}.$$

In order to be able to use the continuity equations we used the **supremizer** enrichment method which allows to project the continuity equation onto the pressure modes, the resulting system is the following

$$\begin{cases} M \dot{\mathbf{a}} = \nu(\mathbf{B} + \mathbf{B}_T) \mathbf{a} - \mathbf{a}^T \mathbf{C} \mathbf{a} + \mathbf{g}^T (\mathbf{C}_{T1} + \mathbf{C}_{T2}) \mathbf{a} - \mathbf{H} \mathbf{b}, \\ \mathbf{P} \mathbf{a} = \mathbf{0}. \end{cases} \quad (1)$$

**Interpolation using Radial Basis Functions (RBF)** is used to approximate the value of the eddy viscosity coefficients vector  $\mathbf{g}$  [2]. First one can notice the following:

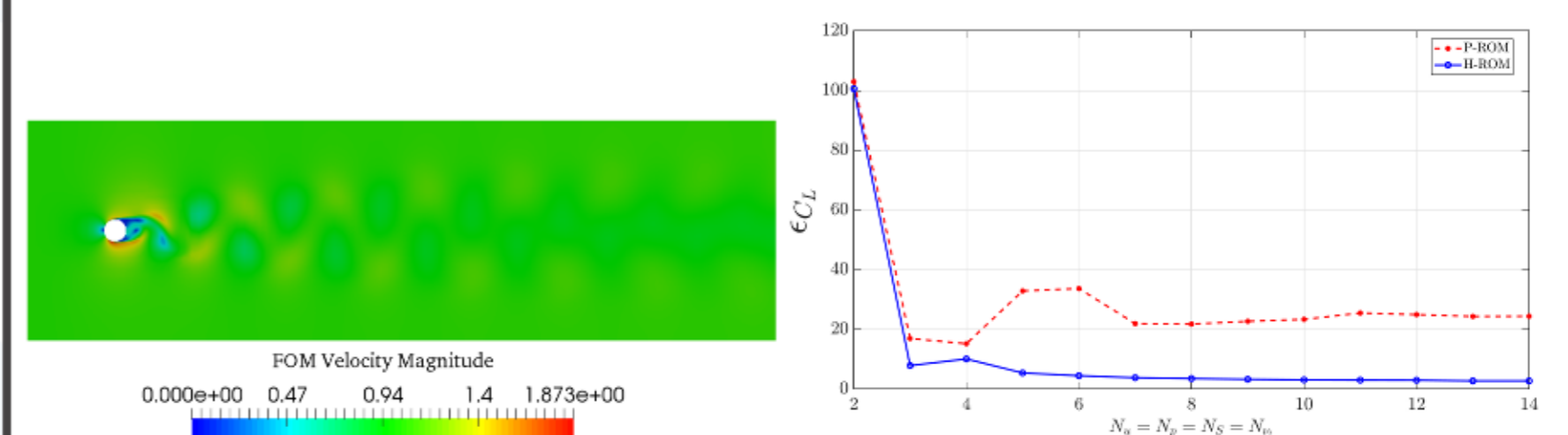
$$\nu_t = f(\bar{\mathbf{u}}) \rightarrow \mathbf{g} = \tilde{f}(\mathbf{a})$$

The interpolation using RBF functions is based on the following formula :

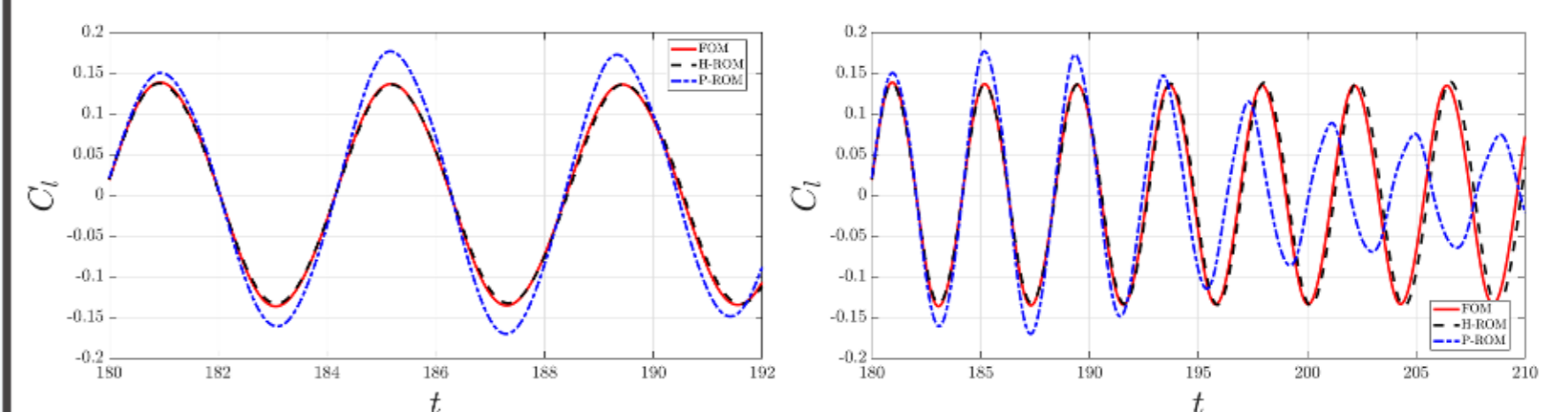
$$g_L(\mathbf{a}) \approx G_L(\mathbf{a}) = \sum_{j=1}^{N_s} w_{L,j} \zeta_{L,j}(\|\mathbf{a} - \mathbf{a}_{L2}^j\|_{\mathbb{R}^{N_u}}), \quad \text{for } L = 1, 2, \dots, N_{\nu_t},$$

## Turbulent flows: results for a circular cylinder

The presented results are for the benchmark case of the flow around a circular cylinder in unsteady state setting. The case is without parameterization so the reduction is done on time (both **reproduction** of the snapshots and **extrapolation** in time). The **Reynolds number** is equal to 10000. The results shown are for the **lift coefficient**  $C_l$ . The FOM results are compared to those of both the **Hybrid ROM (H-ROM)** and the **Projection ROM (P-ROM)** which is based on solving (1) neglecting the turbulent term in RANS equations. A quantitative convergence analysis is shown for the decay of the error with the increase of the number of modes used in the **H-ROM**.



**Left figure:** velocity magnitude FOM field. **Right figure:** lift coefficients  $L^2$  error curve (reproduction case).

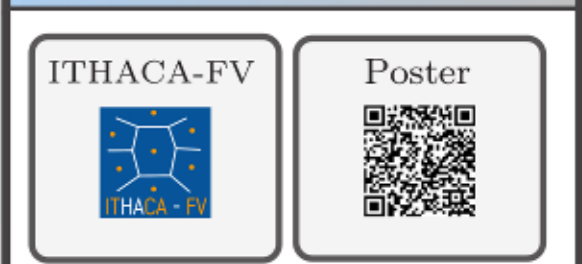


**Left figure:** lift coefficients curves for the FOM, the H-ROM and the P-ROM (reproduction case), 8 modes used. **Right figure:** lift coefficients curves for the FOM, the H-ROM and the P-ROM (extrapolation case), 8 modes used.

## References

- [1] J. S. Hesthaven, G. Rozza, and B. Stamm. *Certified Reduced Basis Methods for Parametrized Partial Differential Equations*. MS&A Series. Springer International Publishing, 2016.
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- [3] G. Stabile and G. Rozza. Finite volume POD-Galerkin stabilised reduced order methods for the parametrised incompressible Navier-Stokes equations. *Computers & Fluids*, 173:273–284, 2018.
- [4] G. Stabile, M. Zancanaro, and G. Rozza. Efficient Geometrical parametrization for Finite-Volume based Reduced Order Methods. *arXiv: https://arxiv.org/abs/1901.06373*, 2019.

## Utilities



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